

1 **Comparing gravity waves in a kilometre-scale run of**
2 **the IFS to AIRS satellite observations and ERA5**

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8 **Key Points:**

- 9 • A kilometre-scale IFS run was resampled as AIRS to compare gravity wave prop-
10 erties using two different methods
- 11 • Gravity waves can be seen in the resampled IFS run and AIRS at similar times
12 and locations
- 13 • Mean amplitudes in the resampled IFS run were found to be significantly lower
14 than in the observations by a factor of ~ 2.4

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15 **Abstract**

16 Atmospheric gravity waves (GWs) impact the circulation and variability of the at-
 17 mosphere. Sub-grid scale GWs, which are too small to be resolved, are parameterized
 18 in weather and climate models. However, some models are now available at resolutions
 19 at which these waves must be resolved and it is important to test whether these mod-
 20 els do this correctly. In this study, a GW resolving run of the ECMWF (European Cen-
 21 tre for Medium-Range Weather Forecasts) IFS (Integrated Forecasting System), run with
 22 a 1.4 km average grid spacing (TCO7999 resolution), was compared to observations from
 23 the Atmospheric Infrared Sounder (AIRS) instrument, on NASA’s Aqua satellite, to test
 24 how well the model resolves these waves. In this analysis, nighttime data were used from
 25 the first 10 days of November 2018 over Asia and surrounding regions. The IFS run is
 26 resampled with AIRS’s observational filter using two different methods for comparison.
 27 The ECMWF ERA5 reanalysis is also resampled as AIRS, to allow for comparison of
 28 how the high resolution IFS run resolves GWs compared to a lower resolution model that
 29 uses GW drag parametrizations. Wave properties are found in AIRS and the resampled
 30 models using a multi-dimensional S-Transform method. Orographic GWs can be seen
 31 in similar locations at similar times in all three data sets. However, wave amplitudes and
 32 momentum fluxes in the resampled IFS run were found to be significantly lower than in
 33 the observations. This could be a result of horizontal and vertical wavelengths in the IFS
 34 run being underestimated.

35 **Plain Language Summary**

36 Small-scale atmospheric waves known as gravity waves (GWs) transport energy and
 37 momentum and affect the dynamics of the atmosphere. At the resolution of the ECMWF
 38 IFS (Integrated Forecasting System) simulation, with an average grid spacing of 1.4 km
 39 (TCO7999 resolution), these waves need to be resolved. GWs in this IFS run were com-
 40 pared to those in observations from the AIRS (Atmospheric Infrared Sounder) instru-
 41 ment on NASA’s Aqua satellite, for nighttime data over Asia and surrounding regions,
 42 during the first 10 days of November 2018, to test how well these waves are resolved in
 43 the model. The high resolution IFS run is resampled as AIRS to remove GWs outside
 44 of the wavelength ranges that can be seen in the observations, allowing the data sets to
 45 be compared. GWs with orographic sources (waves formed by wind flowing over topog-
 46 raphy), can be seen at similar times and in similar locations in the IFS run and obser-
 47 vations, but wave amplitudes in the resampled IFS run were found to be significantly
 48 lower.

49 **1 Introduction**

50 Atmospheric gravity waves (GWs) are small-scale waves which transport energy
 51 and momentum throughout the atmosphere (M. J. Alexander et al., 2010; Fritts & Alexan-
 52 der, 2003). These waves have both direct and indirect effects on the atmosphere: to take
 53 just a few examples, GWs act as a major cause of clear-air turbulence affecting aircraft
 54 (Lane et al., 2009), contribute to ozone depletion in the polar stratosphere (Carslaw et
 55 al., 1998), affect the formation of sudden stratospheric warmings in winter by precon-
 56 ditioning the polar vortex (Albers & Birner, 2014) and affect the timing of the polar vor-
 57 tex breakdown in spring (Polichtchouk et al., 2018). GWs in the stratosphere have also
 58 been shown to impact the Brewer-Dobson Circulation (e.g. Sato & Hirano, 2019). Sources
 59 of GWs include orographic sources (wind flowing over topography) and non-orographic
 60 sources, such as convection and wind shear (Fritts & Alexander, 2003; M. J. Alexander
 61 et al., 2010).

62 Despite their importance for achieving realistic atmospheric circulations, GWs and
 63 their impacts remain notoriously difficult to represent in numerical models. One reason

64 for this is because large portions of the GW spectrum occur at scales below the grid size
 65 of the model, and are therefore unresolved. Instead, the acceleration (or deceleration)
 66 of the background flow at different altitudes due to GW propagation and breaking is rep-
 67 resented by parameterizations, which can be tuned to correct for the unknown momen-
 68 tum forcing due to GWs not resolved by the model. However, these parameterizations
 69 are poorly constrained by observations and contain simplifying assumptions that can lead
 70 to major circulation biases (Butchart et al., 2011; Harvey et al., 2019). Due to compu-
 71 tational constraints, this reliance on GW parameterizations is still widespread in the vast
 72 majority of operational models used for numerical weather prediction (NWP), atmospheric
 73 research and long-term dynamical climate simulations (M. J. Alexander et al., 2010; Plougonven
 74 et al., 2020).

75 In recent decades, ever increasing computational power has allowed models to be
 76 developed with sufficient spatial resolution to resolve ever larger portions of the GW spec-
 77 trum. In some of these specialist offline configurations, the resolution is sufficiently high
 78 that the effects of resolved waves alone are sufficient to achieve realistic circulations in
 79 the middle atmosphere, and GW parameterizations are no longer required (Sato et al.,
 80 2012; Vosper, 2015; Watanabe & Miyahara, 2009; Lund et al., 2020; Wedi et al., 2020).
 81 While these simulations are still prohibitively expensive for operational use, it is likely
 82 that this trend will continue and models will be able to resolve an increasingly large por-
 83 tion of the GW spectrum. This then raises a question: how realistic are the resolved waves
 84 in these high resolution simulations compared to observations?

85 Here this question is investigated for one such model: a high-resolution “kilometre-
 86 scale” configuration of the Integrated Forecasting System (IFS) model developed by the
 87 European Centre for Medium Range Weather Forecasts (ECMWF), as described by Wedi
 88 et al. (2020). This configuration was run at TCo7999 resolution (Wedi et al., 2020), which
 89 is equivalent to an average horizontal grid spacing of around 1.4 km globally, and no GW
 90 parameterizations were used. In this study, the amplitudes, wavelengths and momen-
 91 tum fluxes of resolved GWs in the model stratosphere are compared to 3-D satellite ob-
 92 servations from the Atmospheric Infrared Sounder (AIRS) instrument using the retrieval
 93 of Hoffmann and Alexander (2009). For further comparison, we also investigate resolved
 94 GWs in the ERA5 reanalysis, also produced by ECMWF, which is generated using a 9
 95 km version of the IFS, to understand the impact of the increased resolution of the km-
 96 scale IFS on resolved GW properties compared to operational configurations.

97 In Kruse et al. (2022), four numerical weather prediction models, including the IFS
 98 run with an average grid spacing of ~ 9 km, were compared to AIRS data, which showed
 99 that the models reproduced mountain waves in the observations well, near to the Drake
 100 Passage, but wave amplitudes were lower than those observed. GW momentum fluxes
 101 in the ECMWF operational analyses, produced using the IFS and 4D variational data
 102 assimilation, at a resolution of 0.125° in longitude and latitude (approximately 16 km)
 103 with 91 model levels, were found to be a factor of 5 lower than in Concordiasi balloon
 104 observations (Jewtoukoff et al., 2015). The ECMWF operational analyses was also found
 105 to have lower wave amplitudes compared to AIRS observations by a factor of 2–3 (Hoffmann
 106 et al., 2017), using data from 2003 to 2012. In this time period, the effective horizon-
 107 tal resolution and number of model levels was increased from 39 to 16 km and 60 to 91,
 108 respectively. Okui et al. (2023), found generally good agreement between the amplitudes
 109 and momentum fluxes of GWs in the Japanese Atmospheric GCM for Upper Atmosphere
 110 Research (JAGUAR) (Watanabe & Miyahara, 2009), a high vertical resolution model
 111 which uses no GW drag parameterizations, and AIRS observations.

112 However, making comparisons between observed and simulated GWs is not straight-
 113 forward. This is because no instrument (or model) can observe the full GW spectrum.
 114 The sampling and resolution characteristics of a particular observing instrument (such
 115 as AIRS as used here) limit the range of observable GW horizontal and vertical wave-
 116 lengths, a phenomenon known as the “observational filter” of the instrument (Preusse

117 et al., 2002; M. J. Alexander & Barnet, 2007). Likewise, the spatial resolution of a model
 118 limits its ability to simulate all GW wavelengths. Therefore, to make a fair comparison
 119 between observations of GWs and resolved GWs in a model, we must first sample the
 120 model as if it were observed by the instrument by applying the instrument’s sampling
 121 pattern and horizontal and vertical resolutions to the model output fields (Wright & Hind-
 122 ley, 2018; Hindley et al., 2021). This model-sampled-as-observations dataset can then
 123 be analysed in exactly the same way as the observations and a fair comparison between
 124 the measured GW properties can be made. The approach taken to perform this sam-
 125 pling method however can vary between studies, so here we investigate two different sam-
 126 pling methods to create this dataset: one using a simplified approach described by Hindley
 127 et al. (2021) and the second using the more rigorous, but more computationally-expensive,
 128 approach of Wright and Hindley (2018).

129 The selected km-scale configuration of the IFS was run globally for the period of
 130 November 2018. During this time, significant stratospheric GW activity was observed
 131 in the model, AIRS observations and the ERA5 reanalysis over continental Asia and the
 132 surrounding regions, so we select this region over which to perform our comparison (see
 133 Fig. 1). The region is likely to contain numerous sources of orographic GW activity gen-
 134 erated by surface flow over mountain ranges, such as the Abakanski Khrebet Mountain
 135 range, the Ural mountains, the Pamir mountains and other hotspots as observed by Hoffmann
 136 et al. (2013) and Hindley et al. (2020). GWs in this region have previously been shown
 137 to be strongly visible in AIRS (Hindley et al., 2020) and aircraft (Wright & Banyard,
 138 2020) observations, but not in limb sounder observations (Geller et al., 2013; Ern et al.,
 139 2018), suggesting a strong role for long-vertical-short-horizontal-wavelength GWs of the
 140 type this model should be well-configured to accurately resolve. The region is also likely
 141 to contain non-orographic GW activity from jets, fronts and spontaneous geostrophic
 142 adjustment processes around the edge of the wintertime stratospheric polar vortex. This
 143 region and time period therefore presents an ideal opportunity to investigate the real-
 144 ism of resolved GWs in the high resolution IFS simulation compared to observations and
 145 to the lower resolution reanalyses.

146 The data sets used in this study are described in Section 2. In Section 3, the meth-
 147 ods for resampling the models as AIRS and calculating the GW properties are described.
 148 The results of the comparison between the resampled models and AIRS observations are
 149 presented in Section 4. These results are discussed in Section 5, and the summary and
 150 conclusions are presented in Section 6.

151 2 Data

152 2.1 AIRS

153 Stratospheric temperature data were used from the Atmospheric Infrared Sounder
 154 (AIRS) instrument on NASA’s Aqua satellite (Hoffmann & Alexander, 2009). The Aqua
 155 satellite’s orbit is sun-synchronous and near-polar, with a period of 98.8 minutes. This
 156 allows AIRS to obtain data with near global coverage over a day. AIRS has 2378 chan-
 157 nels which measure infrared radiation in the wavelength range of 3.7–15.4 μm and 4 chan-
 158 nels that measure near-infrared and visible radiation with a range of 0.4–0.94 μm (Parkinson,
 159 2003). AIRS scans from $+49.5^\circ$ to -49.5° across track, with 90 elements and a swath width
 160 of ~ 1780 km and has a horizontal resolution of ~ 13.5 km \times 13.5 km at nadir which
 161 reduces to 41 km \times 21.4 km at the track edge (Chahine et al., 2006). The data are stored
 162 in granules containing 6 minutes of data, with 240 granules for each day (Aumann et al.,
 163 2003).

