

1 **A New Aerosol Dry Deposition Model for Air Quality and Climate Modeling**

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16

17 **Key Points**

18 • New aerosol deposition velocity model agrees better with observations than

19 current models

20 • Impaction on microscale obstacles such as leaf hairs is key process

21 • New aerosol deposition velocity model increases dry deposition of PM_{2.5}

22 compared to the current CMAQ model

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28 **Abstract**

29 Dry deposition of aerosols from the atmosphere is an important but poorly understood
30 and inadequately modeled process in atmospheric systems for climate and air quality.
31 Comparisons of currently used aerosol dry deposition models to a compendia of
32 published field measurement studies in various landscapes show very poor agreement
33 over a wide range of particle sizes. In this study, we develop and test a new aerosol dry
34 deposition model that is a modification of the current model in the Community
35 Multiscale Air Quality (CMAQ) model. The new model agrees much better with
36 measured dry deposition velocities across particle sizes. The key innovation is the
37 addition of a second inertial impaction term for microscale obstacles such as leaf hairs,
38 microscale ridges, and needleleaf edge effects. The most significant effect of the new
39 model is to increase the mass dry deposition of the accumulation mode aerosols in
40 CMAQ. Accumulation mode mass dry deposition velocities increase by almost an order
41 of magnitude in forested areas with lesser increases for shorter vegetation. Peak $PM_{2.5}$
42 concentrations are reduced in some forested areas by up to 40% in CMAQ simulations.
43 Over the continuous United States, the new model reduced $PM_{2.5}$ by an average of 16%
44 for July 2018 at the Air Quality System monitoring sites. For summer 2018 simulations,
45 bias and error of $PM_{2.5}$ concentrations are significantly reduced, especially in forested
46 areas.

47 **Plain Language Summary**

48 Aerosol dry deposition is an important sink for atmospheric particles that are a health
49 hazard and a significant climate forcer. Uncertainties in modeling aerosol dry deposition
50 hamper accurate predictions of air quality and climate. A new aerosol dry deposition

51 model is developed that better agrees with observations of aerosol dry deposition velocity
52 for a variety of vegetation such as forests, grasslands, and water surfaces. This improved
53 aerosol dry deposition model when incorporated into air quality and climate models will
54 improve the accuracy of model predictions.

55

56 **1. Introduction**

57 The lifetime and fate of aerosols in the atmosphere are strongly influenced by wet and
58 dry deposition processes. Thus, the representation of these processes are key elements of
59 atmospheric models for air quality, climate, and ecosystem impacts. The uncertainties in
60 modeling aerosol dry deposition contribute significantly to the uncertainties and errors in
61 direct and indirect radiative forcing that have been identified as some of the most
62 uncertain processes in global climate modeling (IPCC, 2021). Currently, there are a wide
63 variety of aerosol dry deposition models used in atmospheric modeling systems that
64 reflect a great degree of uncertainty. Recently, there have been several studies that
65 compiled observations of aerosol dry deposition in a variety of environments for particle
66 sizes that range from 10s of nanometers to 10s of microns (Saylor et al., 2019; Emerson
67 et al., 2020; Farmer et al., 2021). Saylor et al. (2019) showed that models differ greatly
68 from the observations especially for forested landscapes. Model errors compared to
69 observations vary among different models and overpredict in some size ranges while
70 underpredicting in other size ranges.

71

72 Most size-resolved aerosol dry deposition models used in large-scale air quality and
73 climate models are combinations of mathematical algorithms representing the major

74 processes involved in aerosol deposition as presented by Slinn (1982). All these
75 processes have strong dependencies on particle size and their combination yields a
76 relationship with the dry deposition velocity as a function of particle diameter (i.e.,
77 $V_d(d_p)$). However, since the models in use today have been shown to not agree well with
78 consensus of observations, particularly for forests, these models need to be re-examined
79 and revised. Key questions include: Can the parameterizations of the major processes be
80 revised to improve results or are there key processes that have been neglected?

81

82 The focus of this paper is to address uncertainties in the current aerosol dry deposition
83 modeling and to propose a new model that builds on current forms but includes a key
84 new process that greatly improves agreement with the consensus of observations.

85 Section 2 describes physical processes controlling dry deposition modeling. The
86 proposed new model is described in Section 3. Evaluation of the new model against
87 measurements and discussion are presented in Section 4. Section 5 presents the
88 implementation and evaluation of the new model in the Community Multiscale Air
89 Quality (CMAQ) model (Byun and Schere, 2006) for regional applications. Concluding
90 remarks and future work are given in Section 6.

91

92 **2. Physical processes in modeling dry deposition**

93 The concept of dry deposition velocity (V_d) is that surface flux (F) of a trace atmospheric
94 constituent is directly proportional to its concentration (C) just above the surface as:

$$95 \quad F = V_d \times C \quad (1)$$

96

97 In this way physical and dynamical processes can be isolated from chemical processes.
98 The principal processes involved in aerosol dry deposition include gravitational settling,
99 Brownian diffusion, surface impaction, surface interception, and rebound. All these
100 processes are functions of particle diameter (d_p) and become more effective as particle
101 size increases except Brownian diffusion which is most effective for ultrafine particles.
102 Aerosol particles are transported from the air to the surface simultaneously by turbulent
103 fluxes and gravitational settling. The gravitational settling velocity (V_g) results from a
104 balance of gravitational and viscous drag forces as (Stokes, 1851):

$$105 \quad V_g = \frac{g\rho_p d_p^2 C_c}{18\mu} \quad (2)$$

106

107 where g is the gravitational acceleration, ρ_p is the particle density, μ is the dynamic
108 viscosity of air and C_c is the Cunningham slip correction factor for small particles
109 (Cunningham, 1910). The turbulent fluxes are modeled similarly to gas dry deposition
110 fluxes as combinations of resistances. Flux through the turbulent surface layer is
111 represented by aerodynamic resistance R_a which is the same for gases and aerosols
112 (Pleim and Ran, 2011). However, since the no-slip condition for viscous fluids requires
113 that the velocity is exactly zero at the boundary, there is a very thin quasi-laminar
114 sublayer adjacent to all surfaces. For gases, molecular diffusion across this layer is so
115 efficient that the resistance presented by this layer R_b is rarely the limiting factor relative
116 to R_a and surface resistances. For aerosol, however, R_b is usually the most limiting
117 resistance because diffusion of particles (Brownian diffusion) is much slower than
118 molecular diffusion. Gravitational settling and turbulent fluxes are combined to compute
119 aerosol deposition velocity as (Venkatram and Pleim, 1999),

120
$$V_d = \frac{V_g}{1 - \exp(-V_g(R_a + R_b))} \quad (3)$$

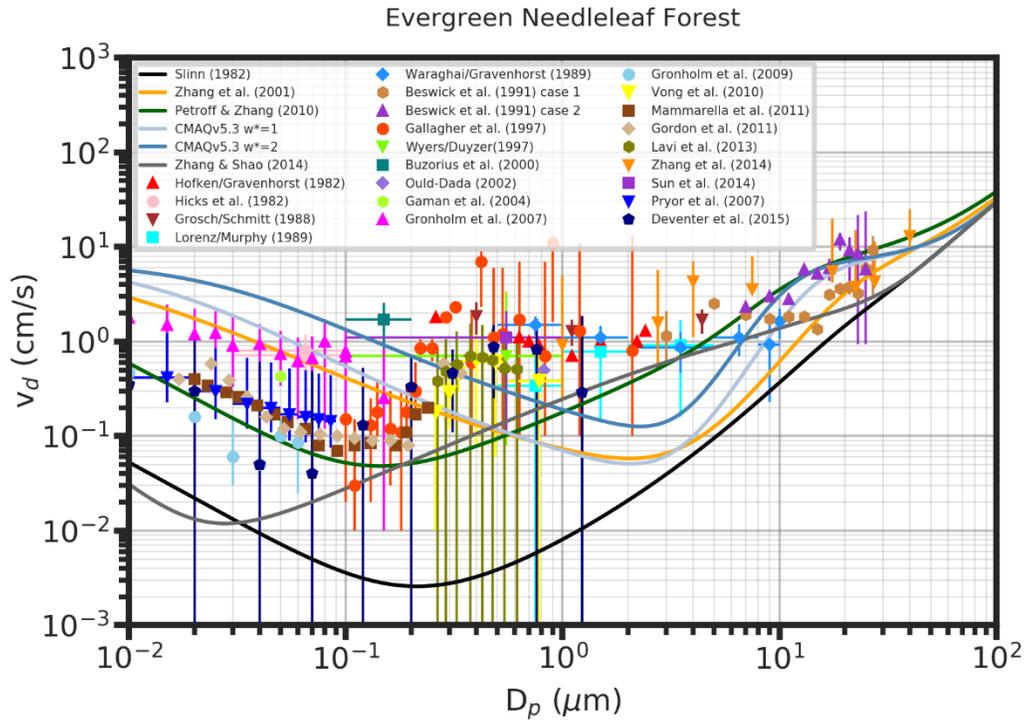
121 The quasi-laminar boundary layer resistance can be expressed in terms of collection
 122 efficiencies E (Slinn, 1982),

123
$$R_b = \frac{1}{LAI \cdot u_* (E_B + E_{im} + E_{in}) R} \quad (4)$$

124 where LAI is leaf area index, u_* is the friction velocity, E_B is Brownian diffusion
 125 collection efficiency, E_{im} is impaction collection efficiency, and E_{in} is interception
 126 collection efficiency. Particles that encounter a surface can either stick or bounce off
 127 which is represented by the rebound factor R in Equation 4 that mostly affects the largest
 128 particles. Most recent aerosol dry deposition models follow the formulations for the
 129 collection efficiencies proposed by Slinn (1982) with a variety of modifications and
 130 extensions.

