

17 **Abstract**

18 Deep convection is the primary influence on weather and climate in tropical regions.
 19 However, understanding and simulating the shallow-to-deep (STD) convective transition
 20 has long been challenging. Here, we conduct high-resolution numerical simulations to
 21 assess the environmental controls on the evolution of isolated convection in the Amazon
 22 during the wet season. Observations and large-scale forcing derived through the constrained
 23 variational analysis approach for the GoAmazon2014/5 experiments are used in the sim-
 24 ulations and model validation. The model consistently reproduces the GOAmazon ob-
 25 servations for precipitation, moisture, and surface fluxes of radiation, latent and sensi-
 26 ble heat. Through sensitivity experiments, we examine the relative importance of mois-
 27 ture and vertical wind shear in controlling the STD convective transition. Reducing the
 28 pre-convective humidity within the lower 1.5 km significantly suppresses vertical devel-
 29 opment and lowers the ice water path. Additionally, the maximum precipitation rate de-
 30 creases almost quadratically with column water vapor. Conversely, a reduction of col-
 31 umn water vapor above 1.5 km by a factor of two or more is necessary to produce a com-
 32 parable decrease in ice water path or precipitation. Moderate low-level wind shear fa-
 33 cilitates the STD transition, leading to an earlier peak of ice water compared to stronger
 34 wind shear or its absence. Although upper-level wind shear negatively influences high
 35 cloud formation, its role in controlling the STD transition is relatively smaller than that
 36 of low-level wind shear. Our results help quantify the role of moisture and wind shear
 37 on the STD transition, but also suggest that dynamic factors may exert a more pronounced
 38 influence.

39 **Plain Language Summary**

40 The Amazon rainforest plays a vital role in the Earth’s climate system. However,
 41 it is not entirely understood how environmental conditions control the evolution from
 42 fair weather conditions to severe thunderstorms in regions of the deep Tropics. We ad-
 43 dress this problem utilizing numerical simulations that capture the interactions between
 44 the forest, atmosphere, and clouds. Atmospheric modeling data developed for the GoA-
 45 mazon2014/5 experiment are used to initialize our Amazon-based simulations. The model
 46 consistently reproduces the Amazon environment throughout the period of our simula-
 47 tions, which covers December 2014. Additionally, we contrast the model results between
 48 the control simulation and experiments in which the moisture or wind is modified to eval-
 49 uate their relative importance to cloud development and precipitation. Lower tropospheric
 50 moisture is critical to cloud growth. The amount of moisture in the air above 1.5 km has
 51 a minor contribution to cloud development and precipitation. Low-level wind of mod-
 52 erate strength facilitates cloud development during the afternoon. The upper-level wind
 53 negatively affects the ice formation in high clouds. These results help strengthen our knowl-
 54 edge of tropical convection, critical for improving numerical model performance.

55 **1 Introduction**

56 Deep convection dominates the weather and climate in the tropics. Nevertheless,
 57 comprehending and simulating the convective processes is a formidable challenge due to
 58 the wide range of spatial and temporal scales involved (Mapes et al., 2009; Moncrieff et
 59 al., 2012; Zhang et al., 2013). Shallow cumulus convection, a small-scale phenomenon
 60 lasting tens of minutes and covering spatial scales of the order of a few kilometers, of-
 61 tentimes evolves into deep convective clouds covering tens of kilometers within typical
 62 time scales of 2 to 4 hours (Wu et al., 2009; Hohenegger & Stevens, 2013; Adams et al.,
 63 2013; Henkes et al., 2021; Powell, 2022). Moreover, deep convection frequently becomes
 64 organized and experiences upscale growth into mesoscale convective systems (MCSs) with
 65 lifetimes spanning hours to a day and ranging in horizontal scale from 100 km to 1,000
 66 km (Houze Jr, 2004). Likewise, land-atmosphere interactions and complex physical pro-

67 cesses ranging from cloud microphysics to the generation of gravity waves are intrinsi-
68 cally tied to deep convection (Silva Dias et al., 2002; Mapes et al., 2006; Mapes & Neale,
69 2011; Jewtoukoff et al., 2013; Gupta et al., 2023).

70 General circulation models (GCMs) rely on parameterizations of convective pro-
71 cesses and typically struggle to reproduce the shallow-to-deep (STD) convective tran-
72 sition over continental regions (Betts, 2002; Betts & Jakob, 2002; Bechtold et al., 2004;
73 Grabowski et al., 2006). Their simulated precipitation peaks much earlier than observed
74 (Lin et al., 2000; Betts, 2002; Collier & Bowman, 2004; Dai & Trenberth, 2004), which
75 is an important source of bias and uncertainty in GCMs to this day (Sherwood et al.,
76 2014; Stevens & Bony, 2013; Icterly et al., 2018; Maher et al., 2018; Freitas et al., 2020,
77 2024). To circumvent the inherent challenges posed by convective parameterizations, cloud-
78 resolving models (CRMs), which explicitly resolve the up- and downdrafts in clouds, have
79 been used to study convective processes over continental and oceanic regions. For ex-
80 ample, M. Khairoutdinov and Randall (2006) conducted the first high-resolution numeri-
81 cal simulations to investigate the STD transition over the Amazon. Their findings high-
82 lighted the importance of cold pools in forcing the development of deep convection, while
83 the impact of vertical wind shear and free tropospheric preconditioning were relatively
84 minor. As part of the EUROCS (EUROpean Cloud Systems study), Derbyshire et al.
85 (2004) evaluated the sensitivity of cumulus convection to free tropospheric humidity. Un-
86 like M. Khairoutdinov and Randall (2006), they observed intense deep precipitating con-
87 vection in moister scenarios, whereas only shallow convection was evident in the driest
88 scenario. Waite and Khouider (2010) conducted idealized numerical simulations over the
89 tropical Atlantic Ocean. Their study emphasized the importance of congestus precon-
90 ditioning, which reduces the impact of entrainment on cloud buoyancy, ultimately lead-
91 ing to the STD transition. In contrast, Hohenegger and Stevens (2013) showed that the
92 transition from congestus to deep convective clouds occurs on shorter time scales than
93 required for congestus clouds to moisten the atmosphere sufficiently. This implies that
94 dynamic factors play a more substantial role in driving convection. While CRM stud-
95 ies offer valuable insights into physical convective processes, they still require validation
96 through high-resolution observations, which have typically been lacking in tropical rain-
97 forests.

98 In the Amazon, important, but often limited field campaigns, have explored dif-
99 ferent aspects of tropical convection. Adams et al. (2015) established the Amazon Dense
100 GNSS Meteorological Network, a one-year campaign to observe the interaction between
101 water vapor fields and deep convection. Adams et al. (2013) also utilized GNSS/GPS
102 data from a long-term single site (July 2008 to December 2011) in Manaus, Brazil to eval-
103 uate the water vapor convergence associated with the STD transition and found a weak
104 and quasi-linear convergence timescale of approximately 8 hours, followed by a robust
105 and non-linear convergence timescale of approximately 4 hours during the STD transi-
106 tion. Later, Adams et al. (2017) employed this dense network data to investigate how
107 vapor fields evolve spatially during during the STD transition. Their results were con-
108 sistent with the 4-hour STD timescale and the spatial evolution was reflective of the wa-
109 ter vapor convergence posited in the single site study. More recently, the Green Ocean
110 Amazon (GOAmazon) 2014/5 Experiment (Martin et al., 2016, 2017) was carried out
111 from 2014 to 2015 in the central Amazon, providing the most comprehensive set of ob-
112 servations of clouds and aerosols in the Amazon to date. Analyzing this dataset during
113 the dry season (June-September), Ghate and Kollias (2016) noted an excess of water va-
114 por above 2 km during the early morning when contrasting locally-driven precipitating
115 days and nonprecipitating days. Conversely, Zhuang et al. (2017) and Tian et al. (2021)
116 observed that deep convective days exhibit relatively higher moisture extending from the
117 surface to mid-levels in all seasons. Schiro et al. (2016) showed a robust correlation be-
118 tween total column water vapor and precipitation in both the central Amazon and the
119 tropical western Pacific. Furthermore, Schiro and Neelin (2019) demonstrated a strong
120 connection between the initiation and likelihood of daytime precipitation and the bound-

121 ary layer and lower free troposphere moisture content. Previous studies do not completely
 122 agree on the relative importance of vertical wind shear. For example, while Zhuang et
 123 al. (2017) indicated that more intense low-level and deep-layer bulk wind shears facil-
 124 itate the STD transition during the dry season (June–September), Chakraborty et al. (2018)
 125 suggested that a more intense low-level shear could inhibit deep convection during the
 126 transition season (August–November), especially if it increases the entrainment of dry
 127 air.

128 In this paper, we conduct high-resolution model simulations to assess the role of
 129 moisture and vertical wind shear in controlling the STD convective transition in the Ama-
 130 zon. First, we focus on model validation employing GoAmazon data for the period of
 131 December 2014. Then, we conduct idealized sensitivity experiments in which either mois-
 132 ture or large-scale wind are modified at different atmospheric levels to assess their rel-
 133 ative importance in the development of deep convection. The paper is structured as fol-
 134 lows: Section 2 shows the study area. Section 3 describes the material and methods. Sec-
 135 tion 4 covers the model validation. Sensitivity experiments for moisture and wind shear
 136 are conducted in section 5. A discussion of the results is given in section 6. Section 7 con-
 137 tains the conclusions.