164 The 3D temperature data used in this study is calculated from AIRS radiance mea-
 165 surements using the retrieval scheme described by Hoffmann and Alexander (2009). This
 166 retrieval has an improved horizontal resolution by a factor of 3, in comparison with AIRS

operational data, in both the along- and across-track directions, allowing more GW features to be seen in the data. The retrieval uses 4 μm and 15 μm AIRS CO_2 emission channels for nighttime, but only the 15 μm channels are used for the daytime retrievals; specifically 12 15 μm channels are used in the retrieval for daytime and nighttime and 23 4 μm channels for nighttime. In daytime, the radiance measurements for the 4 μm channels are affected by non-LTE (local thermodynamic equilibrium) effects due to solar excitation, so these channels are not used. In the middle and upper stratosphere, few of the 15 μm channels are sensitive to temperature perturbations and therefore, GWs, compared to the 4 μm channels. The estimated total retrieval error of the temperature measurements is 1.6–3.0 K for altitudes from 20 to 60 km. The retrieved temperatures have a vertical resolution of ~ 7 –15 km (Hoffmann & Alexander, 2009). Figure 2a–c of Hindley et al. (2019) show estimated AIRS temperature retrieval errors due to noise and vertical resolution with altitude. An altitude range of 27–54 km, was chosen for the AIRS retrieval data used in this study, because nighttime data have relatively low noise and high resolution in this range.

2.2 High resolution IFS simulation (TCO7999)

The high resolution run of the ECMWF IFS used in this study is a global, hydrostatic simulation, based on version CY45R1 of the IFS atmospheric model (ECMWF, 2023), and run at a TCO7999 resolution (Wedi et al., 2020; Polichtchouk et al., 2022). This resolution has a horizontal grid spacing of 1.25 km at the equator, with an average of 1.4 km globally. In this paper, the simulation is referred to as the 1 km IFS run. ECMWF’s operational 10 day forecasts, at the time of writing, use the IFS at a resolution of 9 km with deep convection parameterization.

The CY45R1 version of the IFS has 137 model levels, at heights from 0.01 hPa down to the surface and the spacing between the levels increases with altitude (Wedi et al., 2020). The smallest GWs are likely to be strongly damped by numerical diffusion in the IFS. To prevent wave reflection at the top of the model, the IFS has a weak sponge layer from 10 hPa to the model top, which only has a small effect on resolved waves, and a very strong sponge layer above 1 hPa (Polichtchouk et al., 2023). The contribution of the GW drag parameterizations is designed to reduce as the horizontal resolution of the model is increased, and is zero at an average grid spacing of 1.4 km. At this resolution, the simulation did not use deep convection parameterizations. The 1 km IFS simulation was initialised on 1st November 2018 00:00 UTC, integrated for 4 months, and ran with a time step of 60 s and a model output frequency of 3 hours. The temperature structure and background flow of the 1 km IFS remain similar to IFS simulations run for the same time period at 3.9 km and 7.8 km horizontal resolutions during the the first 15 days of the simulation (Polichtchouk et al., 2022). Polichtchouk et al. (2022, 2023) investigated the effect of the increase in horizontal resolution from ~ 9 to ~ 1 km and the deep convection parameterization and found GWs are still under-resolved at a grid spacing of ~ 9 km, compared to GWs at the ~ 1 km resolution.

In this study, 3 hourly 1 km IFS temperature data were interpolated onto a regular longitude-latitude grid, with a resolution of $0.1^\circ \times 0.1^\circ$. This reduced resolution was chosen to make the data easier to use and should not affect the results, as this is still a significantly higher horizontal resolution than the AIRS retrieval.

2.3 ERA5

The ECMWF ERA5 is a 5th generation global reanalysis, run from 1940 to the present (Copernicus Climate Change Service, 2023). ERA5 uses 4D-Var (4D variational) data assimilation which combines observations, including AIRS data, and hindcasts (past weather forecasts). The observations and hindcasts are combined in space and time within 12 hour assimilation windows (ECMWF, 2021). The hindcasts used in the data assimilation are

217 from the ECMWF IFS CY41R2 (ECMWF, 2023), implemented in 2016, at TCo1279 res-
 218 olution (9 km average horizontal grid spacing globally). ERA5 has the same model lev-
 219 els and sponge layers as in the CY45R1 version of the IFS (ECMWF, 2021, 2020). The
 220 ERA5 temperature data used were downloaded from the Copernicus Climate Data Store,
 221 for every 3 hours during the time period investigated, on a regular latitude-longitude grid
 222 with a resolution of 0.25° (Copernicus Climate Change Service, 2023).

223 3 Methods

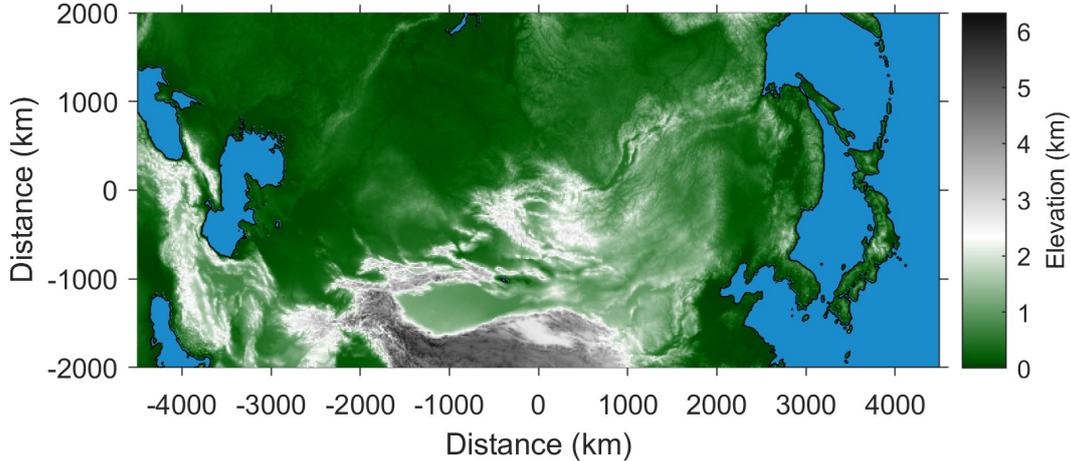


Figure 1. Map of topography in the region investigated on a regular distance grid centred at 52° latitude, 94° longitude. Coastlines are shown in black.

224 Data from the first 10 days of November 2018 for AIRS, the 1 km IFS run and the
 225 ERA5 reanalysis were used, as the 1 km IFS run was initialized on the 1st of this month
 226 at 00:00 UTC. This time period was chosen, because the background temperature and
 227 wind structure, which affect the generation and propagation of GWs, of the 1 km IFS
 228 (Polichtchouk et al., 2022), is expected to remain similar to observations in this period.
 229 This assumption is investigated in section 3.1.3.

230 Where higher magnitude temperature perturbations, indicating GWs, are present,
 231 the data are expected to have a greater variance in general. During this time period, the
 232 AIRS granules with the highest variances were located in Asia and surrounding areas,
 233 suggesting stronger GW activity. Hence, this study focuses on data from this region (shown
 234 in Figure 1). Variances of AIRS temperature perturbations were also used in Hoffmann
 235 et al. (2013) to identify individual GW events.

236 In this study all results use data at 39 km altitude in AIRS and the resampled mod-
 237 els as this is at the centre of the AIRS usable height range (see Figure 2 of Hindley et
 238 al. (2019)). This is also in the altitude range where the AIRS retrieval vertical resolu-
 239 tion is greater and the noise is lower for nighttime data in polar winter, mid latitudes
 240 and the tropics. The results are presented only for nighttime data, due to the the lower
 241 vertical resolution and higher retrieval error of the daytime AIRS retrieval.

242 3.1 Resampling Methods

243 In this study, the observational filter of the AIRS retrieval was applied to the 1 km
 244 IFS run and ERA5, to remove GWs outside of the horizontal and vertical wavelength
 245 ranges in which these waves can be seen in the observations. Data were found from the

246 1 km IFS run and ERA5 at the closest time of the 3 hourly 1 km IFS data used to the
 247 measurement time of each AIRS granule and were resampled as that granule. The mod-
 248 els were not interpolated to the AIRS measurement times as this would smooth out small
 249 scale structures such as GWs (Wright & Hindley, 2018). Two different methods were used
 250 to resample the 1 km IFS run as AIRS. The first resampling method was run on a desk-
 251 top computer, whereas the second is more computationally expensive (Wright & Hind-
 252 ley, 2018) and required the use of HPC.

253 *3.1.1 Method 1*

254 The first method, referred to as method 1 in this paper and applied to both the
 255 1 km IFS run and ERA5, is described by Hindley et al. (2021). The 1 km IFS data, which
 256 was previously interpolated onto a regular longitude-latitude grid with a spacing of 0.1° ,
 257 were selected at the closest 3 hourly time to each AIRS granule. This was first interpo-
 258 lated onto a regular distance grid in the horizontal with a point spacing of 2.7 km, which
 259 is a higher resolution than any part of the original 0.1° longitude-latitude grid spacing
 260 in the region investigated. The data were then smoothed to the approximate horizon-
 261 tal resolution of AIRS at track-centre, using a Gaussian with a FWHM (full width at
 262 half maximum) of $13.5 \text{ km} \times 13.5 \text{ km}$. Following this, the data were interpolated onto
 263 the location of the AIRS granule. The data was then interpolated to a regular distance
 264 spacing in the vertical of 0.1 km from 26 to 55 km altitude, so it could be smoothed to
 265 the vertical resolution of the AIRS retrieval. As the vertical resolution of the retrieval
 266 varies with altitude, the whole volume of data were smoothed in the vertical using a Gaus-
 267 sian function with a different FWHM for each AIRS altitude, from 27 to 54 km, with
 268 a 3 km point spacing. Different arrays of values for the FWHM at each altitude were used
 269 (shown in Figure 2, of Hindley et al. (2019)) for each of the granules containing mostly
 270 nighttime data, depending on whether they are located mostly in the tropics, midlat-
 271 itudes or polar region. The nearest horizontal levels to each altitude were then found and
 272 stored in a separate array, which is the model data resampled as AIRS.

273 ERA5 is also resampled using this method, but since the data has a lower horizon-
 274 tal resolution than AIRS, it is not interpolated to a regular distance grid and smoothed
 275 to the horizontal resolution of AIRS before it is interpolated to the AIRS granule loca-
 276 tion. The 1 km IFS run and ERA5 resampled using this method are referred to as IFS
 277 1 and ERA5 1 in this paper.

278 *3.1.2 Method 2*

279 The second method used to resample the 1 km IFS run as AIRS is described by
 280 Wright and Hindley (2018) and referred to as method 2 in this paper. This involves over-
 281 sampling the model data, onto a grid with a spacing of 1 km in the along and across track
 282 directions and $1/20$ of a decade of pressure in the vertical. These values were selected
 283 based on sensitivity testing discussed in Appendix B of Wright and Hindley (2018). Each
 284 oversampled point was then weighted by the estimated instrument sensitivity at each point
 285 and summed to produce a sample corresponding to each AIRS measurement. This aims
 286 to improve the accuracy in comparison with interpolating the model to the centre of the
 287 satellite measurement volume. Compared to interpolating to a single point, Wright and
 288 Hindley (2018) showed that this method lead to improvements in brightness tempera-
 289 ture measurements in AIRS Level 1 data which are significant for small-scale tempera-
 290 ture perturbations caused by GWs. The 1 km IFS data resampled as AIRS using this
 291 method is referred to as IFS 2.

292 *3.1.3 Temperature divergence of the resampled models and AIRS*

293 Figure 2 shows the point-wise correlation (Figure 2a) and RMSD (Figure 2b) be-
 294 tween the temperature at 39 km altitude in the AIRS observations and resampled mod-

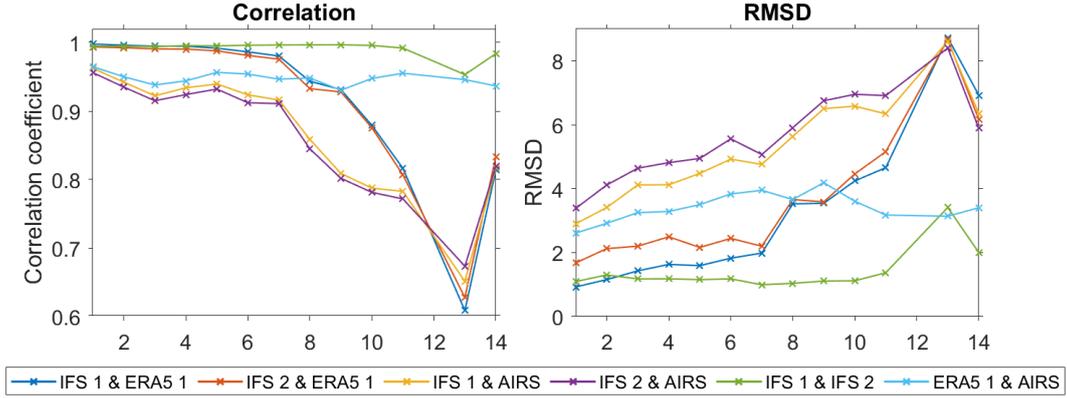


Figure 2. Point-wise correlation coefficients and RMSD between the temperature at 39 km altitude in AIRS, and the 1 km IFS run and ERA5 resampled as AIRS, for each night during the first 14 days of November 2018, within the region shown in Figure 1.