131
 132 Aerosol deposition to vegetated areas is particularly complex and difficult to model given
 133 the wide array of vegetation types with different canopy morphologies and leaf shapes.
 134 Figure 1 shows that several models, including the models currently used in CMAQ, the
 135 Comprehensive Air quality Model with extensions (CAMx) (ENVIRON, 2020), the
 136 Goddard Earth Observing System model with Chemistry (GEOS-Chem) (Bey et al.,
 137 2001), a Unified Regional Air-quality Modeling System (AURAMS) (Gong et al., 2006),
 138 and the Global Environmental Multi-scale model - Modelling Air quality and CHEmistry
 139 (GEM-MACH) (Gong et al., 2015), cannot well represent aerosol deposition for all
 140 particle sizes at evergreen needleleaf forest sites. The current CMAQv5.3 model, which
 141 is a modified version of the earlier CMAQ model described by Pleim and Ran (2011),
 142 and the Zhang (2001) model that is used in CAMx, AURAMS, GEM-MACH and

143 GEOSChem both seem to have the wrong shape with minimum V_d at around 2-3 μm
144 while the measurements indicate minimum at about 0.1 – 0.2 μm . Note that there are two
145 curves for the CMAQv5.3 model because its formulation includes a function of
146 convective velocity scale (w^*). The models of Petroff and Zhang (2010) and Zhang and
147 Shao (2014) have minimum values at much smaller sizes with the Petroff and Zhang
148 (2010) model agreeing well with the aggregate observations in the less than 0.2 μm size
149 range. The recent model described by Emerson et al. (2020) also agrees well with the
150 size of minimum V_d for forest observations as will be shown and discussed in Section 4.
151 However, none of the models seem to capture the rapid increase in deposition velocity
152 seen in the observations from 0.2 to about 0.5 μm and the plateau from 0.5 to about 5 μm .
153 Clearly, the models as currently formulated are not able to produce the S-shaped curve of
154 the observation consensus.



155

156 Figure 1. Measured aerosol dry deposition velocities from literature as functions of
 157 particle size for evergreen needleleaf forest. Symbols represent median values with error
 158 bars that represent estimated uncertainty, usually inter-quartile range. The lines show
 159 predictions by various models assuming $u^* = 0.4$ m/s. Adapted from Saylor et al. (2019).

160

161 Several recent studies have shown that leaf surface texture and leaf shape have significant
 162 influence on aerosol dry deposition. For example, Chen et al. (2017) found that the
 163 needle-shaped leaves of conifers were more effective in general than broad leaves at
 164 $PM_{2.5}$ (particulate matter with an effective aerodynamic diameter less than $2.5 \mu m$)
 165 aerosol dry deposition. They also found that broad leaf species with more grooves or
 166 hairs tended to increase deposition. Several other studies involving both field

167 measurements and wind-tunnel experiments also showed increased deposition for leaves
168 with dense hairs, ridges, grooves or thick epicuticular wax layers (e.g., Weerakkody et al,
169 2017; Chiam et al., 2019; Leonard et al, 2016). Perini et al. (2017) also found enhanced
170 fine aerosol deposition on leaves with thick cuticular waxes but less so for hairy leaves.
171 By measuring aerosols accumulating on 22 species of trees and 25 shrubs, Sæbø et al.
172 (2012) found that leaf properties such as hair and wax cover enhanced aerosol deposition
173 among the broad-leaf species while needle-leaf species were also among the highest
174 aerosol collecting species. Beckett et al. (2000) found greater deposition of 1 μm aerosols
175 on pine needles than broad flat leaves in wind-tunnel studies. They noted that deposition
176 was well correlated with Stokes number which is inversely proportional to the
177 characteristic leaf size which for needles was on the order of 1 mm and about 5 cm or
178 more for broad leaves. Zhang et al. (2021) tested the effects of leaf hair (trichome)
179 density, leaf aspect ratio, petiole (leaf stem) length, and leaf fractal deviation on $\text{PM}_{2.5}$
180 deposition. They found higher trichome density, lower aspect ratio, shorter petiole, and
181 greater leaf fractal deviation all increase $\text{PM}_{2.5}$ dry deposition velocities.

182

183 **3. New model description**

184 The new model, which is intended to replace the current aerosol dry deposition scheme in
185 CMAQ, follows the same general framework that was originally proposed by Slinn
186 (1982) as shown in equations 3 and 4. The aerodynamic resistance is unchanged from the
187 current CMAQ model and is the same for gases and aerosols (Pleim and Ran, 2011).
188 Unlike the current scheme, calculation of the quasi-laminar boundary layer resistance R_b
189 differs for the vegetated and non-vegetated parts of each grid cell. The most important

190 change is to the parameterization of the impaction collection efficiency where a term is
 191 added to better represent the shape of the deposition velocity curve for vegetated areas.

192

193 **3.1 Vegetated areas**

194 For the vegetated fraction of each grid cell the R_b is weighted by LAI to account for the
 195 total leaf surface area density available for deposition as shown in Equation 4. The
 196 Brownian collection efficiency E_B follows Slinn (1982) as:

$$197 \quad E_B = \left(\frac{c_v}{c_d}\right) Sc^{-2/3} \quad (5)$$

198 where Sc is the Schmidt number defined as the ratio of the kinematic viscosity of air
 199 divided by the Brownian diffusivity $Sc = \nu/D_B$ and c_v/c_d is the ratio of viscous drag to
 200 total drag which we specify as 1/3 as deduced by Chamberlain (1966) for grass.

201 A key innovation in the new model to better fit observed dry deposition velocities by
 202 particle size is to represent the impaction efficiency by two terms to account for the
 203 effects of macroscale and microscale obstacles. Impaction efficiency E_{im} is generally
 204 expressed as a function of Stokes number St , which describes the tendency of a particle to
 205 follow fluid flow around obstacles. In the quasi-laminar sub-layer, the relevant flow
 206 velocity is given by the turbulent friction velocity u_* which is the characteristic velocity
 207 of turbulent eddies in the turbulent layer above the quasi-laminar sub-layer. Therefore,
 208 for vegetated surfaces, St is defined as,

$$209 \quad St = \frac{V_g u_*}{gA} \quad (6)$$

210 where A is the characteristic dimension of the obstacles. For the new model, we define
 211 St_l and St_h using different obstacle characteristic dimensions for the leaf scale A_l and

212 microscale A_h representing features such as leaf hairs (trichomes) or other microscale
 213 roughness on leaves. Thus, E_{im} is given as,

$$214 \quad E_{im} = (1 - f_{micro}) \frac{St_l^2}{1+St_l^2} + f_{micro} \frac{St_h^2}{1+St_h^2} \quad (7)$$

215 where f_{micro} is the fraction of total impaction due to the microscale features. The concept
 216 of using macro and microscale obstacle size scales was introduced by Slinn (1982) for
 217 interception processes. Slinn (1982) speculated that the microscale obstacles would
 218 probably not be relevant for impaction because the vegetative hairs or other microscale
 219 obstacles such as cobwebs would be deflected by the flow and not be significant
 220 collectors by impaction. However, testing the two-term approach for both interception
 221 and impaction showed that expressing impaction as in Equation 7 matched the size
 222 dependent deposition velocities, especially for forests, much better than using a similar
 223 expression for interception only. Note that both f_{micro} and A_h are very uncertain
 224 parameters. Slinn (1982) suggested $f_{micro} = 1\%$ but we found better fit to observations
 225 with a slight reduction to $f_{micro} = 0.8\%$. The microscale characteristic obstacle scale is
 226 specified by land use category (LUC) such that $A_h = 0.5 \mu\text{m}$ for needleleaf forest and
 227 grasslands and $A_h = 1.0 \mu\text{m}$ for deciduous forest. The macroscale characteristic obstacle
 228 scale is also specified by LUC with values ranging from 0.5 to 10 mm (Table 1).