138 2 Study Region

139 The Amazon Basin is bordered by significant altitudes (Figure 1a), primarily in
 140 the western region, where some peaks in the Andes Mountains rise well over 6,000 m-
 141 eters in elevation above sea level. However, the simulations are conducted over the GoA-
 142 mazon2014/5 campaign region in the central Amazon, where the topography can be ad-
 143 equately considered as an extensive plain with minimal variations (< 130 m in our do-
 144 main of interest, section 3.2). During the experiment, most of the observations were taken
 145 from the T3 site, located 70 km downwind of Manaus, in Manacapuru (3.21°S , 60.60°W),
 146 a site characterized by a pasture surrounded by forest and close to the intersection of
 147 the Solimões River and Negro River (Figure 1b).

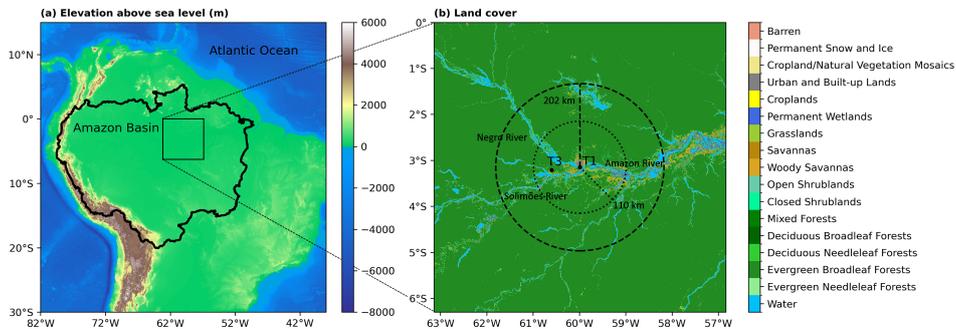


Figure 1. (a) Land topography and ocean depth (NOAA National Centers for Environmental Information, 2022) around the Amazon. (b) Land cover (Friedl et al., 2010) around the GoAmazon2014/5 sites. The dashed circle with a radius of 202 km centered at the T1 site (in Manaus) corresponds to the S-band radar domain. The dotted circle with a radius of 110 km shows the domain of the large-scale forcing developed for the GoAmazon2014/5 Experiment (Tang et al., 2016). We also indicate the Amazon, Solimões, and Negro Rivers on the map. Land cover is from 2014, based on the Moderate Resolution Imaging Spectroradiometer (MODIS) - International Geosphere-Biosphere Programme (IGBP) land cover classification system. The Amazon Basin contour is provided by Mayorga et al. (2012).

148 3 Material and Methods

149 3.1 Data

150 For model validation, we use campaign observations of precipitation, moisture, ra-
 151 diation, and surface latent and sensible heat fluxes. Precipitation is based on the SIPAM
 152 S-band radar measurements (Schumacher & Funk, 2018), which we average over the do-
 153 main of the control runs. Sensible heat flux (H) and latent heat flux (LE) are from the
 154 Quality Controlled Eddy Correlation (QCECOR) Flux Measurement (ARM, 2014b). Sur-
 155 face radiation fluxes are from the Sky Radiation Radiometers (SKYRAD) and Ground
 156 Radiation Radiometers (GNDRAD) (ARM, 2013). Column water vapor (CWV) is cal-
 157 culated from the balloon-borne sounding system (SONDE), which provides the vertical
 158 profiles of thermodynamic conditions 4 times per day during the period of this study,
 159 at 02, 08, 14, and 20 LST (ARM, 2014a).

160 Large-scale atmospheric fields of water vapor mixing ratio, temperature, wind, and
 161 moisture and temperature tendencies are based on 3-hour Constrained Variational Anal-
 162 ysis Data (VARANAL). This assimilation product was developed using atmospheric fields
 163 from the European Centre for Medium-Range Weather Forecasts (ECMWF) ERA-Interim
 164 reanalysis (Dee et al., 2011), which were mainly constrained by the SIPAM S-band radar
 165 precipitation rate and ARM surface fluxes through the column heat and moisture bud-
 166 get analysis (Tang et al., 2016). The VARANAL data represent an average over the anal-
 167 ysis domain centered at T1 site, covering a radius of 110 km (Figure 1b).

168 The Moderate Resolution Imaging Spectroradiometer (MODIS) provides data prod-
 169 ucts of land cover type (MCD12Q1 Version 6) and leaf area index (MCD15A2H Version
 170 6.1) (Friedl et al., 2010). Specifically, we use the land cover product based on the Inter-
 171 national Geosphere-Biosphere Programme (IGBP) land cover classification system. Silt,
 172 clay, and sand content in the soil are based on in-situ measurements of the soil type “Terra
 173 Firme” (Terra-firma) described in Table 1 on Schaefer et al. (2017). Soil temperature
 174 and wetness are based on the NASA Global Land Data Assimilation System (GLDAS)
 175 Noah Land Surface Model L4 3 hourly 0.25 x 0.25 degree V2.1 (Rodell et al., 2004).

176 3.2 Model Configuration

177 We employ the numerical model System for Atmospheric Modeling (SAM), ver-
 178 sion 6.11.8, which solves the anelastic equations of motion and uses liquid water static
 179 energy, total nonprecipitating, and precipitating water as thermodynamic prognostic vari-
 180 ables (M. F. Khairoutdinov & Randall, 2003). The equations are solved using lateral pe-
 181 riodic boundary conditions. A prognostic turbulent kinetic energy 1.5-order closure scheme
 182 is used to parameterize subgrid-scale effects. Different microphysics parameterizations
 183 are available, including the single-moment (Morrison, 2003), double-moment (Morrison
 184 et al., 2005), and Predicted Particle Properties (P3) (Morrison & Milbrandt, 2015) schemes.
 185 The radiative heating can be prescribed or calculated by choosing a radiation scheme,
 186 either the Community Atmosphere Model (CAM3) (Collins et al., 2006) or the Rapid
 187 Radiative Transfer Model (RRTM) (Mlawer et al., 1997) schemes. Surface fluxes can be
 188 prescribed or calculated using Monin-Obukhov similarity theory or a simplified Land Sur-
 189 face Model (LSM) (Lee & Khairoutdinov, 2015), which is only compatible with the CAM3
 190 radiation scheme for the current SAM-LSM version.

191 The baseline configuration for the simulations analyzed in this paper considers a
 192 domain of $200 \times 200 \times 27$ km³. This choice was made primarily to reasonably accommo-
 193 date MCSs, given that they typically span about 100 km (Houze Jr, 2004). The hori-
 194 zontal resolution is 500 m, and the vertical resolution varies: it starts at a minimum of
 195 50 m below 1.5 km and increases to 300 m in the upper troposphere. From there, it grad-
 196 ually stretches up to 500 m at the model’s upper boundary, which reaches 27 km, result-
 197 ing in 128 vertical levels. The temporal resolution is 5 seconds, and instantaneous model

198 fields and statistics are output every 30 minutes. The control simulation uses the P3 mi-
 199 crophysics scheme. The CAM3 radiation scheme is called every 150 seconds. Surface fluxes
 200 are calculated through the LSM (see Section 3.3).

201 The large-scale forcing is based on the VARANAL dataset for the period of De-
 202 cember 2014 in the central Amazon. Winds were nudged with a 2-hour timescale through-
 203 out the simulation. The water vapor mixing ratio was nudged only during the spin-up,
 204 considered as the period from 1-5 December 2014, with a timescale of 6 hours.

205 For the purpose of model validation, we conducted additional simulations where
 206 the only modificatio was the choice of the microphysics scheme: single-moment, double-
 207 moment, or P3 schemes. We also assessed model sensitivity to resolution and domain
 208 size by performing additional simulations at 250 m resolution or with a $400 \times 400 \times 27 \text{ km}^3$
 209 domain.

210 **3.3 Land Surface Model Configuration**

211 The simplified Land Surface Model uses a minimalist set of parameters to solve the
 212 transport of heat, moisture, and radiation in the soil and vegetation and calculate the
 213 transfer of momentum between the surface and the atmosphere (Lee & Khairoutdinov,
 214 2015). It adequately idealizes the land-atmosphere interactions, which fundamentally in-
 215 fluence convection over the Amazon forest (Silva Dias et al., 2002; Betts, 2002). To de-
 216 fine the characteristics of the surface and the vegetation for the LSM, we use the MODIS
 217 data of land cover type during 2014 and mean leaf area index (LAI) for December 2014
 218 over a domain of $200 \times 200 \text{ km}^2$ centered at T3 site (see Figure S1). These are associated
 219 with the period and area of our simulations. The surface LAI was set to $4.4 \text{ m}^2 \text{ m}^{-2}$,
 220 which corresponds to the domain average in satellite observations. In addition, we choose
 221 to cover the surface uniformly with evergreen broadleaf forests, which cover 83% of the
 222 area in observations. However, based on several tests conducted to optimize the LSM
 223 parameters, we modify the default near-infrared visible albedo for vegetation from 0.20
 224 to 0.30, the root length from 150 cm to 200 cm, and the displacement height factor from
 225 0.68 to 0.65. The corresponding displacement height is $0.65 \times 20 \text{ m} = 13 \text{ m}$, where 20
 226 m is the default value of the height of the canopy. These modifications improve the agree-
 227 ment between the observations and the simulated surface radiation, latent and sensible
 228 heat fluxes.