295 els. These are plotted for each night during the first 14 days of November 2018. The data
 296 for the 12th night is missing, because AIRS data were not recorded for most of the re-
 297 gion studied during this night.

298 The point-wise correlation (Figure 2a) between IFS 1 and 2 and AIRS decreases
 299 over time from when the model was initialised up to the 11th night, which is expected
 300 since the free-running 1 km IFS run diverges from the 'truth'. On night 13, the corre-
 301 lation is lower due to a single large anomalous wave covering a large fraction of the re-
 302 gion investigated in the raw model data. The correlation then increases on the 14th night.
 303 As ERA5 assimilates data from observations, including AIRS, the resampled ERA5 data
 304 does not have a decreasing correlation with the AIRS retrieval. Since IFS 1 and 2 are
 305 the same data resampled as AIRS using different methods, the correlation coefficient be-
 306 tween these data sets remains very high, but is lowest on the 13th night.

307 The RMSD (Figure 2b) is greatest between IFS 1 and 2 and AIRS and increases
 308 up to night 11. The RMSD also increases between IFS 1 and 2 and ERA5 1 up to night
 309 11. On the 13th night there is a peak in the RMSD between the 1 km IFS run resam-
 310 pled using both methods and the other data sets, and between IFS 1 and 2 as a result
 311 of the large anomalous wave in the 1 km IFS run data for this night. Due to the corre-
 312 lations and RMSD's shown in Figure 2, data is only used from the 1st – 10th November
 313 2018 for the results presented in Section 4.

314 **3.2 Regridding the data to a regular distance grid and finding temper- 315 ature perturbations**

316 The AIRS granules and model data resampled as each AIRS granule, with data in
 317 the region investigated were regridded onto regular 3D distance grids, as this is required
 318 for 3D spectral analysis. The grids have a horizontal point spacing of ~ 20 km in the across
 319 track direction and ~ 18 km in the along track direction, so that the swath width and
 320 number of across-track and along-track points remain the same after the data has been
 321 regridded, and a vertical spacing of 3 km. Following this, the background is removed from
 322 AIRS and the resampled models using a 4th order polynomial fit in the cross track di-
 323 rection (Wu, 2004; M. J. Alexander & Barnett, 2007).

324 Separately, temperature perturbations were also found for the 1 km IFS run and
 325 ERA5 before resampling as AIRS to allow for comparison to the resampled models and

326 observations (see Section 4.3). The 1 km IFS run and ERA5 were interpolated to a reg-
 327 ular distance grids with a point spacing of 1 km in the vertical and a horizontal point
 328 spacing of 15 km for the 1 km IFS run, and 30 km for ERA5. The background was found
 329 by smoothing both data sets using a Gaussian filter with a convolution kernel size of 11
 330 $\times 11$ points and a standard deviation of 7.15 points. This was then subtracted from the
 331 temperature data to find the perturbations.

332 3.3 Adding AIRS retrieval noise to the resampled models

333 Since temperature perturbations in the AIRS retrieval data can not be separated
 334 from the noise, AIRS noise is added to the temperature perturbations found for the re-
 335 sampled models (ERA5 1, IFS 1 and IFS 2), so that only waves that can still be seen
 336 with AIRS noise added are compared. To find granules containing only noise, the gran-
 337 ules were sorted from lowest to highest variance of the temperature perturbations, and
 338 checked in this order to find granules without subjectively clear visible waves. In Hindley
 339 et al. (2021), temperature perturbations were found for an AIRS overpass containing 2
 340 granules with no waves, which were then randomised at each altitude, and added to the
 341 model resampled as AIRS. This method was also used in Okui et al. (2023) for one AIRS
 342 granule. However, the noise added in Hindley et al. (2021) is uncorrelated pixel-scale noise,
 343 meaning that any noise structures larger than around 30–50 km in the AIRS retrieval
 344 data would not be included. Okui et al. (2023) found that adding noise using this method
 345 resulted in a lower background amplitude due to noise in the resampled model than in
 346 AIRS and may not be suitable for adding AIRS retrieval noise at global scales. There-
 347 fore, in this study a different method was used.

348 30 nighttime AIRS granules containing only noise were selected in total, with 10
 349 in the tropics, mid-latitudes and polar region, respectively. Granules were chosen with
 350 over 10% of the data points in the region investigated. For the tropics and mid-latitudes,
 351 the granules were selected in the first 14 days of November in 2018. As there were not
 352 enough nighttime polar AIRS granules containing only noise in the first two weeks of Novem-
 353 ber 2018, granules were chosen during the first 2 weeks of November in years from 2016
 354 to 2020 for the polar region. The temperature perturbations of the AIRS noise granules
 355 (shown in Supplementary Figures S1–S3) were found using the method described above.
 356 The data in the arrays of temperature perturbations were reversed in the along and across-
 357 track directions separately and saved so that there were 30 noise granules in total for each
 358 group, to increase the number of granules that could be selected. For each granule of model
 359 data resampled as AIRS, an array of noise temperature perturbations was chosen ran-
 360 domly from the corresponding group and added to the resampled model temperature per-
 361 turbations, depending on if the granule contains nighttime data which is mostly from
 362 the tropics ($<30^\circ$), mid-latitudes (30° – 60° latitude) or polar region ($>60^\circ$ latitude).

363 3.4 2D+1 S-Transform

364 The S-Transform (ST) is commonly used for the analysis of GWs (e.g. Fritts et al.,
 365 1998; M. J. Alexander et al., 2008). The 2D+1 ST is based on the 2D S-Transform (Hindley
 366 et al., 2016) and the 3D S-Transform (Wright et al., 2017). The 2D ST and 3D ST are
 367 extensions of the 1D S-Transform (Stockwell et al., 1996).

368 The 2D+1 ST calculates wavelengths using phase shifts between spectral features,
 369 which allows it to measure waves more effectively for 3D data with low resolution in one
 370 dimension compared with the variations in the wave field. Nadir-sensing instruments,
 371 such as AIRS have high horizontal resolution but low vertical resolution. This means there
 372 are a low number of vertical points, for the data from these instruments, in the strato-
 373 sphere in comparison with the point numbers in the horizontal, limiting estimates of the
 374 vertical wavelengths of GWs. For the 2D+1 ST, 2D S-Transforms are found for the data

375 levels in the horizontal and the phase differences between them are calculated (Wright
376 et al., 2021).

377 Using the 2D+1 ST, vertical wavelengths can be calculated more precisely in compar-
378 ison with the 3D ST, and this method is therefore more effective for measuring waves
379 with long vertical wavelengths. A further improvement compared to the 3D ST is that
380 the 2D+1 ST does not quantize vertical wavelengths to Fourier modes, so these wave-
381 lengths vary smoothly in the output. However, the 2D+1 ST is computationally slower.

382 The 2D+1 ST was used to find wave properties for the resampled 1 km IFS run
383 and ERA5 reanalysis, both models before being resampled, and each AIRS granule. The
384 wave amplitude is an output of the 2D+1 ST. The horizontal and vertical wavelengths
385 were calculated using the granule-relative wave frequencies from the 2D+1 ST. The zonal
386 M_x and meridional M_y components of the momentum flux were calculated using the fol-
387 lowing equation derived in Ern et al. (2004),

$$M_x, M_y = -\frac{\rho}{2} \left(\frac{k}{m}, \frac{l}{m} \right) \left(\frac{g}{N_B} \right)^2 \left(\frac{|T'|}{\bar{T}} \right)^2 \quad (1)$$

388 where ρ is the atmospheric density, k , l and m are the wavenumbers in the zonal,
389 meridional and vertical directions respectively, g is the acceleration due to gravity and
390 N_B is the buoyancy frequency, set 0.02 s^{-1} in this study. \bar{T} is the local background tem-
391 perature and $|T'|$ is the amplitude. The wavenumbers are signed to preserve the sign (di-
392 rection) of the zonal and meridional momentum flux components (P. Alexander et al.,
393 2018).

394 The altitude range selected of the model data resampled as AIRS and of the AIRS
395 data used for the 2D+1 ST analysis was 27–54 km. Including altitudes outside of this
396 range with higher noise would affect the wave properties calculated using the 2D+1 ST.
397 The 2D+1 ST was tuned to select waves with horizontal wavelengths ranging from 60
398 to 800 km. Areas where the vertical wavelength is below 6 km or above 45 km were re-
399 moved from the data for the wave properties calculated.

400 4 Results

401 4.1 Example Case Studies

402 Two case studies comparing the resampled models and AIRS are presented in this
403 section. The first case study includes data from AIRS granules with mean times of 19:38–
404 19:50 UTC on the 5th November and the resampled model granules at the closest 3 hourly
405 times to the observations. The temperature perturbations for this example are shown
406 in Figure 3. The top row shows the AIRS swath (Figure 3a) and resampled models with
407 AIRS noise added (Figure 3b–d). The topography of the area is shown in Figure 3e as
408 well as the resampled model swaths before adding noise (Figure 3f–h). $8 \mu\text{m}$ AIRS bright-
409 ness temperatures and ERA5 winds are shown for both case studies in Supplementary
410 Figures S4 and 5.

411 The large wave on the right seen in AIRS and the resampled models (Figure 3a–
412 d) is likely to be orographic, as it is close to a region of higher topography. A curved wave
413 can be seen on the left in AIRS (Figure 3a), which could be convective as it is located
414 close to a region of with a brightness temperature lower than 220 K, indicating deep con-
415 vection (Hoffmann & Alexander, 2010) (see Supplementary Figure S4a). This wave can-
416 not be clearly seen in the resampled IFS (IFS 1 & 2, Figure 3b and c), which suggests
417 that the convective source is missing or significantly reduced in strength in the 1 km IFS
418 run. In ERA5, the wave can be seen but at a lower amplitude, suggesting the convective
419 source was correctly assimilated in the reanalysis.

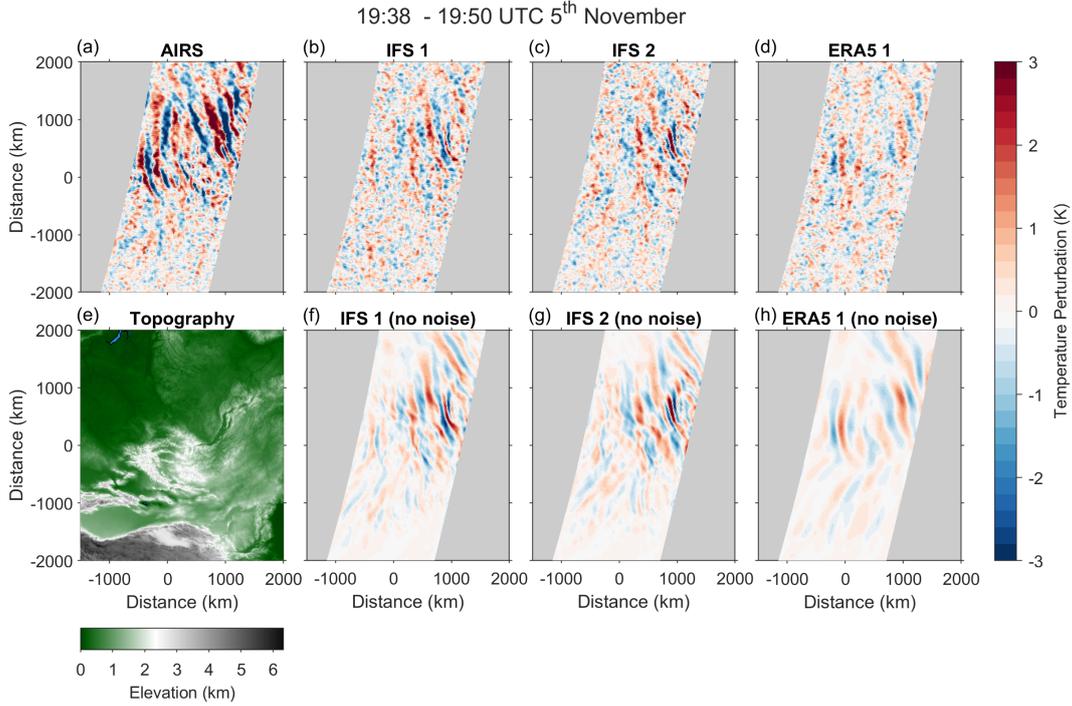


Figure 3. Temperature perturbations for granules in the region shown in Figure 1 with mean times of 19:38–19:50 UTC for AIRS (a) and at the closest time for the resampled models. Panels (b–d) show the resampled models after adding AIRS noise, with the resampled models before adding noise below them (no noise, panels f–h). The section of the topography from Figure 1, for the area is shown in panel (e), with (0,0) at 52° latitude, 94° longitude.