229

230 The third collection efficiency in Equation 4 is interception. Interception is postulated as
 231 capture that occurs when a particle comes within a particle radius of a surface or obstacle.
 232 However, the physical basis of this process is less well defined than the Brownian or
 233 impaction processes. Including interception efficiency as recommended by Slinn (1982)

234 had very little effect in the new model. Therefore, in the new model the interception
 235 collection efficiency is not used.

236 Table 1. Key parameters for new aerosol dry deposition model over different landscapes

LU Type	U^* (m/s)	LAI	f_v	f_{micro}	A_l (mm)	A_h (μm)
Needleleaf forest	0.4	5	93	0.008	2	0.5
Broadleaf forest	0.4	5	93	0.008	10	1.0
Grassland	0.3	2	95	0.002	0.5	0.5
Water	0.2	0	0	NA	NA	NA

237

238 3.2 Non-vegetated areas

239 For non-vegetated areas the definition of E_B is the same as for vegetated areas but the E_{im}
 240 is different,

$$241 \quad E_{im} = 10^{-3/St} \quad (8)$$

242 where,

$$243 \quad St = \frac{\rho_a V_g u_*^2}{g\mu} \quad (9)$$

244 These are the formulations recommended by Slinn (1977) for smooth surfaces where ρ_a
 245 is the air density. For water surfaces, the effects of whitecaps (breaking waves) are
 246 included in an additional term in the E_B expression,

$$247 \quad E_B = (1 - f_{wc}) \left(\frac{C_v}{C_d} \right) Sc^{-2/3} + f_{wc} \frac{u_*}{U_{10}} \quad (10)$$

248 as suggested by Pryor (1999) following Hummelshøj et al. (1992) where U_{10} is the
 249 windspeed at 10 m above the surface. Note that the more complex form in Pryor (1999)
 250 was not used because the term describing the particle capture by spray droplets is always
 251 much smaller than the u^*/U_{10} term. The effects of whitecaps increase rapidly with
 252 increasing windspeed as whitecaps cover more of the water surface. While most models
 253 parameterize the whitecap surface fraction f_{wc} as a function of windspeed, we follow
 254 Albert et al. (2016) who developed a parameterization based on satellite whitecap
 255 fraction data, that is also a function of water surface temperature T_{ws} in Celsius and U_{10} is
 256 in m/s,

$$257 \quad f_{wc} = a(b + U_{10})^2 \quad (11)$$

$$258 \quad \text{where} \quad a = a_1 + a_2 T_{ws} + a_3 T_{ws}^2$$

$$259 \quad b = b_1 + b_2 T_{ws}$$

260 and $a_1 = 8.46 \times 10^{-5}$, $a_2 = 1.63 \times 10^{-6}$, $a_3 = -3.35 \times 10^{-8}$, $b_1 = 3.354$, and $b_2 = -0.062$.

261

262 Another new modification is for the non-vegetated parts of urban LU categories where R_b
 263 is weighted by building area index (BAI). The rationale for this is that buildings add
 264 significant surface area in urban landscapes that is not accounted for by LAI or otherwise.

265 An initial estimate is $BAI = (4 \lambda_f + 1)/(1-f_v)$, where λ_f is frontal area density of building
 266 surfaces and f_v is the vegetated fraction of the grid cell. The logic for this is that the total
 267 building surface area is four times the frontal area. For applications where detailed
 268 building data is not included, default values of BAI are specified by landuse category.

269 For example, BAI is specified as 1.8, 2.0, and 2.3 for the National Land Cover Database
 270 (NLCD) categories of low intensity developed, medium intensity developed, and high

271 intensity developed, respectfully. For all other non-developed categories, BAI is set to 1.
 272 Thus, for unvegetated portions of urban grid cells R_b is defined as,

$$273 \quad R_b = \frac{1}{BAI \cdot u_* (E_b + E_{im})} \quad (12)$$

274 Note that particle rebound effects are not considered ($R=1$) because only the largest
 275 particles are affected and there is much variation by composition, RH, surface type, and
 276 windspeed. Grid cell deposition velocity is the combination of the vegetated and non-
 277 vegetated parts,

$$278 \quad V_d = f_v V_{dveg} + (1 - f_v) V_{dnoveg} \quad (13)$$

279
 280

281 **4. New model evaluation**

282 The new model is evaluated against aerosol dry deposition measurements over different
 283 landscapes. The performance of the new approach is also discussed in terms of the
 284 enhanced processes.

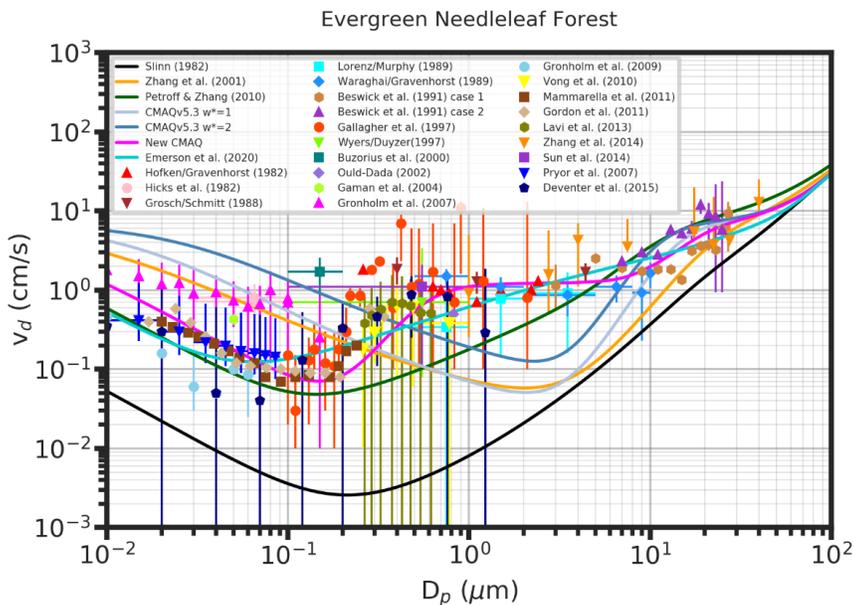
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286 **4.1 Comparison to measurements**

287 Following several recent papers (Saylor et al., 2019; Emerson et al., 2020; Farmer et al.,
 288 2021), we compile aerosol dry deposition measurements by particle size for various
 289 surface types from the published literature to compare to aerosol dry deposition models.
 290 Figure 2 shows the same data for evergreen needleleaf forest as figure 1 with the addition
 291 of the new model described here as well as the recent model described by Emerson et al.
 292 (2020). The parameter values for the model runs for each land use type are shown in
 293 Table 1. The new model, designated *New CMAQ*, is the only one that shows the increase

294 of V_d in the 0.2 – 0.6 μm range and the plateau from about 0.6 μm to 6 μm . Compared to
 295 the CMAQv5.3 model, the new model substantially reduces the underprediction of V_d
 296 from about 0.5 – 5.0 μm . This has profound effects on dry deposition and concentration
 297 of $\text{PM}_{2.5}$ as will be shown in Section 5. The Emerson model generally fits the data very
 298 well since it was developed through changes to six empirical coefficients and exponents
 299 in the Zhang et al (2001) parameterization to optimize agreement with the needleleaf
 300 forest observation data. However, as shown in Figures 2 and 3, the Emerson model does
 301 not replicate the s-shaped curve suggested by the observations for needleleaf and
 302 broadleaf forest through the 0.1 – 5.0 μm range as well as the new model.

303
 304



305
 306 Figure 2. Measured aerosol dry deposition velocities from literature as functions of
 307 particle size for evergreen needleleaf forest. Symbols represent median values with error

308 bars that represent estimated uncertainty, usually inter-quartile range. The lines show
309 predictions by various models assuming $u_* = 0.4$ m/s including the new model (new
310 CMAQ) in magenta and the recent Emerson et al. (2020) model in turquoise. .