229 The soil is simulated using 11 layers from the surface down to a depth of 400 cm.
 230 Clay and sand contents for each layer are based on in-situ measurements on Terra-firma
 231 forests (section 3.1). The initial conditions for soil temperature and soil wetness are based
 232 on GLDAS Noah, which provides information on 4 layers: 0-10, 10-40, 40-100, and 100-
 233 200 cm. The LSM soil layers close to the surface, which experience greater diurnal cy-
 234 cle variation, are interpolated using the nearest neighbor method. The deeper soil lay-
 235 ers are interpolated (and extrapolated) linearly. The initial profile of soil temperature
 236 and wetness is shown in Figure S2.

237 **3.4 Cloud Regime Days**

238 For the sensitivity experiments, we select a set of deep convective (Deep) days from
 239 the control simulation (P3 scheme, with a horizontal resolution of 500 m and a domain
 240 size of $200 \times 200 \times 27 \text{ km}^3$) and perturb the sounding or the large-scale forcing imposed.
 241 For the Deep selection, we require that the domain average of total ice presents a dis-
 242 tinct deepening during the afternoon, characteristic of the STD convective transition.
 243 The chosen Deep days are December 17th, 21st, 23rd, and 26th. For comparison, we also
 244 select a set of shallow cumulus (ShCu) days from the control run. We identify four days
 245 with negligible ice content and minimum surface precipitation: December 9th, 13th, 27th,

246 and 28th. Figure S3 shows the profile of cloud liquid and total ice for our selection of
 247 cloud regime days.

248 4 Model Validation

249 A comparison between simulated and observed CWV, precipitation rate, LE, and
 250 H is shown in Figure 2. Results are shown for simulations using the single-moment, double-
 251 moment, and P3 microphysics schemes, in addition to the higher-resolution (P3/250m)
 252 and larger domain (P3/400km) runs.

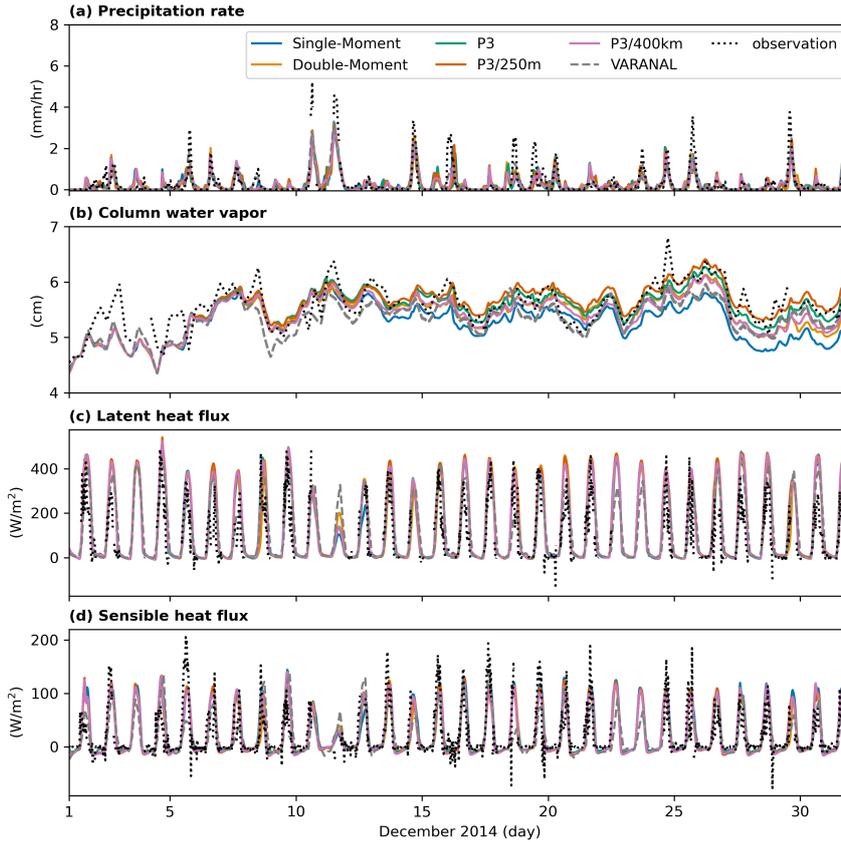


Figure 2. Comparison of modeling results (solid colors), large-scale forcing (dashed gray), and observations (dotted) for (a) Precipitation rate (mm/hr), (b) Column water vapor (cm), (c) Latent heat flux (W/m^2), and (d) Sensible heat flux (W/m^2) averaged over the model domain.

253 Simulated and observed CWV values agree very well in the first week, likely a re-
 254 sult of the nudging imposed during the spin-up period, after which some differences be-
 255 gin to appear. CWV values are generally higher for the P3 scheme and lower for the single-
 256 moment run. The P3 cases exhibit the strongest correlation (Pearson) with the obser-
 257 vations: 0.78 for P3/250m and 0.75 for both P3 and P3/400km (see the Taylor diagrams
 258 in Figure S4). Conversely, the single-moment scheme shows a weaker correlation with
 259 the observed CWV (0.56), while the double-moment scheme correlation is close to that
 260 of the P3 scheme (0.70). For the standard deviation of CWV, the model values range
 261 from 0.24 mm (P3/250m) to 0.28 mm (Single-Moment), while the observations indicate
 262 a value of 0.34 mm. Despite this difference, these statistics suggest that the model can

reproduce observed moisture content reasonably well for at least one month without resorting to any water vapor nudging.

The different simulations closely reproduce the observed surface precipitation rate, with correlations ranging from 0.76 (P3/250m) to 0.79 (P3/400km). The simulations exhibit better agreement for lower precipitation rates, while they tend to underestimate the most intense precipitation events, which are associated with MCSs. Moreover, the model precipitation did not show significant sensitivity to the microphysics, spatial resolution, or domain size. We hypothesize that the model’s underestimation of intense surface precipitation could potentially be attributed to the periodic boundary conditions. These conditions might prevent the advection of MCSs that could have developed in areas outside the domain. Nevertheless, our validation results remain satisfactory, particularly considering our primary focus is locally-driven STD convective transitions.

Observed surface fluxes are reproduced reasonably well in the model runs. LE correlations with observations vary from 0.81 (double-moment) to 0.84 (P3, P3/250m, and P3/400km), while the H correlations range from 0.78 (double-moment) to 0.80 (P3, P3/250m, and P3/400km). The model only slightly overestimates the standard deviation of the observed mean LE, with the difference between model runs and observations being less than 2 W m^{-2} . However, it should be noted that the ECOR flux measurement system provides local measurements of surface fluxes in a grassland region (T3 site, see Figure 1b), while the model provides an average for an area of $200 \times 200 \text{ km}^2$ (or $400 \times 400 \text{ km}^2$ for P3/400km), entirely covered by evergreen broadleaf forest. These differences make the qualitative agreement between model simulations and observations all the more remarkable.

To evaluate the surface radiation budget, Figure 3 compares modeled and observed surface shortwave and longwave fluxes, including both downward and upward components. There is high-frequency variability in the observations that is not present in the model, likely because its values correspond to horizontal averages over the domain, whereas observational values are taken at the T3 site. Nevertheless, the model reproduces the observations satisfactorily for downward/upward surface shortwave and upward longwave fluxes (correlation ranges 0.82-0.86, Figure S5). In the case of downward longwave fluxes, the correlation is weaker, ranging from 0.57 (single and double-moment) to 0.62 (P3, P3/250m, and P3/400km), although these values are reasonable.

Overall, our simulations with different microphysics schemes compared reasonably well with the observations considered in our validation analysis. The exception was the column water vapor, where the P3 scheme showed a stronger correlation with the observations. While neither the higher resolution (P3/250m) nor the larger domain size (P3/400km) simulations demonstrated significant improvements over the P3 case, they significantly increased computational costs. This motivated our choice of the P3 scheme with 500 m horizontal resolution and a $200 \times 200 \text{ km}^2$ domain size as the control run configuration which underlies the results presented below.

5 Sensitivity Experiments

To evaluate the role of moisture and vertical wind shear in the STD convective transition, a series of sensitivity experiments are carried out. First, we perturb the water vapor profile at low levels and the free troposphere to investigate the importance of low- and mid-level preconditioning. For vertical wind shear, we modify the structure of the low or upper-level jets to evaluate the relative importance of wind shear at different levels. The results in this section are associated with the mean composites for the four Deep or ShCu (section 3.4) simulated days.