420 Figure 4 shows the wave properties, derived from the 2D+1 S-Transform, for the
 421 AIRS and resampled model granules shown in Figure 3a–d. Areas of higher amplitude
 422 in Figure 4a–d are seen in similar locations in each data set, but are higher in AIRS (Fig-
 423 ure 4a) and lowest in ERA5 1 (Figure 4d). These regions can also be clearly seen for larger
 424 areas in AIRS than in the resampled models. In ERA5 1, an area of higher amplitude
 425 can be seen at the left of the granule stripe, near the centre of the y axis which is also
 426 seen in AIRS, but is not as clearly seen in IFS 1 and 2 (Figure 4b and c). Longer hor-
 427 izontal and vertical wavelengths can be seen in IFS 2 (Figure 4g and k), than in IFS 1
 428 (Figure 4f and j), and in areas of noise the horizontal wavelengths are lower for all the
 429 data sets (Figure 4e–h). The magnitudes of the zonal and meridional momentum fluxes
 430 are highest in AIRS (Figure 4m and q), and lowest in ERA5 1 (Figure 4p and t).

431 The second case study is shown for AIRS granules with mean times of 19:14–19:26
 432 UTC on the 9th November 2018, further in time from when the 1 km IFS run was ini-
 433 tialised, and the resampled model granules at the closest 3 hourly times. Temperature
 434 perturbations for this case study are shown in Figure 5. As in Figure 3, the resampled
 435 model temperature perturbations are shown with AIRS noise (Figure 5b–d) and before
 436 adding AIRS noise (Figure 5f–h). The GWs shown in Figure 5a–d are likely to have oro-
 437 graphic sources as they are close to regions of higher topography and not located near
 438 to a region of deep convection (shown in Supplementary Figure S4b).

439 In Figure 6, the wave properties are shown for the second case study (the AIRS
 440 swath, and models resampled as the swath in Figure 5a–d). The wave properties found
 441 for the resampled 1 km IFS run appear to agree less well with the AIRS observations,

442 compared to the first case study in Figure 4. By the 9th November, the 1 km IFS run
 443 would have diverged further from reality compared to the first case study on the 5th Novem-
 444 ber, which is closer to the time the simulation was initialised. The amplitudes in ERA5
 445 1 (Figure 6d) are higher than in the example shown in Figure 4d. In the areas where the
 446 amplitude is higher in IFS 1 and 2 (Figure 6b and c), the horizontal wavelengths (Fig-
 447 ure 6f and g) are longer than in AIRS (Figure 6e) for the same locations. In this case
 448 study, the horizontal wavelengths in the area with greater amplitudes in ERA5 1 (Fig-
 449 ure 6h) appear to be more similar to the AIRS observations (Figure 6e) than IFS 1 and
 450 2 (Figure 6f and g). This is also expected as ERA5 assimilates observations, unlike the
 451 1 km IFS run.

452 The vertical wavelengths in IFS 1 and 2 (Figure 6j and k) are shorter than in AIRS
 453 (Figure 6i) for the areas where GWs can be seen. In the areas with higher amplitude in
 454 ERA5 1 (Figure 6d), the vertical wavelengths are generally longer (Figure 6l) than in
 455 the same locations for AIRS (Figure 6i). The magnitudes of the zonal and meridional
 456 momentum fluxes are also highest in AIRS for this case study (Figure 6m and q), due
 457 to the higher wave amplitudes, but they are lower in the resampled 1 km IFS run (IFS
 458 1 and 2 in Figure 6n, o, r and s) than in ERA5 1 (Figure 6p and t). In both case stud-
 459 ies, (Figures 4m–t and 6m–t) the zonal and meridional momentum flux is negative in ar-
 460 eas where the amplitude is highest for all the data sets.

461 4.2 Time mean maps

462 Figure 7a–d shows the mean nighttime amplitudes during the first 10 days of Novem-
 463 ber 2018 for the region in Figure 1. AIRS amplitudes (Figure 7a) are divided by a fac-
 464 tor of 2 before plotting, so that areas with higher amplitudes in the resampled models
 465 can be seen more clearly using the same color bar. Areas of higher amplitude are seen
 466 in similar locations in AIRS and the resampled models. However, the amplitudes are sig-
 467 nificantly higher in AIRS than in the resampled 1 km IFS run (Figure 7c and d) and ERA5
 468 1 (Figure 7b) and are lower in ERA5 1 than in IFS 1 and 2. The regions of higher wave
 469 amplitude are located near to mountain ranges, which can be seen in Figure 1 as areas
 470 of higher elevation, suggesting that the GWs have orographic sources. There is an area
 471 of higher amplitude over the Urals in Russia which can be seen in AIRS and the resam-
 472 pled IFS (1 and 2), but is not clearly seen in ERA5 1. The locations of the peaks in am-
 473 plitude in AIRS and the resampled models are consistent with Hindley et al. (2020) and
 474 Wright and Banyard (2020). The maximum mean amplitudes in IFS 1 (Figure 7c) and
 475 IFS 2 (Figure 7d) are a factor of 2.5 and 2.2 lower than in AIRS respectively, so the max-
 476 imum mean amplitude for the 1 km IFS run is a factor of 2.4 lower than in AIRS, av-
 477 eraging the results from the two resampling methods. The maximum amplitude in ERA5
 478 (Figure 7b) is a factor of ~ 2.8 lower than in AIRS.

479 The mean nighttime zonal and meridional momentum fluxes are shown in Figure
 480 7e–h and 7i–l respectively. Like the amplitude, the magnitude of the mean zonal and merid-
 481 ional momentum flux is significantly higher in AIRS (Figure 7e and i) than in the re-
 482 sampled models (Figure 7f–h and j–l). The zonal momentum flux generally has a higher
 483 magnitude than the meridional momentum flux, which is expected based on the AIRS
 484 retrieval climatology in Hindley et al. (2020) and due to background wind filtering. In
 485 areas where the amplitude is higher, the zonal momentum flux is negative (westward)
 486 and the meridional momentum flux is also negative (southward). This suggests that the
 487 highest amplitude GWs are formed by wind flowing over the northeast–southwest aligned
 488 topography shown in Figure 1.

489 The mean horizontal wavelengths for nighttime data in the first 10 days of Novem-
 490 ber 2018 are also shown in Supplementary Figure S6.

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4.3 Kernel distribution functions

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Kernel distribution functions (KDFs) for the amplitudes, horizontal and vertical wavelengths and momentum flux are shown for nighttime data during the first 10 days of November 2018 in Figure 8. In the first column (Figure 8a, c, e and g), the KDFs for AIRS, and ERA5 and the 1 km IFS run (IFS) before resampling as AIRS are shown and the second column (Figure 8b, d, f and h) shows the KDFs for AIRS and the resampled models. KDFs were chosen to show the distributions of the data rather than probability density functions, due to the noise in AIRS and added noise in the resampled models. The distributions were plotted using a normal kernel smoothing function.

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Areas of noise in the AIRS and resampled model data were reduced by smoothing the amplitude measurements using a 7 by 7 point boxcar filter and removing points in the original data where the amplitude of the smoothed data is below the 70th percentile for nighttime data in each data set (~ 0.90 K for IFS 1 and 2, ~ 1.18 K for AIRS and ~ 0.89 K for ERA5 1). For the KDFs of ERA5 and the 1 km IFS run before being resampled as AIRS, the same method was used to remove areas of the data where the amplitude is below the 70th percentile of AIRS. As the amplitudes are lower in the resampled models than in AIRS, the 70th percentile amplitude cutoffs are lower, so less noise is removed. This means more points from areas containing only noise are included in the KDFs for the resampled models. These areas generally have lower amplitudes and therefore lower momentum fluxes, as well as lower horizontal wavelengths (shown in Figures 4 and 6), so there will be a greater proportion of the resampled model data with lower values for these wave properties.

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After being resampled as AIRS, the horizontal and vertical wavelength spectra of the models are more similar to the AIRS data (Figure 8d and f), compared to the spectra before resampling (Figure 8c and e), suggesting that the resampling methods used allow a fairer comparison between the observations and models to be made. There are a higher proportion of points with shorter horizontal wavelengths in the resampled models than in AIRS (Figure 8d), but this could be due to the 70th percentile amplitude cutoffs used. There is some quantization of the horizontal wavelengths longer than around 200 km, leading to multiple peaks in the KDFs for all data sets in Figure 8c and d. This is a result of these waves being approximated by Fourier modes in the 2D+1 ST, since they are long relative to the data size. For the resampled models, the peaks in the vertical wavelength are offset from the observations by around 2–3 km (Figure 8f). This could be due to the values of the AIRS vertical resolution, used for both resampling methods, being overestimated, but the vertical wavelengths in the resampled models could be too short, and the different 70th percentile amplitude cutoffs used for each data set could also affect these results.

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Before resampling there are a greater proportion of higher amplitude GWs in the 1 km IFS run, up to ~ 17 K, than in the AIRS observations, where the KDF tails off at ~ 15 K, and significantly higher fraction of lower amplitude waves in ERA5, where the KDF tails off at around ~ 9 K (Figure 8a). After the AIRS observational filter is applied to the models, the wave amplitudes are generally higher in AIRS, and there is a greater fraction of data points with lower amplitudes in ERA5 1, where the KDF decreases to ~ 4 K, than in IFS 1 and 2, which have KDFs that decrease to ~ 6 K (Figure 8b). Whilst the horizontal momentum flux KDFs are similar for AIRS and the 1 km IFS run before resampling (Figure 8g), they are generally higher in AIRS than in the resampled models (Figure 8h), as the momentum flux is proportional to vertical wavelengths and the square of the amplitudes. The results shown in Figure 8b, d and f suggest that there are high amplitude GWs with longer horizontal or vertical wavelengths in the AIRS retrieval data that are not present in the resampled models.

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4.4 Bivariate histograms

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Bivariate histograms are shown in Figure 9 to compare the wave properties in the resampled models and AIRS. The color bars show the normalised density, i.e. the number of counts in each bin divided by the total number of counts. Areas of noise in the data used for the bivariate histograms were also reduced using the same method as for the kernel distribution functions (Figure 8), described by the previous section. The number of points used to plot the histograms for each comparison can be found in Table S1 in the appendix. These values vary as points are only included if both data sets do not have a missing value in the point location. Values will be missing if the vertical wavelength in the point location is lower than 6 km or greater than 45 km, or the amplitude at that location is below the 70th percentile amplitude cutoff for each data set. Table S1 also shows the fraction of points above (f.a) and below (f.b) the 1:1 line (grey dashed line) in Figure 9.

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The amplitudes in nighttime AIRS data are significantly higher than in the resampled models (Figure m, q and u), with f.b ranging from 0.815 (AIRS & IFS 2) to 0.859 (AIRS & ERA5 1) (Table S1). Stripes with no data can be seen in the bivariate histograms of the horizontal wavelengths (Figure 9b, f, j, n, r, and v), as the horizontal wavelengths are quantized at longer wavelengths. There are also more points where the AIRS data has a longer vertical wavelength than in the resampled models (Figure 9o, s and w) with f.b ranging from 0.557 (AIRS & ERA5 1) to 0.668 (AIRS & IFS 2). The histograms also show higher momentum fluxes in AIRS (Figure 9p, t and x) as a result of the higher GW amplitudes and vertical wavelengths. For the momentum flux (Figure 9d, h, l, p, t and x), the points in the bivariate histograms are very spread out suggesting there is little point-wise correlation between the data sets except from between IFS 1 and 2 (Figure 9d). The data points are also quite spread out for the vertical wavelength plots (Figure 9c, g, k, o, s, and w) indicating a low point-wise correlation. Data points are closer to the 1:1 line for IFS 1 and 2 for higher amplitude values (Figure 9a and d) and momentum flux. The fraction of points above and below the 1:1 line for IFS 1 & IFS 2 (Figure 9a–d) are similar for all wave properties shown (see Table S1) with the greatest difference in the fractions for the horizontal wavelength where f.a is 0.557 and f.b is 0.443.

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5 Discussion

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The methods used in this study have allowed a fairer comparison between the models and AIRS observations compared to previous work. However some issues still remain, including how noise is selected and added, AIRS’s observational filter and the amplitude cutoffs used.

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The AIRS temperature perturbations containing only noise were selected by ordering the granules from lowest to highest variance and selecting granules manually for different regions in nighttime data, during the same period of the year as the data used in this study. This meant the noise added to the resampled models was better correlated to the location of each resampled model granule and could include large noise structures compared to previous methods used (e.g. in Hindley et al., 2021; Okui et al., 2023). However this method of selecting noise granules would be too time consuming for a longer data set, so using machine learning to identify whether granules contain only noise or GWs could be a better approach.