311

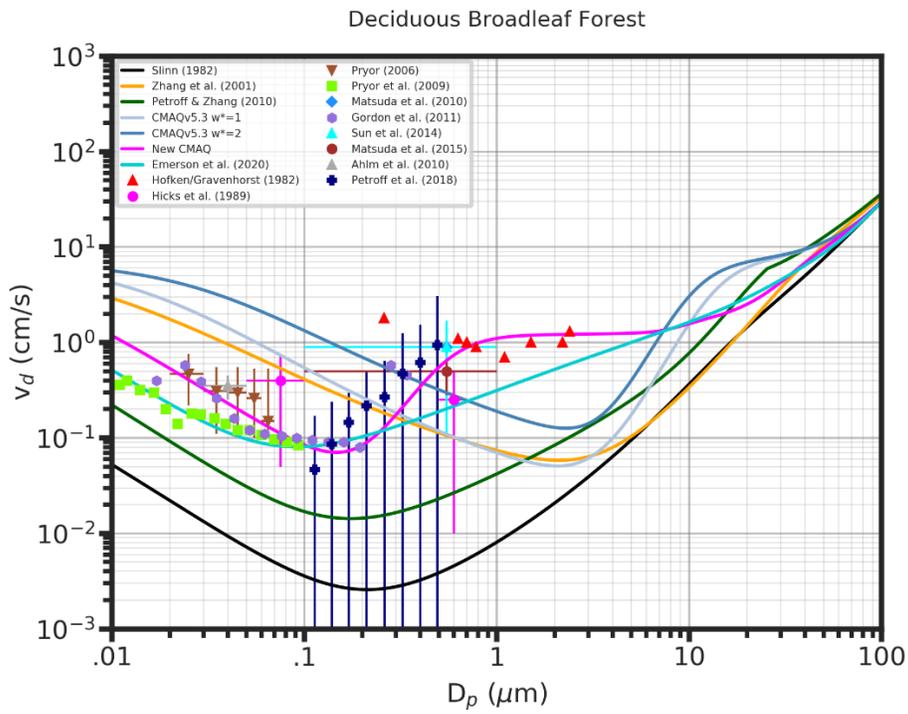
312 There are not as many measurement studies in the literature for other landuse/vegetation
313 types. Figure 3 shows the observation data and models for deciduous broadleaf forest.

314 The consensus of dry deposition measurements to broadleaf forests suggests a similar
315 shape to the V_d vs d_p curve where there is an increase of V_d in the approximately 0.2 to
316 0.6 μm range with a plateau up to about 6 μm . Again, the new model is the only one that
317 replicates the shape of this curve and does not greatly underestimate V_d in the 0.5 - 4.0
318 μm range.

319

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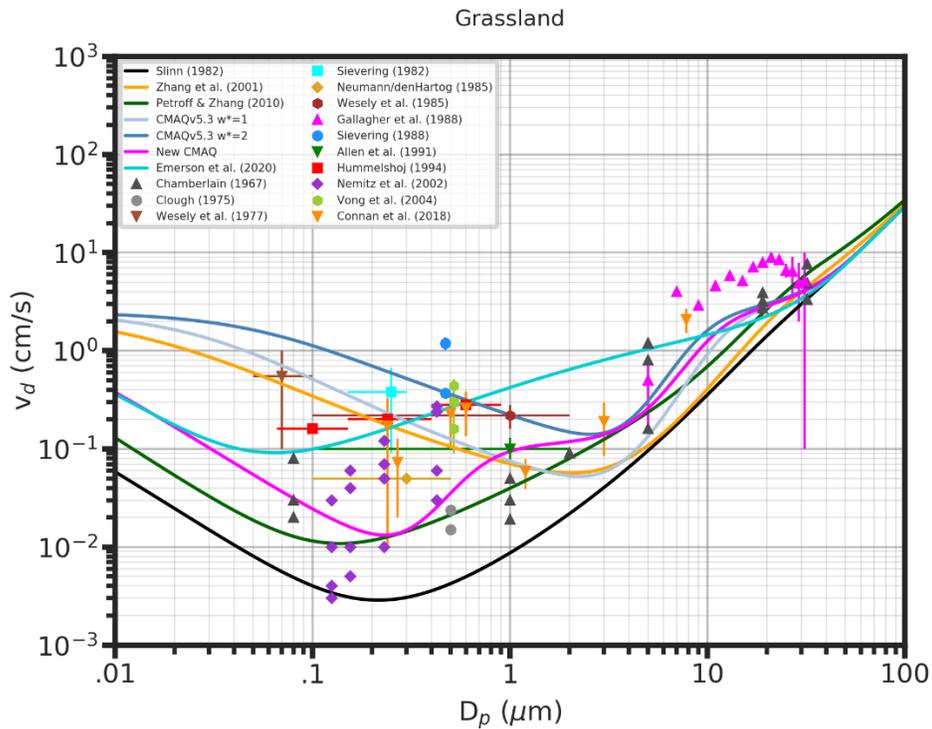
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322

323 Figure 3. Same as figure 2 but for deciduous broadleaf forest.

324



325

326

327

328 Figure 4. Same as figure 2 but for grasslands and $u_* = 0.3$ m/s.

329

330 Figure 4 shows the new model and other models for grasslands compared to

331 measurements. While there are a lot of measurements for grasslands, there seems to be

332 much less consensus among them, even within the same studies, than for forests. With

333 the degree of scatter in the measurements there isn't clear guidance for parameter

334 selection. While the rationale for microscale impaction may also apply to grass since

335 grass leaves often have leaf hairs or trichomes and serrated edges, the evidence from the

336 measurements is less clear. Therefore, the parameter values selected for the microscale

337 impaction scaling factor f_{micro} are set to smaller values (see Table 1) than for forests,

338 which seem to better fit with the measurements. Running box models on an hourly or

339 sub-hourly basis using detailed field measurements may add some clarity to model

340 performance and refinement of parameters.

341

342 The large scatter among the reported measurements for grassland is likely due the variety

343 of grass species which can have very different characteristics including length. For

344 example, the measurements reported by Vong et al. (2004) were made over rye grass 0.75

345 - 1 m tall, Allen et al. (1991) measured deposition to short grass of 3 – 7 cm in length,

346 Conan et al. (2018) used artificial grass, and Nemitz et al. (2002) measured in a moorland

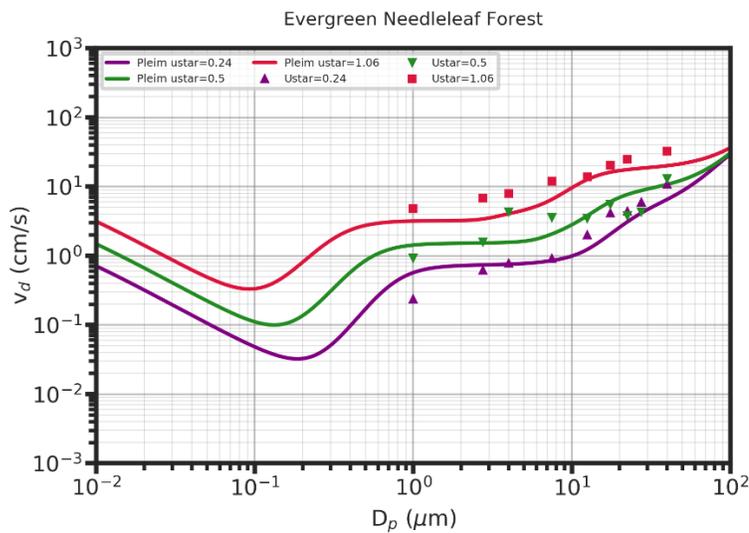
347 which is characterized by hummocks and hollows with vegetation including peat moss

348 (*Sphagnum*) and several species of tall grasses.

349

350 Virtually all studies found strong dependences of V_d on u^* with V_d increasing with
 351 increasing u^* (Pryor et al., 2008). Some have suggested that V_d/u^* is a more robust
 352 quantity for analysis and comparison (e.g., Conan et al., 2018) but many studies did not
 353 report this. Zhang et al. (2014) measured dry deposition velocities of dust particles (1 –
 354 40 μm) in a wind tunnel at various wind velocities over different surfaces. To simulate
 355 deposition to trees, evergreen branches were planted in the test section surface. The new
 356 model set up for evergreen needleleaf forest is shown (Figure 5) to compare well to the
 357 tree wind tunnel experiments at three windspeeds with measured friction velocities of
 358 0.24, 0.5, and 1.06 m/s.

359



360

361 Figure 5. New model compared to wind tunnel experiments for tree surfaces (Zhang et
 362 al. 2014) at three friction velocities.