Figure 4 shows the composite of cloud liquid (r_l), total ice (r_i), and rainwater (r_r) mixing ratios for the Deep and ShCu days averaged over the model domain. In addition

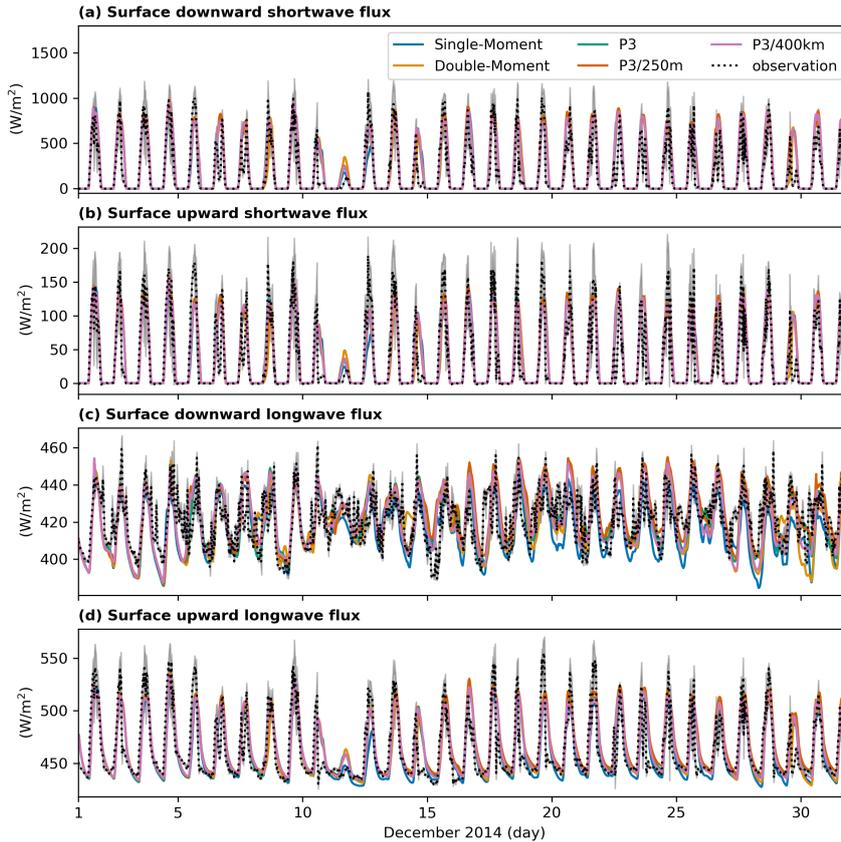


Figure 3. Similar to Figure 2, but for (a) downward shortwave, (b) upward shortwave, (c) downward longwave, and (d) upward longwave flux at the surface (W/m^2) averaged over the model domain.

313 to liquid water, we include the convective boundary layer (CBL) height (magenta dashed
 314 line), defined as the height at which the buoyancy flux reaches its first local minimum.
 315 The figure shows that the CBL height follows closely cloud base and the lifting conden-
 316 sation level, with values reaching a maximum of 1.30 km at 14:45 LST on Deep days and
 317 1.41 km at 16:15 on ShCu days. Both regimes exhibit a peak in r_l associated with shal-
 318 low convection, below 3 km between 10-14 LST. Additionally, Deep days show two peaks
 319 in r_i . The first occurs between 12-14 LST at upper levels (> 8 km), associated with deep
 320 convection driven by surface heating (Martin et al., 2016; Tang et al., 2016; Zhuang et
 321 al., 2017; Tian et al., 2021). The second peak occurs a few hours later, between 16-18
 322 LST, associated with the late afternoon STD convection transition triggered by these
 323 land-atmosphere interactions.

324 5.1 Low-Level Moisture Experiment

325 For the low-level moisture experiment, a moisture perturbation within the lower
 326 1.5 km of the domain is introduced. We introduce this perturbation by multiplying the
 327 water vapor mixing ratio by a constant factor. To ensure smoothness of the vertical pro-
 328 file, we linearly decrease the factor to 0 between altitudes of 1.25 km and 1.75 km. For
 329 each of the Deep days selected, the model is restarted from the control run at 02 LST
 330 and the perturbation was applied. Finally, the factors are selected such that the CWV
 331 for the perturbed profile at 02 LST drops by 1, 2, 3, and 4 mm.

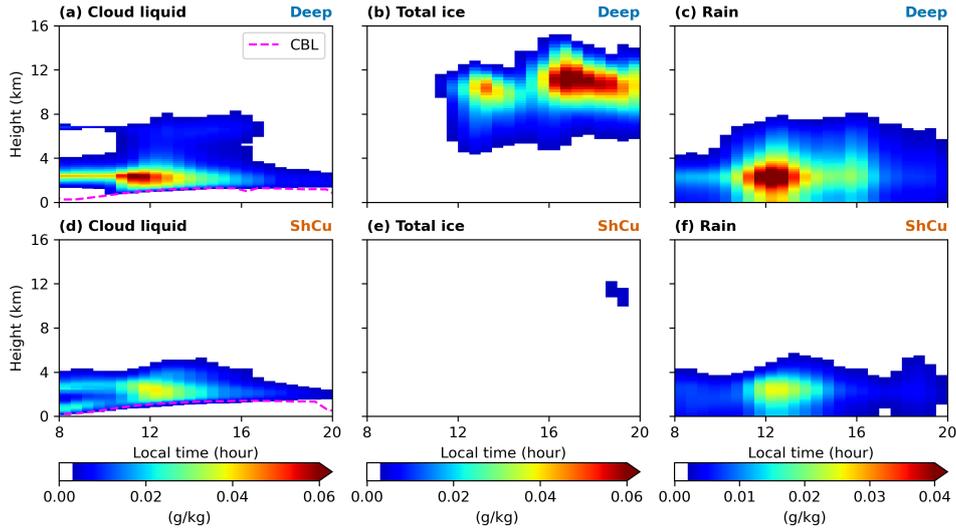


Figure 4. Composites for the Deep (top) and ShCu (bottom) regime days showing the diurnal cycle of domain-averaged (a,d): Cloud liquid water (colormap) and convective boundary layer (magenta dashed line), (b,e): Total ice, and (c,f): Rain content. The Deep selection includes the 17th, 21st, 23rd, and 26th of December 2014, while the ShCu selection comprises the days of the 9th, 13th, 27th, and 28th (section 3.4).

332 Figure 5 depicts the 30-minute average water vapor profiles at 02:15 LST, 08:15
 333 LST, and 14:15 LST, along with the daytime CWV for both moisture sensitivity exper-
 334 iments (low-level and free troposphere, with the latter described in the next section). The
 335 low-level dry perturbations (continuous lines) diminish from nighttime to afternoon due
 336 to latent heat flux and moisture tendencies. Above the perturbed region, the mixing ra-
 337 tio values are remarkably similar, indicating minimal vertical mixing. At 14 LST, be-
 338 fore the late afternoon STD convective transition, the lower tropospheric water vapor
 339 for the case where CWV drops by 3 mm (BL_{3mm}) is similar to those for the ShCu days.
 340 In terms of CWV, the experiments BL_{3mm} and BL_{4mm} demonstrate lower CWV values
 341 compared to ShCu days in the early morning. However, while the ShCu composite re-
 342 mains relatively stable throughout the diurnal cycle, the Deep composite exhibits wa-
 343 ter vapor convergence, leading to higher CWV values than the ShCu days for all exper-
 344 iments in the afternoon.

345 Figure 6 shows the magnitude (colors) and relative (contours) difference between
 346 experiments and control case for cloud liquid, total ice, and rain domain-averaged mix-
 347 ing ratios. In addition, the liquid water path, ice water path, and surface precipitation
 348 (lines) are also included. Cloud liquid water is reduced up to 75% near the cloud base,
 349 with a more extensive impact observed for drier scenarios from 10 to 12 LST. Above 3
 350 km, the amount of liquid water experiences a significant increase of up to 100% during
 351 the afternoon in drier scenarios, reflecting a greater presence of warm clouds. Addition-
 352 ally, the peak in liquid water in drier cases occurs later, shifting from 11:45 LST in the
 353 control case to 13:15 LST for experiments BL_{3mm} and BL_{4mm} .

354 Ice water content shows a significant sensitivity to low-level dry perturbations. For
 355 example, the control case exhibits an ice water path maximum of 110.1 g m^{-2} , declin-
 356 ing to 67.2 g m^{-2} , 56.2 g m^{-2} , 40.0 g m^{-2} , and reaching a minimum of 9.5 g m^{-2}
 357 with a decrease in CWV by 1, 2, 3, and 4 mm at 02 LST, respectively. The relative differences
 358 compared to the control case are 33.5%, 44.4%, 60.4%, and 90.6%, respectively, indicat-
 359 ing a non-linear decrease in the ice water path as a function of the change in CWV.