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A limitation of this work is that only GWs in the 1 km IFS run with wavelengths in the AIRS’s observational filter can be compared to AIRS observations. The AIRS retrieval data has a low vertical resolution. In future work, data from instruments with different observational filters, such as limb sounders or satellites using GPS radio occultation, could be used to validate some of the resolved GWs in the 1 km IFS run with

590 wavelength ranges outside of AIRS’s observational filter. Limb sounders have a low hor-
 591 izontal resolution, but higher vertical resolution than nadir sounders like AIRS.

592 Whilst the location and timing of the GWs agree well in the resampled 1 km IFS
 593 run and AIRS observations, the mean amplitudes were found to be significantly lower
 594 in the resampled 1 km IFS run, by a factor of ~ 2.4 , but higher than in the lower reso-
 595 lution ERA5 reanalysis. As a result of this, the horizontal momentum flux is also lower
 596 in the resampled models compared to the observations. Kruse et al. (2022) found that
 597 GW amplitudes in a lower resolution run of the IFS, with a grid-spacing, of ~ 9 km, were
 598 lower than in AIRS observations. High amplitude GWs, seen in the AIRS observations,
 599 are not found in the resampled 1 km IFS data (Figure 8), but are present in the 1 km
 600 IFS run before resampling suggesting these waves have wavelengths outside of AIRS’s
 601 observational filter.

602 Amplitude cutoffs were used to reduce areas of noise included in the AIRS retrieval
 603 and resampled model data for the kernel distribution functions (Figure 8) and bivari-
 604 ate histograms (Figure 9), but this could not remove all areas of noise without also re-
 605 moving areas of low amplitude GWs. This means that these results will be affected by
 606 the remaining noise. These cutoffs were chosen by finding the 70th percentile of all night-
 607 time data during the first 10 days of November 2018 for each data set. Due to the lower
 608 wave amplitudes in the resampled models, the 70th percentiles were lower than for the
 609 AIRS retrieval data, so more areas of noise are included in the results for the resampled
 610 models.

611 The two methods used to resample the 1 km IFS lead to quite similar results and
 612 worked effectively to smooth the model data to AIRS’s resolution, resulting in more sim-
 613 ilar distributions of GW horizontal and vertical wavelengths, compared to before the ob-
 614 servational filter was applied (Figure 8c–f). However, the peaks of the distributions of
 615 vertical wavelength for the resampled models were found to be around 2–3 km lower than
 616 for AIRS (Figure 8f). This could be a result of the AIRS resolution values, used to smooth
 617 the model data, being overestimated, but could also be due to differences in the verti-
 618 cal wavelengths in the resampled models compared to the observations. These results
 619 will also be affected by the noise in the data.

620 **6 Summary and Conclusions**

621 In this study, gravity wave (GW) properties in a ~ 1.4 km gravity-wave-resolving
 622 run (TCO7999 resolution) of the IFS are compared to AIRS observations over Asia and
 623 surrounding regions, using nighttime data during the first 10 days of November 2018.
 624 The results show a good level of fidelity for the model by comparison to the observations,
 625 but with important differences, discussed below.

626 Two different methods were used to resample the 1 km IFS run to facilitate this
 627 comparison, the first method by smoothing to AIRS’ resolution and then interpolating
 628 to the measurement location, and the second by oversampling the model and then pro-
 629 ducing a weighted average of the oversampled points. Although small differences are seen,
 630 they generally produce quite similar results. Since method 2 is significantly more com-
 631 putationally expensive, method 1 may be better suited for comparing models to AIRS.
 632 This result does not necessarily hold in the general case: the large- and vertically-deep
 633 volume of nadir measurements, such as those from AIRS, is likely less sensitive to foot-
 634 print positioning and morphology than finer-vertical wider-horizontal measurements such
 635 as from limb sounders, and this will be investigated further in future work.

636 Based on these results, the output from the ERA5 reanalysis was also resampled
 637 as AIRS using method 1, to see how well the 1 km IFS run resolves GWs in compari-
 638 son to this lower resolution (and slightly chronologically older) model with assimilative
 639 capabilities. Noise derived from wave-free AIRS observations was also added to the sim-

640 ulated data to produce a fairer comparison with the very noisy observations, following
 641 experience in Okui et al. (2023) which showed the significant effect such noise has on 1:1
 642 comparisons. Finally, the 2D+1 S-Transform analysis of Wright et al. (2021) was used
 643 to find the wave properties for each data set.

644 The results of this analysis lead to the following conclusions:

- 645 1. GWs in the 1 km IFS run can be seen at similar locations and times, and with
 646 similar wave morphology to AIRS, suggesting that the model works well in this
 647 regard. ERA5 waves are in general less morphologically consistent with observa-
 648 tions, and in particular often have inconsistently long horizontal wavelengths,
 649 but do occur at similar locations and times to the observations in many if not most
 650 cases.
- 651 2. Measured amplitudes and momentum fluxes are significantly lower in both resam-
 652 pled models than in AIRS data, with ERA5 amplitudes slightly lower (and thus
 653 less observationally-consistent) than those in the 1 km model. This difference is
 654 large, with a long tail of high-amplitude AIRS measurements (Figure 8b) which
 655 in turns drives a similar difference in momentum fluxes. Investigation of the raw
 656 model data shows that many high-amplitude waves in the 1 km IFS run have wave-
 657 lengths too horizontally-short for AIRS to observe (see e.g. Figure 8), which are
 658 thus not seen in the resampled model. Given that the overall amplitude and mo-
 659 mentum flux distributions (Figure 8a,g) in the raw models are broadly similar to
 660 AIRS, this may suggest that wave activity in the model has plausible total am-
 661 plitudes and fluxes, but skewed to much shorter wavelengths than in the true GW
 662 spectrum.
- 663 3. Vertical wavelengths in both ERA5 and the 1 km IFS run are significantly shorter
 664 than in AIRS observations, even after resampling to match the observational res-
 665 olution. This difference is typically $\sim 2-3$ km, i.e. approximately 10-20% of the ob-
 666 served wavelengths. A similar effect is seen in horizontal wavelength, with resam-
 667 pled model wavelengths showing a sharp peak at short wavelengths while observed
 668 horizontal wavelengths peak at a slightly higher value but exhibit a much longer
 669 and flatter distribution. Both of these conclusions are difficult to decouple from
 670 the effects of (both real and synthesised) measurement noise, and further work is
 671 needed to address this question more carefully.

672 This work highlights the importance of carefully applying the observational filter
 673 of the observing platform to models before comparing GWs in simulations to those in
 674 observations, which is shown to be necessary for producing a meaningful comparison in
 675 this study. This is important for accurate testing of how well GWs are resolved in high
 676 resolution models, with further implications for parameterization development, as this
 677 increasingly frequently uses high-resolution models of this nature as a ‘truth’ for tun-
 678 ing purposes.

679 Open Research

680 The AIRS temperature data used in the study were computed from AIRS radiances
 681 using the retrieval scheme described in Hoffmann and Alexander (2009). The 3D AIRS
 682 temperature retrieval can be obtained at [https://datapub.fz-juelich.de/slcs/airs/
 683 gravity_waves](https://datapub.fz-juelich.de/slcs/airs/gravity_waves). The ECMWF ERA5 reanalysis data at 0.25° resolution can be down-
 684 loaded from the Copernicus Climate Data Store at [https://cds.climate.copernicus
 685 .eu/cdsapp#!/dataset/reanalysis-era5-pressure-levels?tab=overview](https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-pressure-levels?tab=overview) (Copernicus
 686 Climate Change Service, 2023). For the 1 km IFS run, the size of the raw model out-
 687 put on the native grid is a few hundred TB, so it is not possible for all of the data to be
 688 made available. However, the post processed data will be retained and is available on
 689 request. Code written in MATLAB (available at <https://uk.mathworks.com/products/>

690 matlab.html) was used to resample the models as AIRS, analyse the gravity wave prop-
 691 erties and produce the figures. The MATLAB code used is available at [https://github](https://github.com/Emily-Lear/Comparing_the_1_km_IFS_run_to_AIRS_observations.git)
 692 .com/Emily-Lear/Comparing_the_1_km_IFS_run_to_AIRS_observations.git.

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 704 ness data.

705 References

- 706 Albers, J. R., & Birner, T. (2014). Vortex Preconditioning due to Planetary and
 707 Gravity Waves prior to Sudden Stratospheric Warmings. *J. Atmos. Sci.*, *71*,
 708 4028–4054. doi: 10.1175/JAS-D-14-0026.1
- 709 Alexander, M. J., & Barnet, C. (2007). Using satellite observations to constrain
 710 parameterizations of gravity wave effects for global models. *J. Atmos. Sci.*, *64*,
 711 1652–1665. doi: 10.1175/JAS3897.1
- 712 Alexander, M. J., Geller, M., McLandress, C., Polavarapu, S., Preusse, P., Sassi, F.,
 713 ... Watanabe, S. (2010, JUL). Recent developments in gravity-wave effects
 714 in climate models and the global distribution of gravity-wave momentum flux
 715 from observations and models. *Quart. J. Roy. Meteor. Soc.*, *136*(650, A),
 716 1103–1124. doi: {10.1002/qj.637}
- 717 Alexander, M. J., Gille, J., Cavanaugh, C., Coffey, M., Craig, C., Eden, T., ...
 718 Dean, V. (2008, May). Global estimates of gravity wave momentum flux from
 719 high resolution dynamics limb sounder observations. *Journal of Geophysical*
 720 *Research*, *113*(D15), 0148–0227. Retrieved from [https://doi.org/10.1029/](https://doi.org/10.1029/2007jd008807)
 721 [2007jd008807](https://doi.org/10.1029/2007jd008807) doi: 10.1029/2007jd008807
- 722 Alexander, P., Schmidt, T., & de la Torre, A. (2018). A method to determine grav-
 723 ity wave net momentum flux, propagation direction, and “real” wavelengths:
 724 A gps radio occultations soundings case study. *Earth and Space Science*, *5*(6),
 725 222–230. Retrieved from [https://agupubs.onlinelibrary.wiley.com/doi/](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2017EA000342)
 726 [abs/10.1002/2017EA000342](https://doi.org/10.1002/2017EA000342) doi: <https://doi.org/10.1002/2017EA000342>
- 727 Aumann, H., Chahine, M., Gautier, C., Goldberg, M., Kalnay, E., McMillin, L.,
 728 ... Susskind, J. (2003, Feb). Airs/amsu/hsb on the aqua mission: de-
 729 sign, science objectives, data products, and processing systems. *IEEE*
 730 *Transactions on Geoscience and Remote Sensing*, *41*(2), 253–264. doi:
 731 10.1109/TGRS.2002.808356
- 732 Butchart, N., Charlton-Perez, A. J., Cionni, I., Hardiman, S. C., Haynes, P. H.,
 733 Krüger, K., ... Yamashita, Y. (2011). Multimodel climate and variability
 734 of the stratosphere. *J. Geophys. Res.*, *116*(D5), n/a–n/a. (D05102) doi:
 735 10.1029/2010JD014995
- 736 Carslaw, K. S., Wirth, M., Tsias, A., Luo, A. B. P., Dörnbrack, A., Leutbecher,
 737 M., ... Peter, T. (1998). Increased stratospheric ozone depletion due
 738 to mountain-induced atmospheric waves. *Nature*, *391*, 675–678. doi:
 739 <https://doi.org/10.1038/35589>
- 740 Chahine, M. T., Pagano, T. S., Aumann, H. H., Atlas, R., Barnet, C., Bblaisdell, J.,