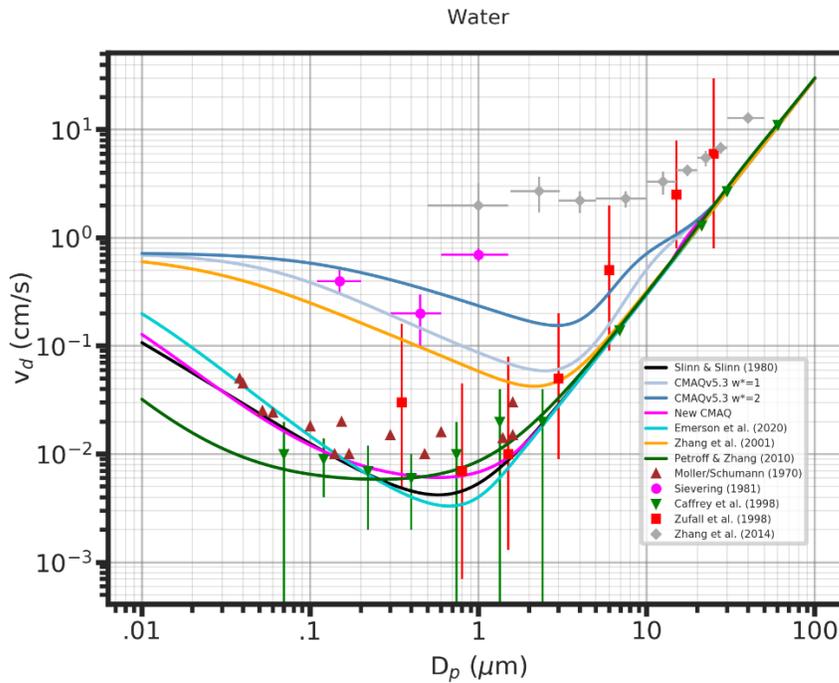
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364 Figure 6 shows models compared to field measurements for water surfaces. Most of the
365 measurements over water show that dry deposition velocities are much lower in the
366 accumulation size range than for vegetated surfaces. The measurements by Sievering
367 (1981) seem to be outliers with much higher deposition velocities in the $0.15 - 1 \mu\text{m}$
368 range. These measurements, which used the momentum gradient technique that assumes
369 similarity between aerosol fluxes and momentum fluxes, can be particularly uncertain
370 when the surface has bluff bodies such as waves than can induce form drag on
371 momentum flux. The other outlier study is by Zhang et al. (2014) which is a wind tunnel
372 study where dry deposition velocity is estimated by particle dynamic analysis (PDA).
373 The measurements over water in the wind tunnel were found to agree quite well with the
374 Slinn and Slinn (1980) (SS80) model for three different friction velocities when R_b is set
375 to zero (Zhang et al., 2014, Figure 11). The authors hypothesize that this is due to waves
376 and spray droplets. However, given the dramatic dissimilarity from most of the other
377 studies, there could be an issue of scaling wave-wind dynamics to a wind tunnel where
378 the water is very shallow with restricted fetch.

379

380 The new model agrees well with the measurements other than the 2 outlier studies and is
381 most similar to the SS80 model. The main difference in formulation between the new
382 model and the SS80 model for water is the inclusion of the effects of whitecaps in
383 enhancing deposition velocity. Since Equation 11 does not have any dependence on d_p
384 the whitecaps effects are effectively a lower limit on V_d which raises the trough of the
385 curve.

386



387

388 Figure 6. Same as Figure 2 but for water surfaces.

389

390 4.2 Discussion

391 Figure 7 shows the size dependence of the components of the dry deposition model for

392 needleleaf forest. Clearly, the most effective component responsible for the S-shape

393 curve for vegetated surfaces in the 0.2 – 10 μm range is the impaction on microscale

394 features. The impaction collection efficiency acts as a threshold process where $E = 1$

395 where $St > \sim 3$ for the formulation shown in Equation 7 and ramps down as St and d_p get

396 smaller. The Stokes number, as represented by Equation 6, is a ratio of the inertial

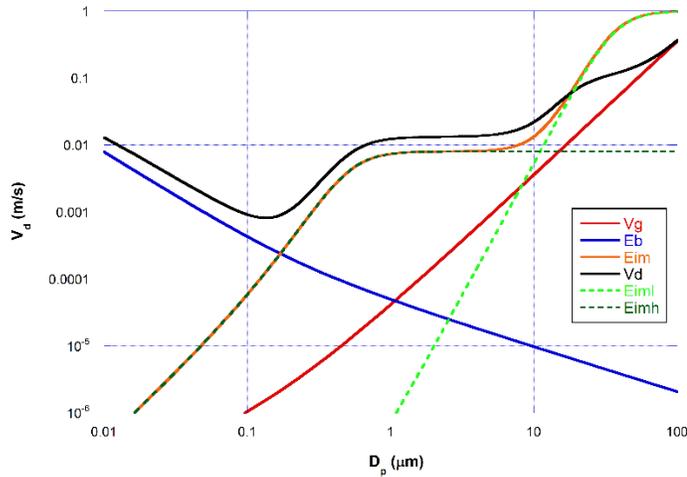
397 stopping distance of a particle to a characteristic length scale of an obstacle. For the

398 macroscale term, which is common to nearly all current aerosol deposition models, the

399 characteristic length represents the effects of the leaves. Figure 7 shows that the leaf

400 scale impaction term has greatest effect on particles larger than about 5 μm . Note that for

401 broadleaf forests the effects of this term are limited to even larger particles because the
402 characteristic length scale is greater for broad leaves than for needle leaves.
403
404 The stopping distance for quasi-laminar sublayers on leaves can be estimated as $V_g u^*/g$.
405 For $u^* = 0.4$ m/s the stopping distance decreases with decreasing d_p reaching $1 \mu\text{m}$ at $d_p =$
406 $0.75 \mu\text{m}$. A characteristic length of about $0.5 \mu\text{m}$ and a scaling factor f_{micro} of about
407 0.008 result in a good fit to the measured data as shown in Figure 2. A physical
408 explanation of this process is that there exist microscale features on many leaves and
409 stems that act as obstacles in the quasi-laminar sublayers. Many studies have shown
410 increased deposition of $\text{PM}_{2.5}$ for broadleaf species that have dense hairs, ridges, grooves,
411 or thick epicuticular wax layers as summarized above in Section 2. While needleleaf
412 species generally don't have leaf hairs, they are often seen to be particularly efficient at
413 $\text{PM}_{2.5}$ collection. A possible explanation is that the quasi-laminar sublayer grows from
414 the leading edge of a surface and therefore will be thinner near the edge. Since the needle
415 shape presents far more edge to the flow per area than broad leaves the deposition of sub-
416 micron sized particles is more efficient. The hypothesis that more edge per area increases
417 deposition is supported by the results of Zhang et al (2021) that showed increased $\text{PM}_{2.5}$
418 deposition to leaves with lower aspect ratio and greater fractal deviation. A physical
419 interpretation of the scaling factor is that the edge effects of needle shaped leaves is only
420 affecting a small portion of the leaf area. For broadleaf forest, the scaling factor accounts
421 for the fraction of species that have dense leaf hairs or other microscale features and the
422 sparsity of these features per leaf area.



423

424 Figure 7. Size dependence of the components of the dry deposition model for needleleaf
425 forest.

426

427 5. CMAQ implementation and evaluation

428 The new aerosol dry deposition model is implemented in the latest version of the CMAQ
429 modeling system. The coupled version of the Weather Research and Forecasting
430 (WRFv4.0.2) model and the CMAQv5.3 model are used to evaluate the new approach in
431 air quality simulations over different-resolution domains for summer conditions in 2018.

432

433 5.1 Model implementation

434 Since CMAQ uses log-normal size distributions to represent Aitken, accumulation, and
435 coarse modes, aerosol deposition velocities need to be integrated over the log-normal size
436 distributions to calculate the 0th, 2nd, and 3rd moments which represent the number,
437 surface area, and volume of each mode, respectively. Therefore, the terms in the model

438 that have explicit dependence on particle diameter D are integrated following Binkowski
 439 and Shankar (1995)

$$440 \quad \hat{X}_k = \frac{1}{M_k} \int_{-\infty}^{\infty} X D^k (\ln D) d \ln D \quad \text{where } M_k = N D_g^k \exp\left(\frac{k^2}{2} \ln^2 \sigma_g\right) \quad (14)$$

441 where k is the moment (0,2,3) and N is the particle number concentration. The only
 442 terms in the new model that are explicit functions of particle diameter are Brownian
 443 diffusivity and gravitational settling velocity. For Brownian diffusivity the integrated
 444 form is,

$$445 \quad \hat{D}_{Bk} = D_B(D_g) \left\{ \exp\left(\frac{(-2k+1)}{2} \ln^2 \sigma_g\right) + 1.246 K n_g \exp\left(\frac{(-4k+4)}{2} \ln^2 \sigma_g\right) \right\} \quad (15)$$

446 For gravitational settling velocity the integrated form is:

$$447 \quad \hat{V}_{gk} = V_g(D_g) \left\{ \exp\left(\frac{(4k+4)}{2} \ln^2 \sigma_g\right) + 1.246 K n_g \exp\left(\frac{(2k+1)}{2} \ln^2 \sigma_g\right) \right\} \quad (16)$$

448 where D_g is the geometric mean diameter, σ_g is the geometric standard deviation, and the
 449 Knudsen number is $K n_g = 2\lambda/D_g$ where λ is the mean free path. For the modal model, the
 450 dry deposition velocity is computed as described in Section 3 but with \hat{D}_{Bk} and \hat{V}_{gk} for k
 451 $= 0,2,3$ for the three moments of each of the three modes replacing D_B and V_g .