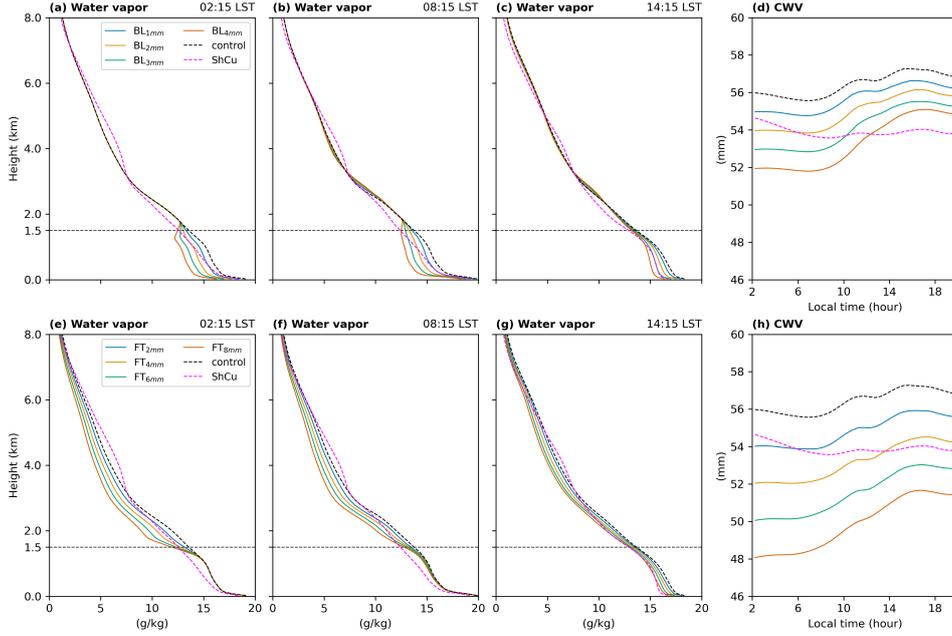


Figure 5. Moisture perturbation. Specific humidity (g/kg) profile at (a,e) 02:15 LST, (b,f) 08:15 LST, (c,g) 14:15 LST, and time series of (d,h) column water vapor (mm) for low-level (upper panels, solid colors) and free troposphere (lower panels, solid colors) moisture experiments. The Deep composite (control) is the dashed black line, and the ShCu composite is the dashed magenta line.

360 For the rain content, a decrease in drier scenarios is observed from early morning
 361 to early afternoon. While the reduction remains insignificant when reducing CWV by
 362 1 mm, contours of 50% and 75% emerge in the drier cases. In terms of precipitation rate,
 363 the control case shows a peak of 0.86 mm hr^{-1} , declining to 0.78, 0.68, 0.55, and 0.42
 364 mm hr^{-1} with a decrease in CWV by 1, 2, 3, and 4 mm at 02 LST, respectively. The
 365 decrease in the maximum precipitation rate resulting from changes in CWV can be ef-
 366 fectively modeled by a quadratic function (Figure S6). Note that despite the observed
 367 sensitivity of ice content to low-level dry perturbations, where ice content becomes neg-
 368 ligible in the driest scenario, the model still produces significant amounts of warm pre-
 369 cipitation for all experiments.

370 5.2 Free Troposphere Moisture Experiment

371 The experiments conducted in the free troposphere are analogous to the low-level
 372 moisture experiments, differing only in that the perturbation is applied above 1.5 km.
 373 We select multiplicative factors such that the CWV for the perturbed profile at 02 LST
 374 drops by 2, 4, 6, and 8 mm, respectively. These changes in CWV correspond to double
 375 what is applied in the previous section, and our choice is motivated by the weaker sen-
 376 sitivity of free troposphere humidity to convection, as we will show below.

377 The dry perturbations in the free troposphere diminish throughout the day, simi-
 378 larly to the low-level perturbations (Figure 5). The drier scenarios also exhibit slightly
 379 drier conditions at lower levels by early afternoon compared to the control run. The free
 380 troposphere experiments start with a lower CWV value than that of shallow days. How-
 381 ever, the case $\text{FT}_{2\text{mm}}$ already presents higher CWV values than the shallow composite
 382 in the early morning, while $\text{FT}_{4\text{mm}}$ shows higher values than shallow days only in the

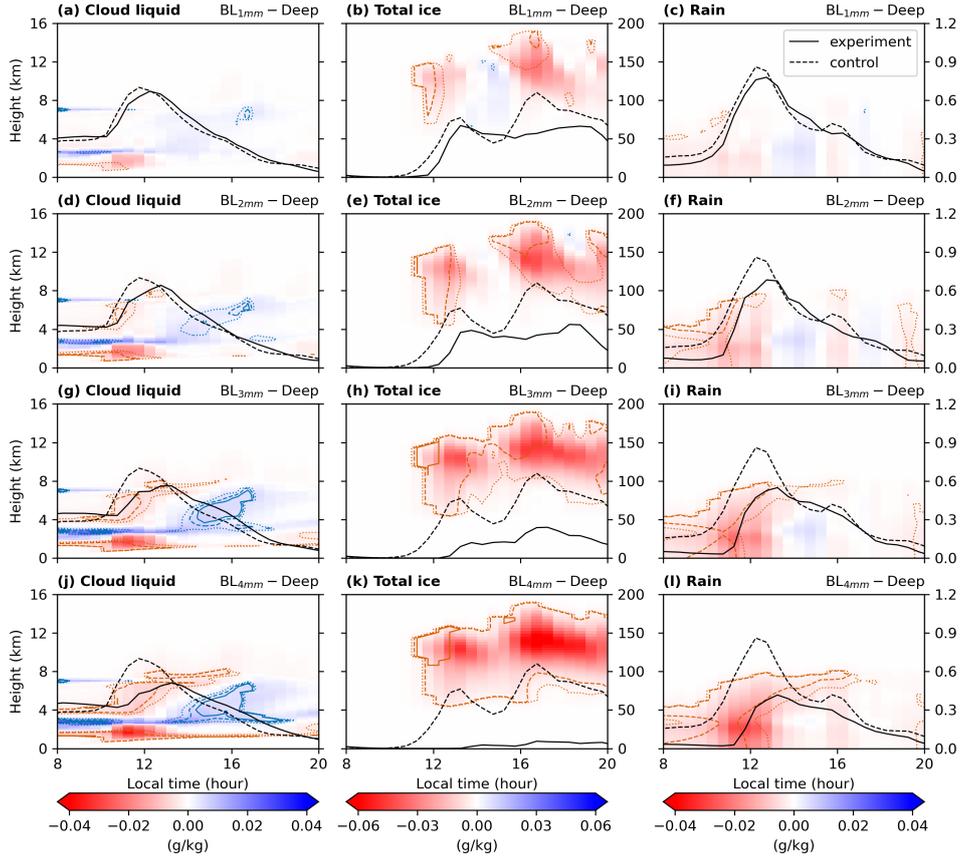


Figure 6. Low-level moisture experiment. The composites show the diurnal cycle of domain-averaged anomalies for (a,d,g,j) Cloud liquid water (g/kg), (b,e,h,k) Total ice (g/kg), and (c,f,i,l) Rain (g/kg) mixing ratios. Each row corresponds to a different decrease in CWV, (a-c) 1 mm, (d-f) 2 mm, (g-i) 3 mm, and (j-l) 4 mm. The colors indicate the absolute difference between each experiment and the control, while the contours show relative differences of 50% (dotted), 75% (dashed), and 100% (solid). Liquid and ice water paths (right axis, g m^{-2}) are presented along with cloud liquid and total ice, respectively, and surface precipitation rate (right axis, mm hr^{-1}) is shown alongside rainwater. The solid line represents the experiment, while the dashed line represents the control runs conducted during the Deep days.

383 afternoon (around 14 LST). For the drier cases, $\text{FT}_{6\text{mm}}$ and $\text{FT}_{8\text{mm}}$, their CWV remains
 384 lower than that of the shallow days throughout the simulation.

385 Figure 7 presents the results for the free-troposphere moisture experiments and control
 386 runs during the Deep days. While applying a dry perturbation above 1.5 km leads
 387 to a reduction in cloud liquid water and liquid water path throughout the troposphere,
 388 this impact is relatively minor compared to what is observed in the low-level experiments.

389 Cloud ice water exhibits a greater sensitivity to the free troposphere perturbations.
 390 However, the impact is still relatively minor compared to the perturbations at low levels.
 391 For instance, when the perturbation in the free troposphere and at low levels leads
 392 to a 2 mm drop in CWV, the maximum ice water path is 99.1 g m^{-2} and 56.2 g m^{-2} ,
 393 respectively. Similarly, with a 4 mm drop in CWV, the maximum ice water path is 60.9
 394 g m^{-2} and 9.5 g m^{-2} for the perturbations in the free troposphere and at low levels,
 395 respectively. Moreover, the driest free troposphere case ($\text{FT}_{8\text{mm}}$) still exhibits significant