- 741 ... Zhou, L. (2006, 07). AIRS: Improving Weather Forecasting and Provid-
 742 ing New Data on Greenhouse Gases. *Bulletin of the American Meteorological*
 743 *Society*, 87(7), 911-926. doi: 10.1175/BAMS-87-7-911
- 744 Copernicus Climate Change Service. (2023). *ERA5 hourly data on pressure levels*
 745 *from 1979 to present*. European Centre for Medium-Range Weather For-
 746 casts (ECMWF), [https://cds.climate.copernicus.eu/cdsapp#!/dataset/](https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-pressure-levels?tab=overview)
 747 [reanalysis-era5-pressure-levels?tab=overview](https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-pressure-levels?tab=overview). ([Accessed 15 Aug
 748 2023])
- 749 ECMWF. (2020). *L137 model level definitions*. [https://confluence.ecmwf.int/](https://confluence.ecmwf.int/display/UDOC/L137+model+level+definitions)
 750 [display/UDOC/L137+model+level+definitions](https://confluence.ecmwf.int/display/UDOC/L137+model+level+definitions). ([Accessed 28 Feb 2022])
- 751 ECMWF. (2021). *ERA5*. <https://confluence.ecmwf.int/display/CKB/ERA5>.
 752 ([Accessed 15 Aug 2023])
- 753 ECMWF. (2023). *IFS Documentation*. [https://www.ecmwf.int/en/publications/](https://www.ecmwf.int/en/publications/ifs-documentation)
 754 [ifs-documentation](https://www.ecmwf.int/en/publications/ifs-documentation).
- 755 Ern, M., Preusse, P., Alexander, M. J., & Warner, C. D. (2004). Absolute values
 756 of gravity wave momentum flux derived from satellite data. *J. Geophys. Res.*,
 757 109, D20103. doi: 10.1029/2004JD004752
- 758 Ern, M., Trinh, Q. T., Preusse, P., Gille, J. C., Mlynczak, M. G., III, J. M. R., &
 759 Riese, M. (2018, April). GRACILE: a comprehensive climatology of atmo-
 760 spheric gravity wave parameters based on satellite limb soundings. *Earth*
 761 *System Science Data*, 10(2), 857–892. doi: 10.5194/essd-10-857-2018
- 762 Fritts, D. C., & Alexander, M. J. (2003). Gravity wave dynamics and ef-
 763 fects in the middle atmosphere. *Reviews of Geophysics*, 41, 1003. doi:
 764 10.1029/2001RG000106
- 765 Fritts, D. C., Riggins, D. M., Balsley, B. B., & Stockwell, R. G. (1998). Recent re-
 766 sults with an mf radar at mcmurdo, antarctica: Characteristics and variability
 767 of motions near 12-hour period in the mesosphere. *Geophys. Res. Lett.*, 25(3),
 768 297–300. doi: 10.1029/97GL03702
- 769 Geller, M., Alexander, M. J., Love, P., Bacmeister, J., Ern, M., Hertzog, A., ...
 770 Zhou, T. (2013). A Comparison between Gravity Wave Momentum Fluxes in
 771 Observations and Climate Models. *Journal of Climate*, 26, 6383–6405. doi:
 772 10.1175/JCLI-D-12-00545.1
- 773 Harvey, V. L., Randall, C. E., Becker, E., Smith, A. K., Bardeen, C. G., France,
 774 J. A., & Goncharenko, L. P. (2019). Evaluation of the mesospheric polar
 775 vortices in waccm. *Journal of Geophysical Research: Atmospheres*, 124(20),
 776 10626-10645. doi: <https://doi.org/10.1029/2019JD030727>
- 777 Hindley, N. P., Smith, N. D., Wright, C. J., Rees, D. A. S., & Mitchell, N. J. (2016,
 778 June). A two-dimensional stockwell transform for gravity wave analysis of
 779 AIRS measurements. *Atmospheric Measurement Techniques*, 9(6), 2545–2565.
 780 doi: 10.5194/amt-9-2545-2016
- 781 Hindley, N. P., Wright, C. J., Gadian, A. M., Hoffmann, L., Hughes, J. K., Jackson,
 782 D. R., ... Ross, A. N. (2021). Stratospheric gravity waves over the mountain-
 783 ous island of south georgia: testing a high-resolution dynamical model with 3-d
 784 satellite observations and radiosondes. *Atmospheric Chemistry and Physics*,
 785 21(10), 7695–7722. Retrieved from [https://acp.copernicus.org/articles/](https://acp.copernicus.org/articles/21/7695/2021/)
 786 [21/7695/2021/](https://acp.copernicus.org/articles/21/7695/2021/) doi: 10.5194/acp-21-7695-2021
- 787 Hindley, N. P., Wright, C. J., Hoffmann, L., Moffat-Griffin, T., & Mitchell, N. J.
 788 (2020, November). An 18-year climatology of directional stratospheric gravity
 789 wave momentum flux from 3-d satellite observations. *Geophysical Research*
 790 *Letters*, 47(22), e2020GL089557. doi: 10.1029/2020gl089557
- 791 Hindley, N. P., Wright, C. J., Smith, N. D., Hoffmann, L., Holt, L. A., Alexander,
 792 M. J., ... Mitchell, N. J. (2019). Gravity waves in the winter stratosphere
 793 over the southern ocean: high-resolution satellite observations and 3-d spectral
 794 analysis. *Atmospheric Chemistry and Physics*, 19(24), 15377–15414. doi:
 795 10.5194/acp-19-15377-2019

- 796 Hoffmann, L., & Alexander, M. J. (2009). Retrieval of stratospheric temperatures
797 from Atmospheric Infrared Sounder radiance measurements for gravity wave
798 studies. *J. Geophys. Res.*, *114*, D07105. doi: 10.1029/2008JD011241
- 799 Hoffmann, L., & Alexander, M. J. (2010, 10). Occurrence frequency of convective
800 gravity waves during the north american thunderstorm season. *J. Geophys.*
801 *Res.*, *115*(D20). doi: 10.1029/2010JD014401
- 802 Hoffmann, L., Spang, R., Orr, A., Alexander, M. J., Holt, L. A., & Stein, O. (2017).
803 A decadal satellite record of gravity wave activity in the lower stratosphere
804 to study polar stratospheric cloud formation. *Atmospheric Chemistry and*
805 *Physics*, *17*(4), 2901–2920. Retrieved from [https://acp.copernicus.org/](https://acp.copernicus.org/articles/17/2901/2017/)
806 [articles/17/2901/2017/](https://acp.copernicus.org/articles/17/2901/2017/) doi: 10.5194/acp-17-2901-2017
- 807 Hoffmann, L., Xue, X., & Alexander, M. J. (2013). A global view of stratospheric
808 gravity wave hotspots located with Atmospheric Infrared Sounder observa-
809 tions. *J. Geophys. Res.*, *118*, 416–434. doi: 10.1029/2012JD018658
- 810 Jewtoukoff, V., Hertzog, A., Plougonven, R., de la Camara, A., & Lott, F. (2015).
811 Comparison of gravity waves in the southern hemisphere derived from balloon
812 observations and the ecmwf analyses. *Journal of the Atmospheric Sciences*,
813 *72*(9), 3449–3468. doi: 10.1175/JAS-D-14-0324.1
- 814 Kruse, C. G., Alexander, M. J., Hoffmann, L., van Niekerk, A., Polichtchouk, I.,
815 Bacmeister, J. T., ... Stein, O. (2022). Observed and modeled mountain
816 waves from the surface to the mesosphere near the drake passage. *Journal*
817 *of the Atmospheric Sciences*, *79*(4), 909 - 932. Retrieved from [https://](https://journals.ametsoc.org/view/journals/atsc/79/4/JAS-D-21-0252.1.xml)
818 journals.ametsoc.org/view/journals/atsc/79/4/JAS-D-21-0252.1.xml
819 doi: <https://doi.org/10.1175/JAS-D-21-0252.1>
- 820 Lane, T. P., Doyle, J. D., Sharman, R. D., Shapiro, M. A., & Watson, C. D. (2009).
821 Statistics and dynamics of aircraft encounters of turbulence over greenland.
822 *Monthly Weather Review*, *137*(8), 2687 - 2702. Retrieved from [https://](https://journals.ametsoc.org/view/journals/mwre/137/8/2009mwr2878.1.xml)
823 journals.ametsoc.org/view/journals/mwre/137/8/2009mwr2878.1.xml
824 doi: <https://doi.org/10.1175/2009MWR2878.1>
- 825 Lund, T. S., Fritts, D. C., Wan, K., Laughman, B., & Liu, H.-L. (2020). Numer-
826 ical simulation of mountain waves over the southern andes. part i: Mountain
827 wave and secondary wave character, evolutions, and breaking. *Journal of the*
828 *Atmospheric Sciences*, *77*(12), 4337 - 4356. doi: [https://doi.org/10.1175/](https://doi.org/10.1175/JAS-D-19-0356.1)
829 [JAS-D-19-0356.1](https://doi.org/10.1175/JAS-D-19-0356.1)
- 830 Okui, H., Wright, C. J., Hindley, N. P., Lear, E. J., & Sato, K. (2023). A com-
831 parison of stratospheric gravity waves in a high-resolution general circu-
832 lation model with 3-d satellite observations. *Journal of Geophysical Re-*
833 *search: Atmospheres*, *128*(13), e2023JD038795. Retrieved from [https://](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2023JD038795)
834 agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2023JD038795 doi:
835 <https://doi.org/10.1029/2023JD038795>
- 836 Parkinson, C. (2003). Aqua: an earth-observing satellite mission to examine wa-
837 ter and other climate variables. *IEEE Transactions on Geoscience and Remote*
838 *Sensing*, *41*(2), 173–183. doi: 10.1109/TGRS.2002.808319
- 839 Plougonven, R., de la Cámara, A., Hertzog, A., & Lott, F. (2020). How does
840 knowledge of atmospheric gravity waves guide their parameterizations? *Quar-*
841 *terly Journal of the Royal Meteorological Society*, *146*(728), 1529–1543. doi:
842 [10.1002/qj.3732](https://doi.org/10.1002/qj.3732)
- 843 Polichtchouk, I., Shepherd, T. G., Hogan, R. J., & Bechtold, P. (2018, May).
844 Sensitivity of the brewer-dobson circulation and polar vortex variability to
845 parameterized nonorographic gravity wave drag in a high-resolution atmo-
846 spheric model. *Journal of the Atmospheric Sciences*, *75*(5), 1525–1543. doi:
847 [10.1175/jas-d-17-0304.1](https://doi.org/10.1175/jas-d-17-0304.1)
- 848 Polichtchouk, I., van Niekerk, A., & Wedi, N. (2023). Resolved gravity waves in
849 the extratropical stratosphere: Effect of horizontal resolution increase from
850 $o(10)$ to $o(1)$ km. *Journal of the Atmospheric Sciences*, *80*(2), 473 - 486. Re-

- 851 trieved from [https://journals.ametsoc.org/view/journals/atasc/80/2/](https://journals.ametsoc.org/view/journals/atasc/80/2/JAS-D-22-0138.1.xml)
852 JAS-D-22-0138.1.xml doi: <https://doi.org/10.1175/JAS-D-22-0138.1>
- 853 Polichtchouk, I., Wedi, N., & Kim, Y.-H. (2022). Resolved gravity waves in the
854 tropical stratosphere: Impact of horizontal resolution and deep convection
855 parametrization. *Quarterly Journal of the Royal Meteorological Society*,
856 148(742), 233-251. Retrieved from [https://rmets.onlinelibrary.wiley](https://rmets.onlinelibrary.wiley.com/doi/abs/10.1002/qj.4202)
857 .com/doi/abs/10.1002/qj.4202 doi: <https://doi.org/10.1002/qj.4202>
- 858 Preusse, P., Dörnbrack, A., & Eckermann, S. (2002). Space-based measurements of
859 stratospheric mountain waves by CRISTA 1. Sensitivity, analysis method, and
860 a case study. *J. Geophys. Res.*, 107, 8178. doi: 10.1029/2001JD000699
- 861 Sato, K., & Hirano, S. (2019). The climatology of the brewer–dobson circulation
862 and the contribution of gravity waves. *Atmospheric Chemistry and Physics*,
863 19(7), 4517–4539. Retrieved from [https://acp.copernicus.org/articles/](https://acp.copernicus.org/articles/19/4517/2019/)
864 19/4517/2019/ doi: 10.5194/acp-19-4517-2019
- 865 Sato, K., Tateno, S., Watanabe, S., & Kawatani, Y. (2012). Gravity Wave Charac-
866 teristics in the Southern Hemisphere Revealed by a High-Resolution Middle-
867 Atmosphere General Circulation Model. *J. Atmos. Sci.*, 69, 1378-1396. doi:
868 10.1175/JAS-D-11-0101.1
- 869 Stockwell, R. G., Mansinha, L., & Lowe, R. P. (1996, Apr). Localization of the
870 complex spectrum: the s transform. *IEEE Transactions on Signal Processing*,
871 44(4), 998-1001. doi: 10.1109/78.492555
- 872 Vosper, S. B. (2015). Mountain waves and wakes generated by south georgia: impli-
873 cations for drag parametrization. *QJRMS*, 141(692), 2813-2827. doi: 10.1002/
874 qj.2566
- 875 Watanabe, S., & Miyahara, S. (2009). Quantification of the gravity wave forcing
876 of the migrating diurnal tide in a gravity wave–resolving general circulation
877 model. *Journal of Geophysical Research: Atmospheres*, 114(D7), 0148-0227.
878 doi: <https://doi.org/10.1029/2008JD011218>
- 879 Wedi, N. P., Polichtchouk, I., Dueben, P., Anantharaj, V. G., Bauer, P., Boussetta,
880 S., ... Vitart, F. (2020). A baseline for global weather and climate simulations
881 at 1 km resolution. *Journal of Advances in Modeling Earth Systems*, 12(11),
882 e2020MS002192. doi: <https://doi.org/10.1029/2020MS002192>
- 883 Wright, C. J., & Banyard, T. P. (2020). Multidecadal measurements of utls grav-
884 ity waves derived from commercial flight data. *Journal of Geophysical Re-*
885 search: *Atmospheres*, 125(19), e2020JD033445. Retrieved from [https://](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2020JD033445)
886 agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2020JD033445
887 (e2020JD033445 10.1029/2020JD033445) doi: [https://doi.org/10.1029/](https://doi.org/10.1029/2020JD033445)
888 2020JD033445
- 889 Wright, C. J., & Hindley, N. P. (2018). How well do stratospheric reanalyses re-
890 produce high-resolution satellite temperature measurements? *Atmos. Chem.*
891 *Phys.*, 18(18), 13703–13731. doi: 10.5194/acp-18-13703-2018
- 892 Wright, C. J., Hindley, N. P., Alexander, M. J., Holt, L. A., & Hoffmann, L.
893 (2021). Using vertical phase differences to better resolve 3d gravity wave
894 structure. *Atmospheric Measurement Techniques*, 14(9), 5873–5886. Re-
895 trieved from <https://amt.copernicus.org/articles/14/5873/2021/> doi:
896 10.5194/amt-14-5873-2021
- 897 Wright, C. J., Hindley, N. P., Hoffmann, L., Alexander, M. J., & Mitchell, N. J.
898 (2017). Exploring gravity wave characteristics in 3-d using a novel s-transform
899 technique: Amsu/aqua measurements over the southern andes and drake
900 passage. *Atmospheric Chemistry and Physics*, 17(13), 8553–8575. doi:
901 10.5194/acp-17-8553-2017
- 902 Wu, D. L. (2004). Mesoscale gravity wave variances from amsu-a radiances. *Geophys.*
903 *Res. Lett.*, 31(12), 1944-8007. (L12114) doi: 10.1029/2004GL019562