452

453 Figure 8 demonstrates the relationships among the dry deposition velocity moments for
 454 both the CMAQv5.3 model and the new model plotted against geometric mean diameter
 455 D_g compared to the non-integrated models vs particle diameter D_p applied for needleleaf
 456 forest. The 0th moment, which represents the number of the modal distribution, shows
 457 the effects of integration are to increase the \hat{V}_d over the V_d for all sizes except in the
 458 plateau range ($\sim 1 - 3 \mu\text{m}$) of the new model. The 3rd moment is similar to the 0th moment

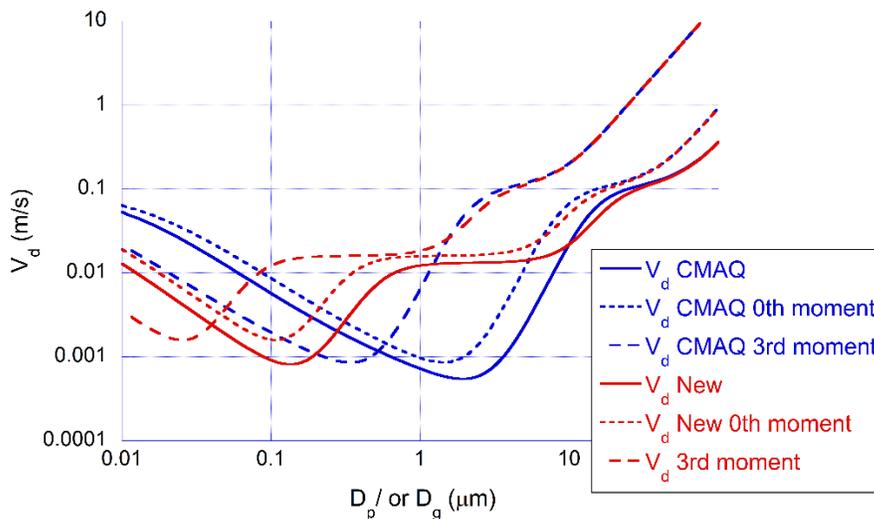
459 but with a shift to smaller D_g because the larger end of the distribution contributes more
 460 to volume than number.

461 Since the 3rd moment is proportional to mass, the 3rd moment of the dry deposition
 462 velocity represents the dry deposition velocity for mass concentration of aerosols. From
 463 Figure 8 it is evident that the mass deposition velocity for the new model is about an
 464 order of magnitude greater in the accumulation mode ($D_g \sim 0.1 - 0.4 \mu\text{m}$) than for the
 465 CMAQv5.3 model in forested areas. The effects of this increased dry deposition are
 466 assessed for CMAQ simulations across the conterminous US (CONUS) and at high
 467 resolution in the NE U.S.

468

469

470



471

472 Figure 8. The dry deposition velocities for the 0th and 3rd moments of the CMAQv5.3

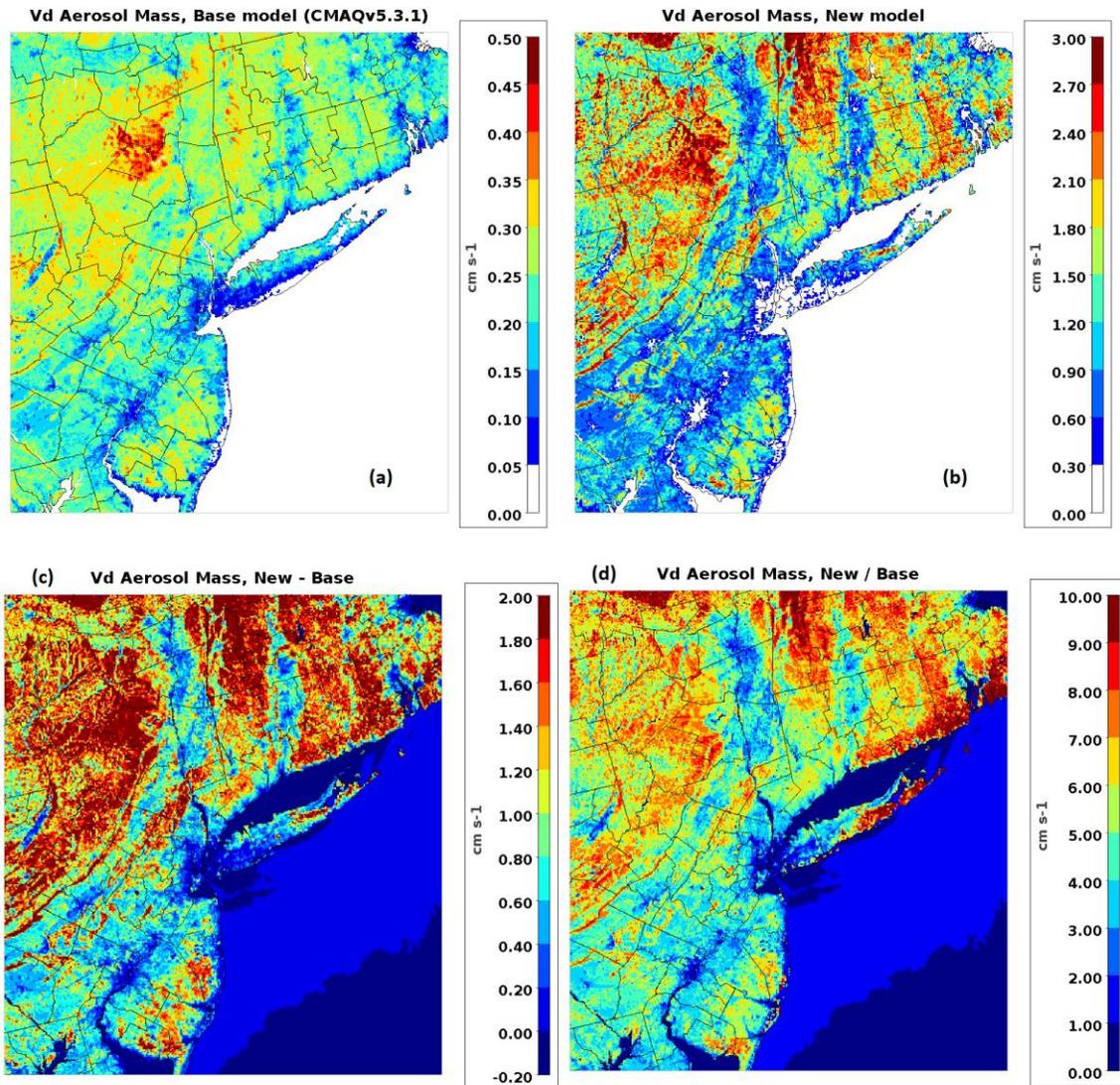
473 model and the new model compared to the non-integrated dry deposition velocities for

474 both models applied to needleleaf forest. The x-axis represents D_g for the \hat{V}_d plots and D_p
475 for the V_d plots.

476

477 **5.2 WRF-CMAQ simulations**

478 The coupled WRFv4.0.2/CMAQv5.3.1 model system was run in both the base
479 configuration and with the new aerosol dry deposition model for several months in
480 summer 2018. These simulations were based on modeling of the Long Island Sound
481 Tropospheric Ozone Study (LISTOS) which was an intensive multi-institutional field
482 study during the summer of 2018 (Karambelas 2020). The base model was run for three
483 resolutions (12 km, 4 km, 1.33 km) where the outer domain covered the CONUS with
484 one-way nested domains over the northeast (NE) states (4 km) and the New York/New
485 Jersey/Connecticut (NY-NJ-CT) region (1.33 km). Detailed description of the model
486 configuration and evaluation are presented by Torres-Vazquez et al. (2022). Model
487 simulations using the new aerosol dry deposition model (NEW) were conducted for the
488 12 km CONUS domain for July 2018 and for the 1.33 km domain for July and August
489 2018. In both cases, identical simulations using the base CMAQv5.3 model (BASE)
490 were also run. Initial conditions for the 12 km CONUS runs were from base case runs on
491 June 21, 2018, that were started on January 1, 2018. The 1.33 km runs were initialized
492 on July 1, 2018, from base case 1.33 km runs that started on May 2, 2018. All runs used
493 the same boundary conditions and emissions as described by Torres-Vazquez et al.
494 (2022).