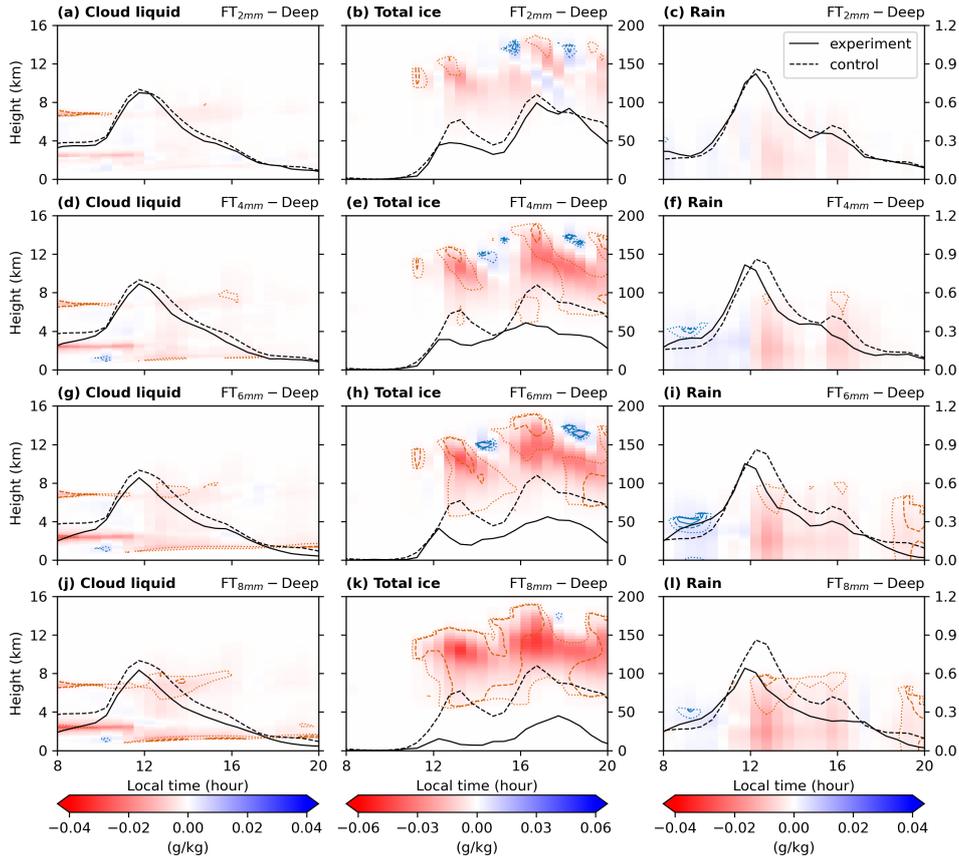


Figure 7. Same as Figure 6, but for the free troposphere moisture experiment instead, where the decreases in CWV were 2, 4, 6, and 8 mm.

396 ice water path values, with a maximum of 45.0 g m^{-2} at 17:45 LST, being even greater
 397 than the value observed for the experiment where CWV is reduced by 3 mm at low levels
 398 ($BL_{3\text{mm}}$, 40.0 g m^{-2}).

399 Rain content also shows a reasonable sensitivity to the free troposphere perturbations,
 400 although it is minor compared to low-level perturbations. For example, when the
 401 perturbation in the free troposphere and at low levels leads to a 2 mm drop in CWV,
 402 the peak of precipitation is 0.82 mm hr^{-1} and 0.68 mm hr^{-1} , respectively. Similarly, with
 403 a 4 mm drop in CWV, the peak for the free troposphere experiment remains the same
 404 (0.82 mm hr^{-1}) while the low-level perturbation shows 0.42 mm hr^{-1} . Finally, the driest
 405 free troposphere case ($FT_{8\text{mm}}$) shows a maximum of 0.64 mm hr^{-1} , which better relates
 406 with experiment $BL_{2\text{mm}}$ (0.68 mm hr^{-1}).

407 5.3 Wind Jet Experiment

408 To assess the impact of wind shear on the deepening of convective clouds in the
 409 Amazon, we perform sensitivity experiments in which the low- or high-level jets are mod-
 410 ified. Each jet's intensity, amplitude, and altitude are changed separately. The modified
 411 wind profiles are used to force the model, with nudging applied with a timescale of 2 hours
 412 throughout the simulation.

413 In order to have better control of the shape of the wind profiles and easily gener-
 414 ate sensitivity tests, we approximate the imposed wind using an analytical formula. More

415 specifically, considering the shape of the wind speed, we model this quantity as the su-
 416 perposition of two Gaussian functions, each representing a jet. The average wind speed
 417 is fitted to this function, and the fitting parameters are given in Table S1. The wind di-
 418 rection was fitted to a piece-wise linear function. (Figure 8b). The wind direction is con-
 419 stant in the bottom ~ 2 km. It veers clockwise at a constant rate of about 14° km^{-1} from
 420 2 to 15 km, and counterclockwise at a rate of $-28^\circ \text{ km}^{-1}$ from 15 to 20 km, and it is con-
 421 stant above 20 km (not shown). The wind speed and direction imposed for control runs
 422 are shown in Figure 8a-b (solid black line).

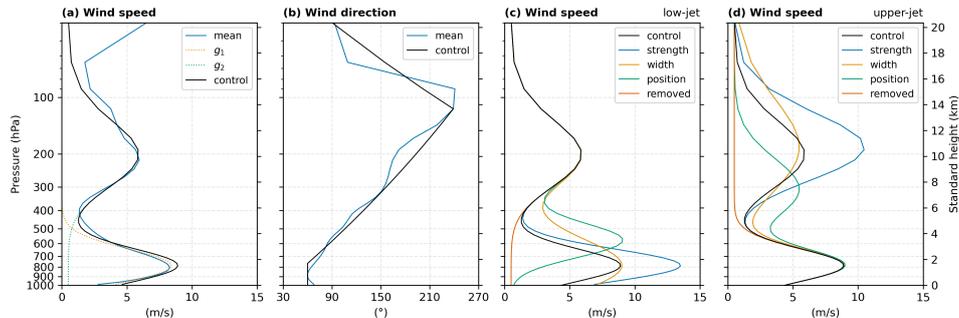


Figure 8. Wind profiles for the jet experiments control run showing: (a) large-scale wind speed and (b) wind direction, as measured (blue) and as idealized (black). Sensitivity experiments perturbed the wind speed profile of either the (c) low-level or (d) upper-level wind jets by intensifying (blue), widening (orange), shifting (green), or removing the jet (red).

423 For each jet, we conduct 4 experiments in which we (1) increase the jet strength,
 424 (2) increase the width, (3) shift the peak position, and (4) remove the jet entirely while
 425 keeping the wind direction constant in all cases. The modified wind profiles are shown
 426 in Figures 8c-d. For each experiment, the model is restarted from the control run at 02
 427 LST of each one of the four Deep days selected, and the modified wind profiles are ap-
 428 plied to force the model with a 2-h nudging timescale.

429 Figure 9 presents the results for the low-level jet experiments. While the jet po-
 430 sition significantly affects cloud liquid water, the jet width has a negligible impact. The
 431 higher position possibly enhances the low-level updrafts, leading to an increase of r_l above
 432 2 km, particularly in the late afternoon (around 16 LST), where r_l can increase by as
 433 much as 100%. Removing the low-level jet also exhibits a similar impact on cloud liq-
 434 uid water, although the anomalies are smaller than those associated with the higher jet.
 435 The stronger jet only slightly impacts the cloud liquid water. Although the jet influences
 436 the cloud water profile, the integrated liquid water path is similar in all experiments and
 437 does not significantly differ from the control.

438 The cloud ice content is more significantly affected by the low-level jet. The con-
 439 trol experiment exhibits peaks in ice water path at 12:45 LST and 16:15 LST. The stronger
 440 jet simulation shows positive anomalies before 12 LST and after 16 LST, resulting in an
 441 increase in r_i of up to 100%, with negative anomalies observed in between. For the wider
 442 jet, a positive anomaly dominates throughout the diurnal cycle, especially between 12-
 443 17 LST, when r_i increases up to 100% above 12 km. However, there is only a modest
 444 increment in ice water path. The higher jet shows a decrease ranging from 50-75% in r_i
 445 around 12 LST and an increase afterwards. There is a delay of the convective activity,
 446 with a suppression at 12 LST and an increase around 16 LST. When the low-level jet
 447 is removed, the ice water path is substantially reduced, both at 12 and 16 LST (up to
 448 75%), with the maximum being reached only around 17:45.