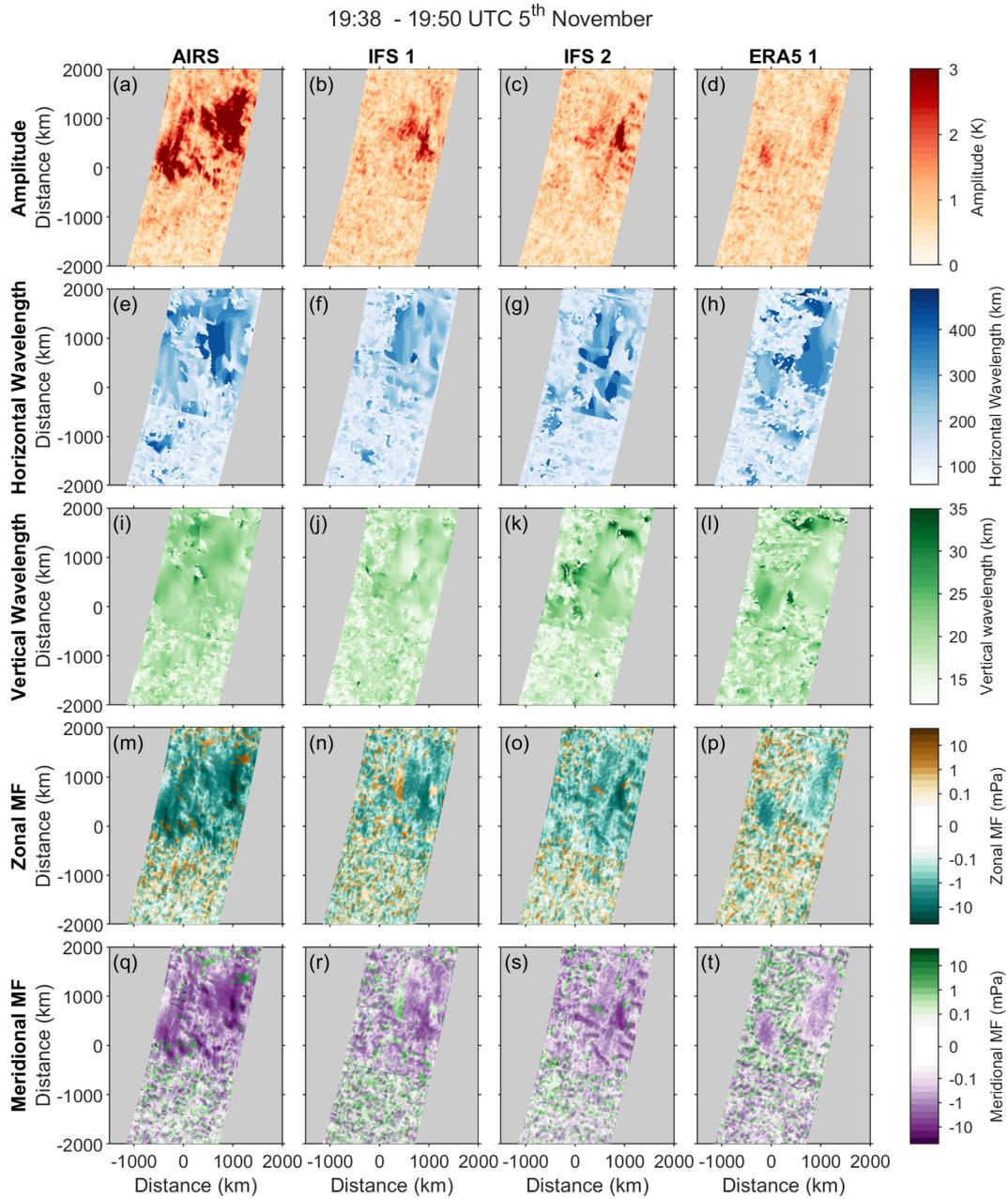


Figure 4. Wave properties for the granules shown in Figure 3, including the amplitude, horizontal and vertical wavelengths and the zonal, and meridional momentum flux (MF). The zonal and meridional momentum flux are shown on a log color scale. Data points were removed where the vertical wavelength is below 6 km or above 45 km.

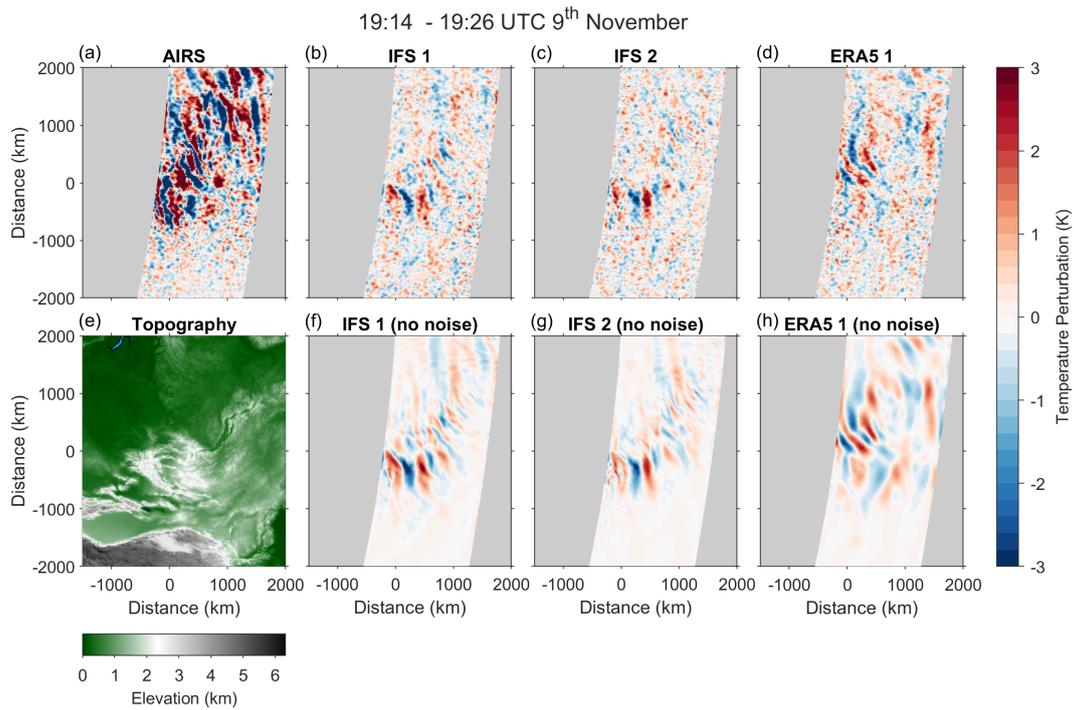


Figure 5. As in Figure 3, but for AIRS granules with mean times from 19:14 to 19:26 UTC on the 9th November 2018 (a) and the resampled models at the closest times (b–d) with and (f–h) without AIRS noise added.

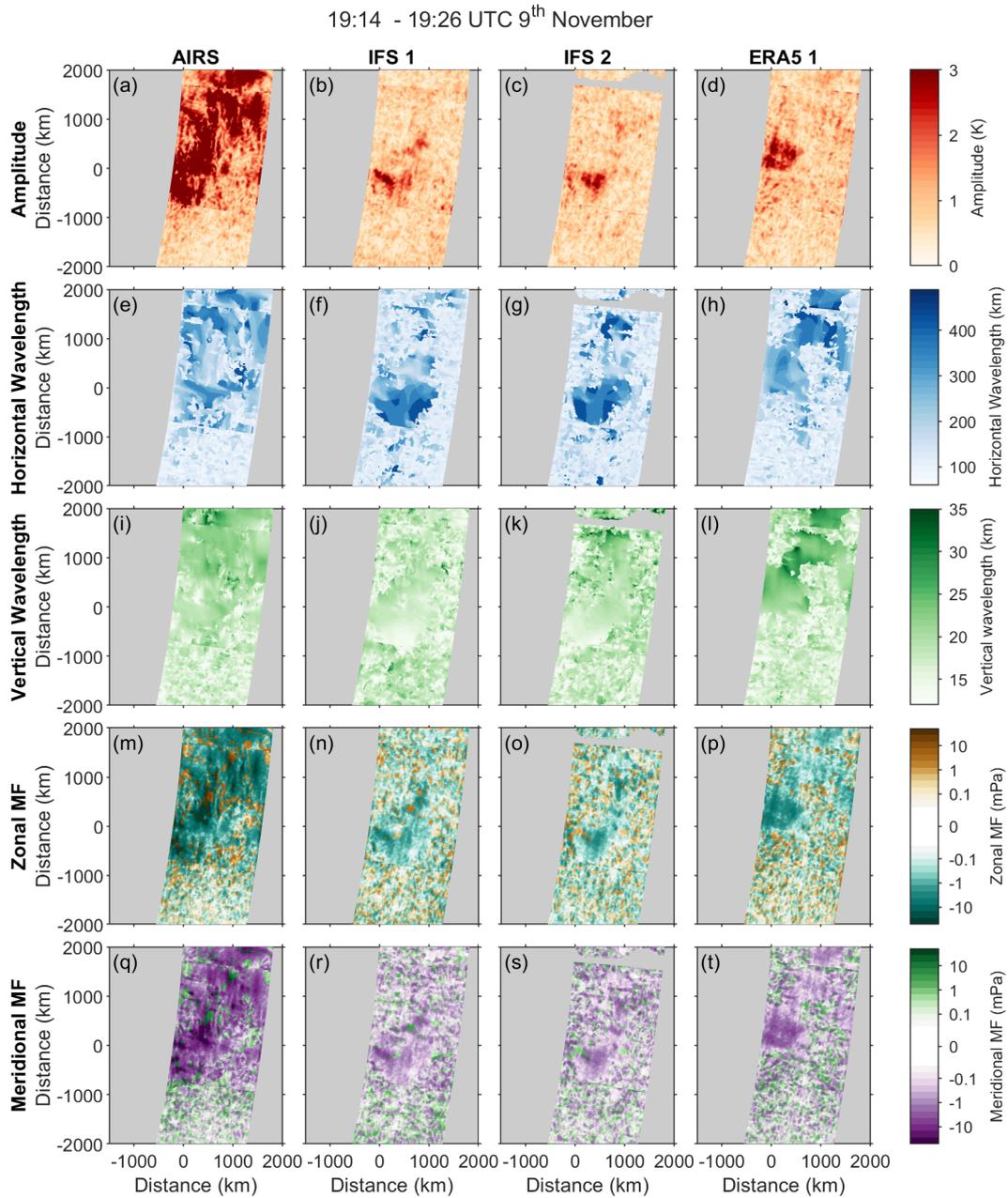


Figure 6. As in Figure 4, but for wave properties for the granules shown in Figure 5. The zonal and meridional momentum flux (MF) are shown on a log color scale.