495

496

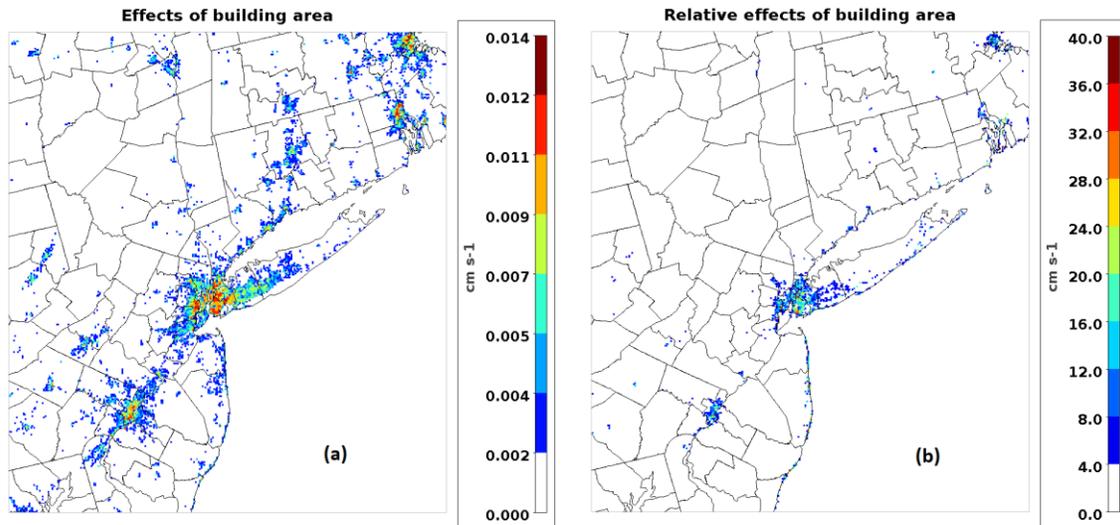
497 Figure 9. Accumulation mode mass dry deposition velocity from WRF-CMAQ on July
 498 10, 2018 at 18 UTC (2 pm LT) for (a) BASE, (b) NEW, (c) NEW-BASE and (d)
 499 NEW/BASE over the New York/New Jersey/Connecticut 1.33-km domain. Note that the
 500 scales for (a) and (b) are different. b)

501

502 Figure 9 shows the dry deposition velocity for accumulation mode mass on July 10, 2018,
 503 at 18 UTC (2 pm LT) for BASE, NEW, NEW-BASE and NEW/BASE over the NY-NJ-

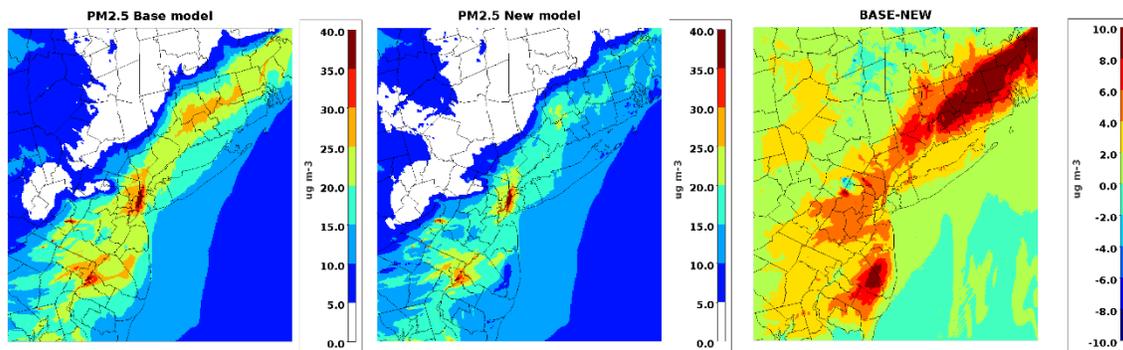
504 CT 1.33 km resolution domain. July 10 was a particularly polluted day in the NYC area.
505 Thus, it is an interesting case to study the effects of the new aerosol dry deposition
506 model. The plots in Figure 9a and 9b both show how variations in land use strongly
507 influence aerosol dry deposition for both BASE and NEW with greatest V_d in the forested
508 areas. This is due to the combinations of large roughness length resulting in low
509 aerodynamic resistance and large LAI. The greatest effects of the new model are
510 indicated by the difference in the plot scales (BASE plots 0-0.5 cm/s; NEW plots 0-3.0
511 cm/s) which reflects an almost order of magnitude increase in V_d for accumulation mode
512 mass in forested areas. Figures 9c and 9d, NEW – BASE and NEW/BASE, respectively,
513 demonstrate that the largest differences are in the forested areas of Pennsylvania (PA),
514 NY, Massachusetts (MA), CT, and Rhode Island (RI). The ratio of NEW/BASE (Figure
515 9d) is a factor of 8-10 in the heavily forested areas while it is only 1-3 in developed areas
516 depending on the intensity of development. The inclusion of the building area as in
517 Equation 12 increases V_d for accumulation mode mass in developed areas but only by a
518 small amount as shown in Figure 10a. However, because the V_d in developed areas is
519 small compared to vegetated areas, inclusion of the building effects has a substantial
520 relative impact on the V_d (Figure 10b).

521



522

523 Figure 10. Difference (a) and relative difference (b) in V_d for accumulation mode mass
 524 between a NEW model run using BAI as in Equation 12 and a NEW mode run where BAI
 525 = 1 for July 10, 2018, at 18 UTC over the New York/New Jersey/Connecticut 1.33-km
 526 domain.



527

528 Figure 11. $\text{PM}_{2.5}$ concentration on July 11, 2018, at 00 UTC (July 10, 8 pm LT) for
 529 BASE model (a) NEW model (b), and BASE - NEW (c) over the 1.33 km resolution
 530 domain.

531

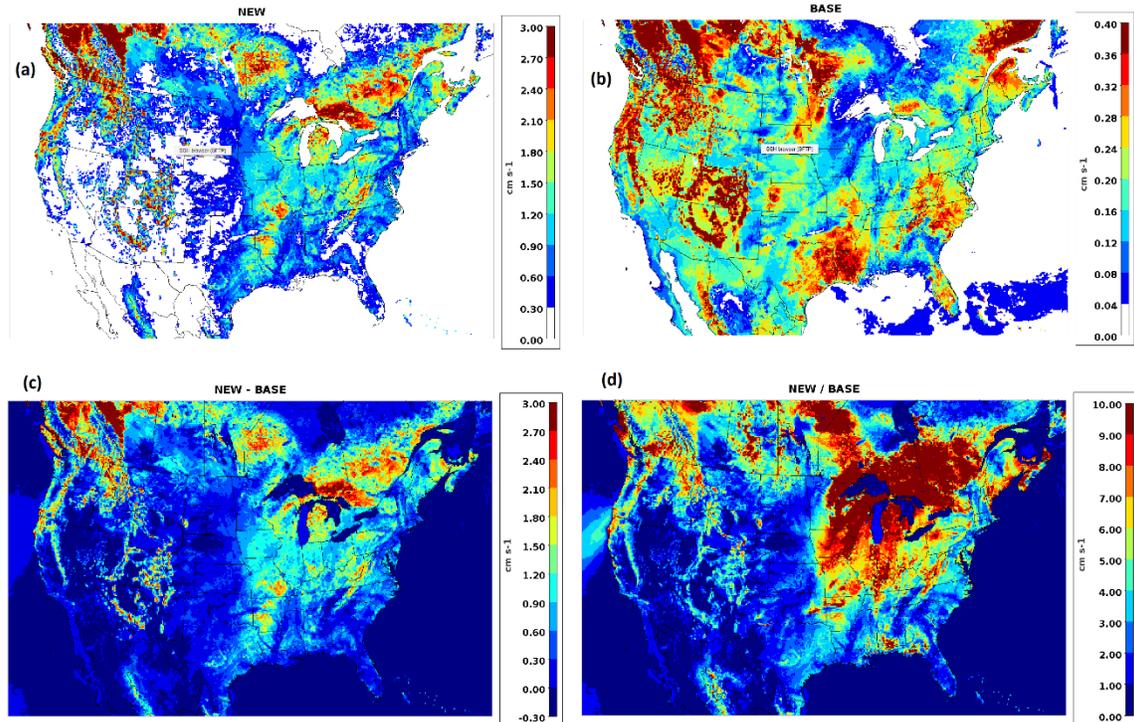
532 Figure 11 shows the consequences of the new aerosol deposition velocities, shown in
 533 Figure 9, on $\text{PM}_{2.5}$ concentrations 6 hours later when the highest concentrations occurred

534 in the region. The $PM_{2.5}$ concentrations at this time are greatest in the urban centers of
535 NY city (NYC) and Philadelphia with high concentrations downwind to the northeast into
536 CT and RI. The biggest effect of the new aerosol dry deposition is to substantially reduce
537 the downwind concentrations in southern New England which is mostly forested.
538 Another relative concentration maximum in the BASE model in southeastern NJ in an
539 area known as the Pine Barrens, which is dense forest with a high fraction of needleleaf
540 trees, is mostly absent in the NEW model run. Thus, the NEW model has larger effects
541 on $PM_{2.5}$ concentrations in forested areas than other areas such as in urban areas. The
542 peak concentrations in NYC only reduced by 10% while in some of the forested areas the
543 NEW simulation reduces $PM_{2.5}$ concentration by more than 40%.