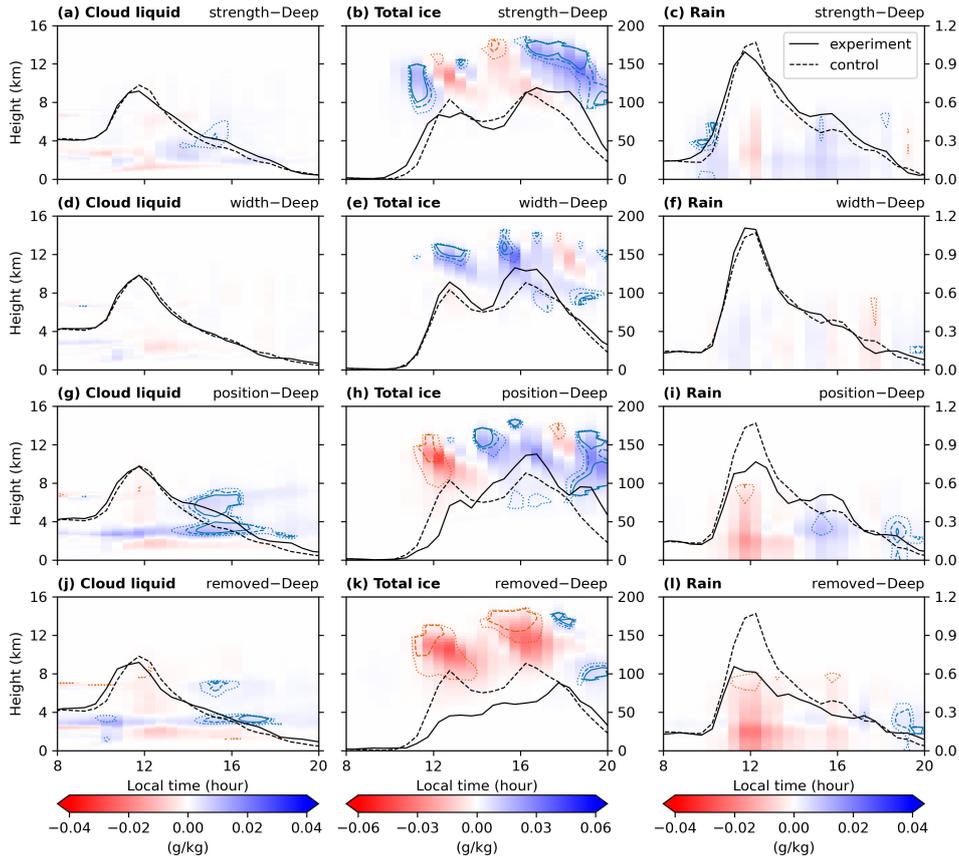


Figure 9. Low-level jet experiment. The composites show the diurnal cycle of domain-averaged anomalies for (a,d,g,j) Cloud liquid water (g/kg), (b,e,h,k) Total ice (g/kg), and (c,f,i,l) Rain (g/kg) mixing ratios. Each row corresponds to a different low-level jet perturbation, (a-c) strength, (d-f) width, (g-i) position, and (j-l) removed entirely. The colors indicate the absolute difference between each experiment and the control, while the contours show relative differences of 50% (dotted), 75% (dashed), and 100% (solid). Liquid and ice water paths (right axis, g/m^2) are presented along with cloud liquid and total ice, respectively, and surface precipitation rate (right axis, mm hr^{-1}) is shown alongside rainwater. The solid line represents the experiment, while the dashed line represents the control runs conducted during the Deep days.

449 Rain content follows the changes in total ice mixing ratio, with negative (positive)
 450 anomalies where ice decreases (increases). However, changes in r_r are less significant, with
 451 smaller areas showing changes greater than 50%. While the surface precipitation shows
 452 minimal impact in the stronger and wider jet experiments, it exhibits notable differences
 453 for the other two cases. For the higher jet, precipitation decreases around noon and in-
 454 creases in the late afternoon, following the changes in ice water path. In the absence of
 455 a jet, precipitation is particularly reduced between 12-15 LST, with the noon peak be-
 456 ing roughly 40% lower than that observed in the control runs.

457 Figure 10 presents the results for the upper-level jet experiments. There is a strik-
 458 ing contrast with the low-level jet results. The upper-level jet affects convection only in
 459 the upper troposphere, with negligible impacts on liquid water. Overall, the experiments
 460 show an alternating increasing and decreasing ice content pattern. This is related to a
 461 delay in convection, which can be more easily noticed on the ice water path curves. Both

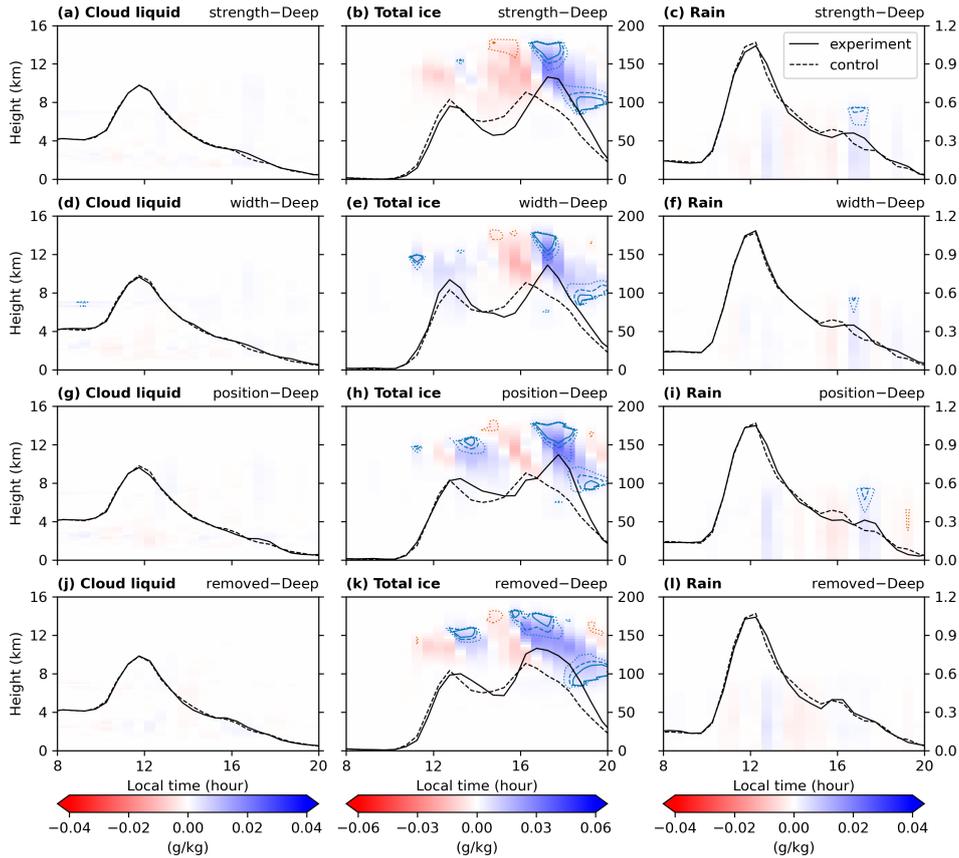


Figure 10. Same as Figure 9, but for the upper-level jet experiment instead.

462 the noon and afternoon peaks are displaced to later times, and the afternoon peak also
 463 gets more intense. The exception is the wider jet experiment, where the noon peak slightly
 464 increases and is not delayed. The changes are more significant after 17 LST when all ex-
 465 periments exhibit an increase of up to 100% in ice water. In the case of the removed jet,
 466 the afternoon peak occurs 30 minutes later (16:45 LST). For the stronger and wider jet
 467 experiments, the peak is delayed by 1 hour (17:15 LST), and for the jet with a relatively
 468 lower position, the peak occurs 1.5 hours later (17:45 LST). Nonetheless, there is no sig-
 469 nificant change in the rain content and surface precipitation.

470 6 Discussion

471 While several studies in the literature have employed CRMs to simulate convective
 472 properties in both continental (M. F. Khairoutdinov & Randall, 2003; M. Khairout-
 473 dinov & Randall, 2006; Henderson & Pincus, 2009; Cecchini et al., 2022) and ocean (Blossey
 474 et al., 2007; M. F. Khairoutdinov et al., 2009; Liu et al., 2015; Blossey et al., 2021) re-
 475 gions, we conduct simulations specifically for the central Amazon by coupling a CRM
 476 with a LSM, thus explicitly representing biosphere-atmosphere feedbacks, which plays
 477 a crucial role on convection over the Amazon tropical forest (Silva Dias et al., 2002). Vilà-
 478 Guerau de Arellano et al. (2020) utilized a large-eddy model combined with a different
 479 LSM to investigate the diurnal cycle of energy, moisture, and carbon dioxide from clear to
 480 cloudy conditions. Nevertheless, their study relied only upon a specific day during the
 481 Amazon dry season and did not address deep convection. Although the recent work by
 482 Gonçalves et al. (2022) also utilized a CRM coupled with an LSM to simulate convective

483 tion in the central Amazon, they did not evaluate the evolution of moisture content and
484 surface latent and sensible heat fluxes. This restricts the validation of their simulations
485 in reproducing these crucial convective properties, which we have addressed in our study.
486 Specifically, we evaluate the model performance using single-moment, double-moment,
487 and P3 microphysics schemes. Only the CWV exhibits a noticeable sensitivity to the mi-
488 crophysics, where the P3 scheme shows the strongest correlation (Pearson, 0.8) with ob-
489 servations. Varying the horizontal resolution (from 500 to 250 m) and domain size (from
490 $200 \times 200 \times 27$ to $400 \times 400 \times 27$ km³) reveals minimal sensitivity to the model results. For
491 the surface fluxes, the P3 microphysics scheme indicates a correlation of about 0.8 for
492 both latent heat and sensible heat fluxes. Overall, our results demonstrate that our sim-
493 ulations satisfactorily reproduce the convective properties in the central Amazon.

494 Our sensitivity experiments indicate that the humidity in the early morning at lower
495 levels plays a crucial role in the late afternoon STD convective transition in the Ama-
496 zon. By reducing water vapor in the lowest 1.5 km, the diurnal peak of the ice water path
497 substantially drops with changes in CWV, whereas a change of 2 mm causes a reduc-
498 tion of ice in the range of 50-75%, and a change of 4 mm results in negligible ice dur-
499 ing the simulation. Moreover, the maximum precipitation rate demonstrates an approx-
500 imately quadratic decrease with variations in the low-level CWV. Schiro et al. (2016) ex-
501 amined the relationship between precipitation and CWV by calculating the condition-
502 ally averaged precipitation to CWV using local observations of both variables at corre-
503 sponding times. They observed that the probability and intensity of precipitation can
504 be roughly characterized by an exponential function of CWV magnitude. It is essential
505 to highlight that our approach differs from that of Schiro et al. (2016). We used the domain-
506 average for maximum afternoon precipitation, while CWV precedes the diurnal cycle.
507 This procedural difference somewhat justifies why we observed a quadratic relationship
508 instead of an exponential one, as there is no contradiction between these results. Fur-
509 thermore, our findings indicated that achieving a comparable reduction in ice water path
510 or precipitation, as observed in the low-level experiments, requires reducing the column
511 water vapor by a factor of two or more in the free troposphere.