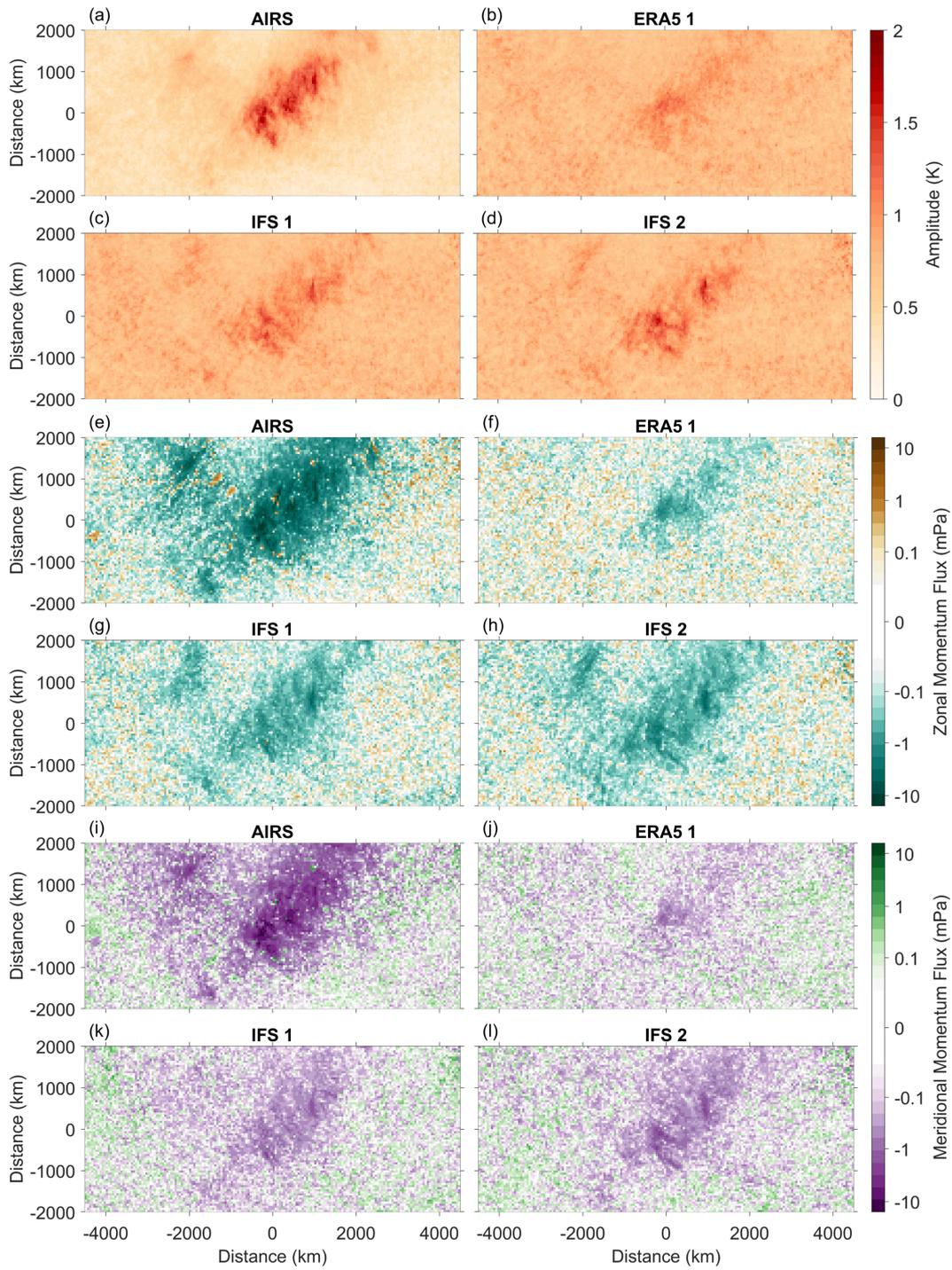


Figure 7. Mean amplitude (a–d) and mean zonal (e–h) and meridional (i–l) momentum flux in the region shown in Figure 1 for nighttime data in the first 10 days of November 2018. This is plotted on a regular distance grid with a point spacing of 50 km by 50 km. AIRS amplitudes in (a) are divided by 2. The zonal and momentum flux are shown on a log color scale.

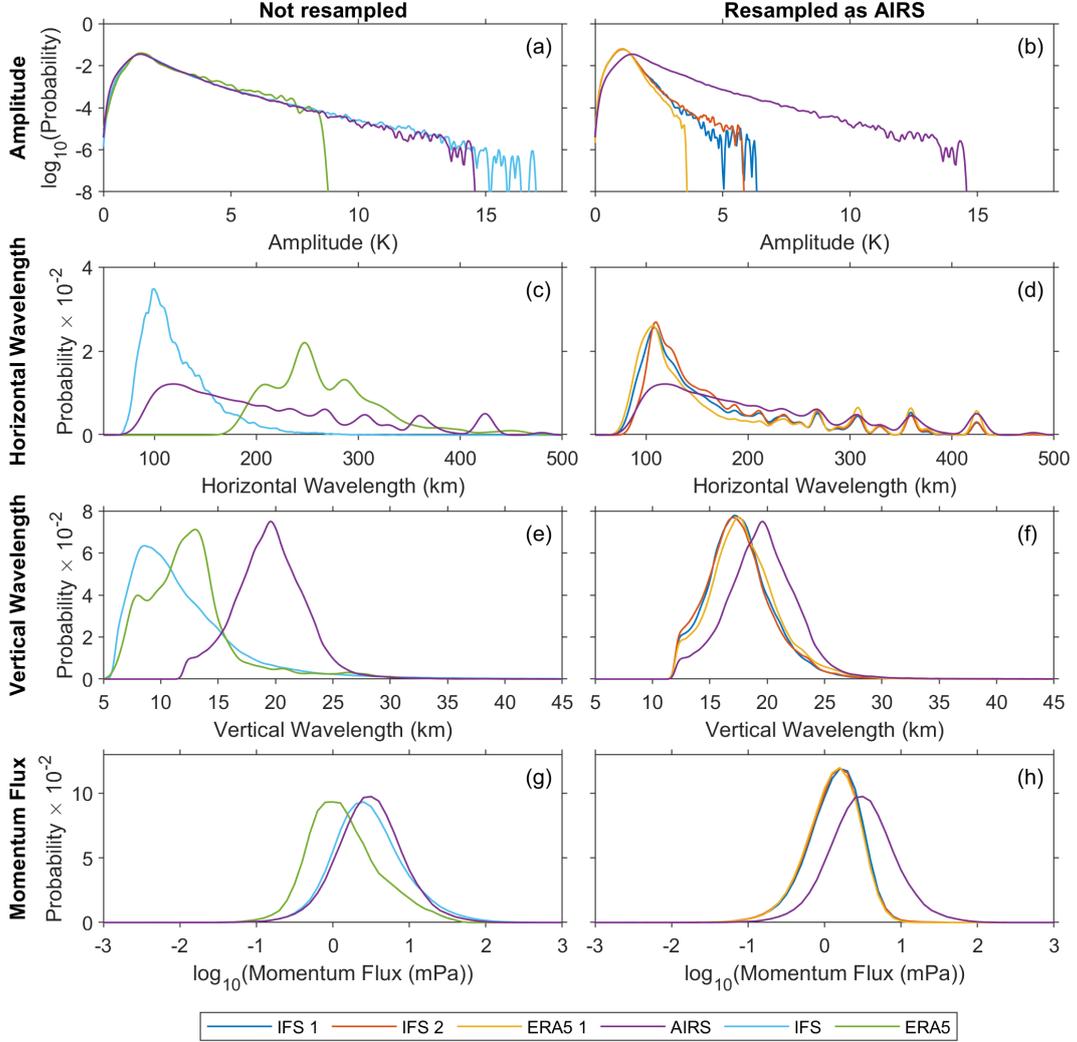


Figure 8. Kernel distribution functions (KDFs) for the wave properties in nighttime data at 39 km altitude for AIRS and the models before resampling (the 1 km IFS run (IFS) and ERA5) (panels a, c, e and g) and for the resampled 1 km IFS run (IFS 1 and 2) and ERA5 resampled as AIRS (ERA5 1) shown with the KDFs for AIRS in panels (b, d, f and h). The KDFs for the amplitude have been logged to base 10. Noise is reduced by using a 70th percentile amplitude cutoff for the resampled models and AIRS, and AIRS’s amplitude cutoff is also used for ERA5 and the 1 km IFS run before resampling.

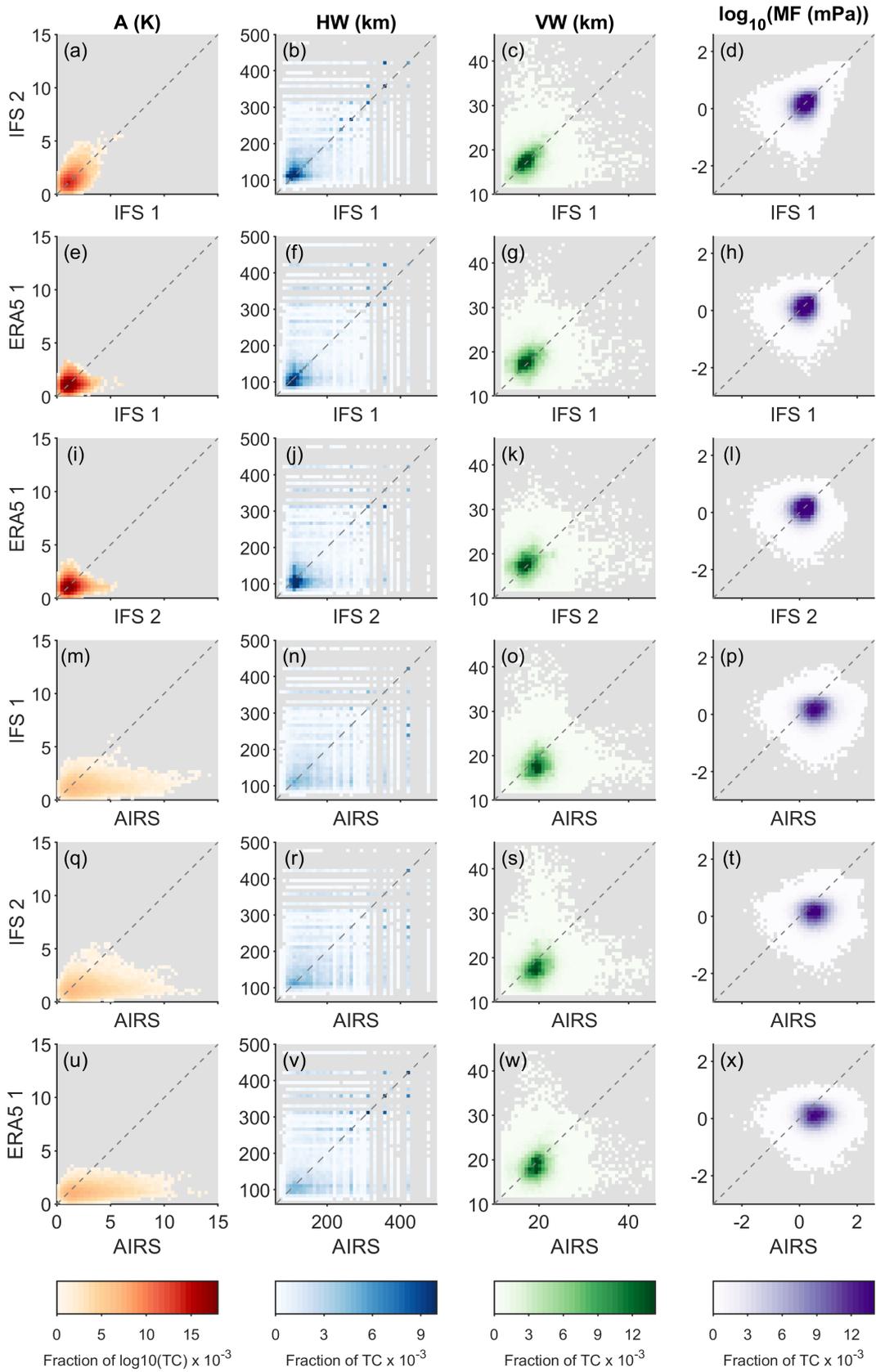


Figure 9. Bivariate histograms of wave properties (amplitude (A), horizontal wavelength (HW), vertical wavelength (VW) and momentum flux (MF)) using nighttime data from the first 10 days of November 2018 in the region shown in Figure 1 at 39 km altitude. The color scales show the fraction of the total bin counts (TC) for each subplot.