512 While M. Khairoutdinov and Randall (2006) conducted experiments that differed
513 from those designed in this study and were limited to a single idealized case during TRMM-
514 LBA on February 23, 1999, they similarly indicated that free troposphere precondition-
515 ing plays a minor role in convection in the Amazon. Based on GoAmazon2014/5 obser-
516 vations, Ghate and Kollias (2016) noted that locally-driven precipitating days during the
517 dry season show an early morning water vapor excess above the boundary layer while
518 Zhuang et al. (2017); Tian et al. (2021) found that deep convective days exhibit a moister
519 environment extending from the surface to higher levels, regardless of the season. Schiro
520 and Neelin (2019) showed that the onset and probability of the STD transition are closely
521 linked to both lower-free-tropospheric moisture (700–900 hPa) and boundary layer mois-
522 ture. Conversely, MSC likelihood rises with higher lower-free-tropospheric humidity, while
523 the relationship with boundary layer moisture is less distinct. The relative importance
524 of moisture to convection can also vary based on the regions being studied. Focusing on
525 the Tropical Western Pacific region on Nauru Island, Holloway and Neelin (2009) found
526 a strong correlation between observed precipitation and moisture variability in the free
527 troposphere, with limited variability in the boundary layer. Additionally, Bretherton et
528 al. (2004) also highlighted the importance of free-tropospheric humidity to convection
529 over the Tropical Oceans.

530 Vertical wind shear primarily impacts the peak timing of ice water in our simula-
531 tions. Furthermore, our findings indicate that convection is enhanced during the after-
532 noon when the low-level wind is idealized using a jet of larger width, moderate strength,
533 and with a relatively higher peak position from around 2 to 4 km. Conversely, the upper-
534 level wind has a minor influence on convective intensity. M. Khairoutdinov and Randall
535 (2006) designed experiments employing an idealized large-scale wind forcing and a free

536 wind shear environment. Similar to our results, the STD transition was not prevented
 537 by removing vertical wind shear. Cecchini et al. (2022) also conducted numerical exper-
 538 iments to quantify the impact of vertical wind shear in the central Amazon, specifically
 539 targeting shallow cumulus convection during a typical day in the dry season. By intro-
 540 ducing incremental changes in the large-scale wind speed across the entire vertical do-
 541 main, the authors observed a weakening of convective intensity, suggesting that verti-
 542 cal wind shear prevents the STD convective transition. Here, we have identified that the
 543 vertical level of wind shear significantly influences its impact on convection. Moreover,
 544 while stronger vertical wind shear suppresses the initial phase of convection, the STD
 545 convective transition still occurs, albeit with a delay ranging from a few minutes to an
 546 hour.

547 In contrast to prior observational studies, our modeling results offer quantitative
 548 insights into the role of vertical wind shear in Amazonian convection. For example, Zhuang
 549 et al. (2017) observed that ShCu days are linked to stronger mid-level wind shear dur-
 550 ing the wet season. We observe that when the upper-level jet is shifted from around 12
 551 to 8 km, which is related to mid-level wind shear, the trigger for the STD transition is
 552 only delayed by about 1.5 hours. While Chakraborty et al. (2018) indicated that more
 553 intense low-level shear is associated with shallow convection during the transition sea-
 554 son, we observe that stronger low-level shear has little influence on cloud liquid water
 555 and might only provoke a delay in the late afternoon STD transition.

556 Thus, our findings suggest that a moderate shear environment might more efficiently
 557 separate downdrafts and updrafts within the cloud while concurrently organizing the con-
 558 vergence of low-level water vapor within the cloud layers. A wider jet leads to a smoother
 559 and more gradual shift in wind shear, extending from the surface to higher altitudes, thereby
 560 also organizing the water vapor convergence from below the cloud base to higher levels.
 561 Given the dependence of water vapor convergence on low-level humidity, convection demon-
 562 strates heightened sensitivity to boundary layer humidity. Meanwhile, the upper-level
 563 jet primarily impacts the extensive cloud anvil, exerting a relatively minor influence on
 564 ice content above 8 km. Finally, these results collectively suggest that dynamic factors
 565 may exert a more pronounced influence on convection in the Amazon.

566 7 Conclusions

567 While numerous observational studies have explored the environmental controls on
 568 convection in the Amazon (Itterly et al., 2016; Ghate & Kollias, 2016; Zhuang et al., 2017;
 569 Schiro et al., 2016, 2018; Chakraborty et al., 2018; Tian et al., 2021; Giangrande et al.,
 570 2023), we have specifically addressed this problem through high-resolution idealized sim-
 571 ulations. We employ the System for Atmospheric Modeling (SAM) model coupled with
 572 a LSM to perform simulations for the Amazon region in December 2014. The model is
 573 forced with the large-scale fields from the variational analysis, and the observations from
 574 the GoAmazon2014/5 experiment are used to validate the model results. The LSM de-
 575 fault input parameters are modified according to in-situ and satellite observations over
 576 the Amazon region, and fine-tuning tests focused on improving the model agreement with
 577 the observations. The simulations consistently reproduce the observations for precipi-
 578 tation, column water vapor (CWV), surface latent and sensible heat fluxes, and surface
 579 radiation fluxes. Sensitivity tests demonstrate that simulations conducted using a single-
 580 moment microphysics scheme drifted towards a drier state, while simulations with the
 581 P3 microphysics scheme more closely reproduce the observed water budget. For a more
 582 detailed validation of the LSM, having more comprehensive observations of the soil prop-
 583 erties (e.g., temperature and wetness down to 4 m) would be necessary.

584 In light of recent observational studies addressing the shallow-to-deep (STD) con-
 585 vective transition (Ghate & Kollias, 2016; Zhuang et al., 2017; Tian et al., 2021), our study
 586 has the advantage of conducting idealized sensitivity experiments in which only one en-

587 vironmental control—moisture or vertical wind shear at low or high levels—is modified
 588 at a time. This approach efficiently isolates their influence in controlling convection. The
 589 pre-convective humidity at low levels had the greatest impact on convection. The diurnal
 590 peak in the ice water path robustly decays with changes in CWV within the lower
 591 1.5 km. To have a comparable impact on the diurnal cycle of convection, it is necessary
 592 to reduce free tropospheric CWV by approximately twice the amount in the lower lev-
 593 els. Vertical wind shear mainly affects the ice water peak timing. A wider low-level jet
 594 of moderate strength possibly facilitates the STD convective transition by organizing low-
 595 level water vapor convergence and potentially separating downdrafts and updrafts within
 596 the cloud. The upper-level wind shear has a minor influence over convection in the Ama-
 597 zon.

598 While our results provide quantitative information on the role of moisture and wind
 599 shear in convection, we suggest that sensitivity experiments be conducted using differ-
 600 ent cloud-resolving models. For instance, SAM uses periodic lateral conditions, artifi-
 601 cially impacting the numerical results. Using multiple models can aid in evaluating the
 602 robustness of the findings and identifying potential model biases. Although our sensi-
 603 tivity experiments identified that the maximum afternoon precipitation rate decreases
 604 roughly quadratically with changes in pre-convective CWV, particularly in the low-level
 605 experiment, it is important to note that this relationship was derived from only four val-
 606 ues of moisture perturbation. This limitation restricts the significance of the findings.
 607 The robustness of this conclusion should be further addressed, with a particular focus
 608 on understanding the associated mechanism for this relationship. Finally, we also rec-
 609 ommend that future studies conduct specific experiments to investigate the role of wa-
 610 ter vapor convergence and the effects of large-scale wind direction on deep convection.

611 8 Data Availability

612 The GoAmazon2014/5 observations are publicly available at [https://www.arm.gov/
 613 research/campaigns/amf2014goamazon](https://www.arm.gov/research/campaigns/amf2014goamazon). The large-scale forcing data based on the vari-
 614 ational analysis for the GoAmazon2014/5 experiment is available at the ARM Archive:
 615 http://iop.archive.arm.gov/arm-iop/0eval-data/xie/scm-forcing/iop_at_mao/.
 616 The Moderate Resolution Imaging Spectroradiometer (MODIS) data for land cover and
 617 leaf area index can be downloaded through the Application for Extracting and Explor-
 618 ing Analysis Ready Samples (*AppEEARS*, [https://appeears.earthdatacloud.nasa
 619 .gov/](https://appeears.earthdatacloud.nasa.gov/)). The Global Land Data Assimilation System (GLDAS) data for soil tempera-
 620 ture and soil wetness are available at [https://disc.gsfc.nasa.gov/datasets/GLDAS
 621 _NOAH025_3H.2.1/summary](https://disc.gsfc.nasa.gov/datasets/GLDAS_NOAH025_3H.2.1/summary). For further assistance concerning the model input files and
 622 the necessary modifications in the SAM's source code to perform the moisture pertur-
 623 bations as described in section 5.1 and section 5.2, please refer to the author.

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