544

545 The afternoon (18 – 20 UTC) deposition velocity for accumulation mode mass was
546 averaged for July 2018 for BASE and NEW model runs for the CONUS at 12 km
547 resolution (Figure 12). Again, it is evident that the largest effects of the NEW model are
548 in the forested areas mostly in the NE, western mountains and across the boreal forests of
549 Canada. In some areas of the northeastern U.S. and southeastern Canada the NEW dry
550 deposition velocities are 7-10 times the BASE values (Figure 12d). Note that Figure 12a
551 shows some discontinuities at the Canadian border. This is due to the hybrid land use
552 data that is a combination of higher resolution NLCD for the CONUS and the lower
553 resolution Moderate Resolution Imaging Spectroradiometer (MODIS) land use data for
554 elsewhere (Torres-Vazquez et al., 2022; Appel et al. 2014). In some of the sparsely
555 vegetated areas in the west the difference between NEW and BASE is quite small while
556 in the plains and predominately agricultural areas the difference is moderate.

557



558

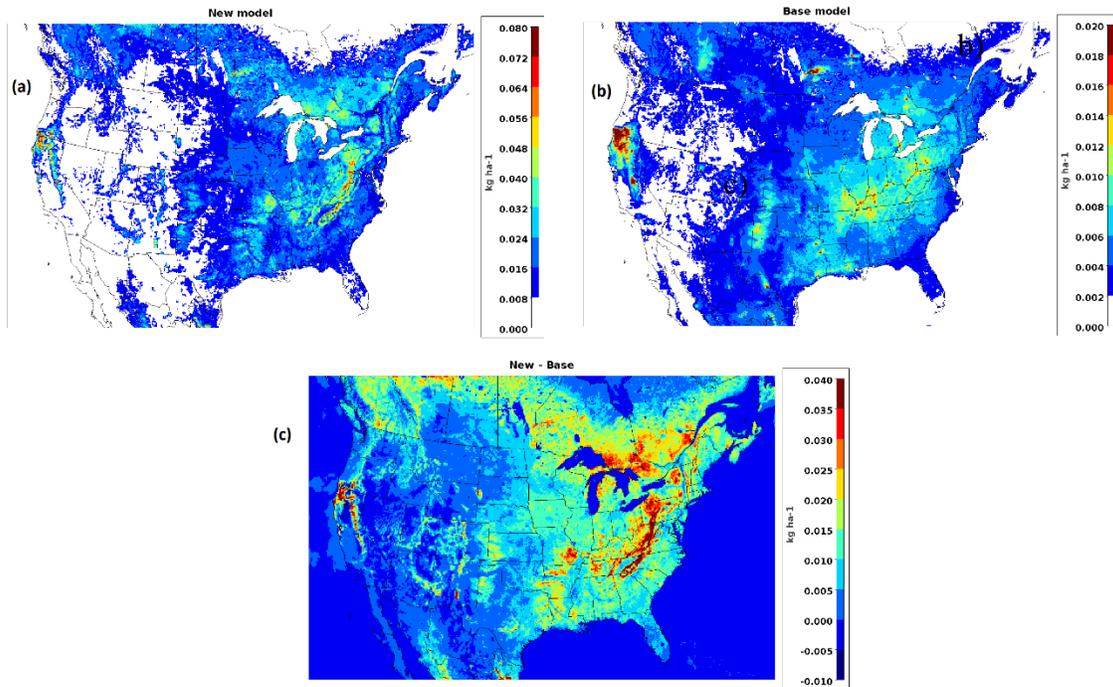
559 Figure 12. Afternoon (18 – 20 UTC) average accumulation mode mass dry deposition
 560 velocity for July 2018 on the CONUS 12 km domain for (a) BASE, (b) NEW, (c) NEW-
 561 BASE and (d) NEW/BASE. Note that the scales for (a) and (b) are different.

562

563 Increases in dry deposition velocity for accumulation mode aerosol using the NEW
 564 model increase the loading of aerosol species to land ecosystems. For example, Figure
 565 13 shows that accumulated dry deposition mass of accumulation plus Aitken mode
 566 (approximately $< 2.5 \mu\text{m}$) ammonium aerosol for July 2018 is much greater for the NEW
 567 model than the BASE model. Thus, the NEW model has much greater predictions of
 568 nutrient loading of the aerosol components, especially to forested ecosystems. The NEW
 569 model increases deposition of other aerosol species as well that have may have health
 570 effects on livestock and wildlife. Predictions of exposure to hazardous chemicals, which

571 may affect human health through ingestion of soil or contaminated produce, also increase
572 with the NEW model.

573



574

575 Figure 13. Accumulated dry deposition (kg/ha) over July 2018 of ammonium aerosol in
576 accumulation plus Aitken modes. NEW model (a), BASE model (b) and NEW – BASE
577 (c). Note that the scale for NEW model (a) is four times the scale for BASE model (b)

578

579 **5.3 Simulation evaluation**

580 The WRF-CMAQ simulations at both 1.33 km grid resolution in the NYC area and 12
581 km grid resolution for the CONUS were evaluated for PM_{2.5} at the U.S. Environmental
582 Protection Agency (EPA) Air Quality System (AQS) sites (AQS;
583 <http://www.epa.gov/aqs>). Evaluations of model runs using the NEW aerosol dry
584 deposition model were compared to the BASE model for summer season when the

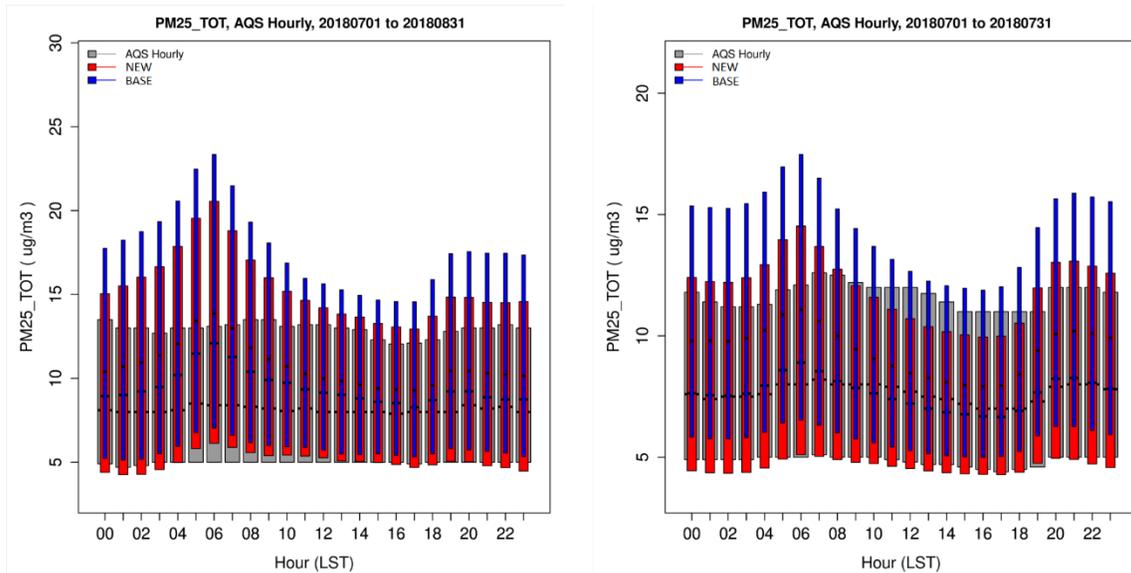
585 differences between NEW and BASE are greatest. Figure 14 shows diurnal bar charts for
586 July and August 2018 for all AQS sites in the 1.33 km grid resolution domain (see figure
587 9 for size of domain) and for July 2018 at AQS sites in the 12 km grid resolution CONUS
588 domain. The colored bars represent the 25th and 75th percentiles of the PM_{2.5}
589 concentration distributions for NEW, BASE, and AQS. The black line in each bar
590 indicates the median value. For the fine grid NE domain both NEW and BASE are high
591 compared to AQS for every hour but NEW is closer to AQS. The diurnal pattern of the
592 modeled concentrations is different from the observation with the greatest concentrations
593 around 6 am and lowest in late afternoon at about 17 LT while the AQS concentrations
594 show very little diurnal variation. This diurnal pattern is typical for modeled
595 concentrations of species with a large ground emitted fraction (i.e., NO_x, CO, and PM_{2.5}).
596 This is a known issue that is related to the suppressed vertical mixing at night and the
597 much greater mixing during the day. The near dawn peak results from the combination
598 of suppressed mixing and high emissions during the morning rush hour. Note that this
599 issue has been improved in recent years through updates to the PBL scheme in the WRF-
600 CMAQ system (e.g., Toro et al., 2021). Another contributing factor is the uncertainty in
601 emissions and especially the hourly attribution of emissions.

602

603 Evaluation of the 12 km CONUS simulations shows similar diurnal pattern with peak
604 concentrations around sunrise and lowest concentrations in afternoon (EDT). The AQS
605 concentrations are more constant over the day than the models but with a slight variation
606 with a similar diurnal pattern. The BASE model is again higher than the observations for

607 every hour, but the NEW model is slightly high during the night and slightly low during
 608 the day.

609



610

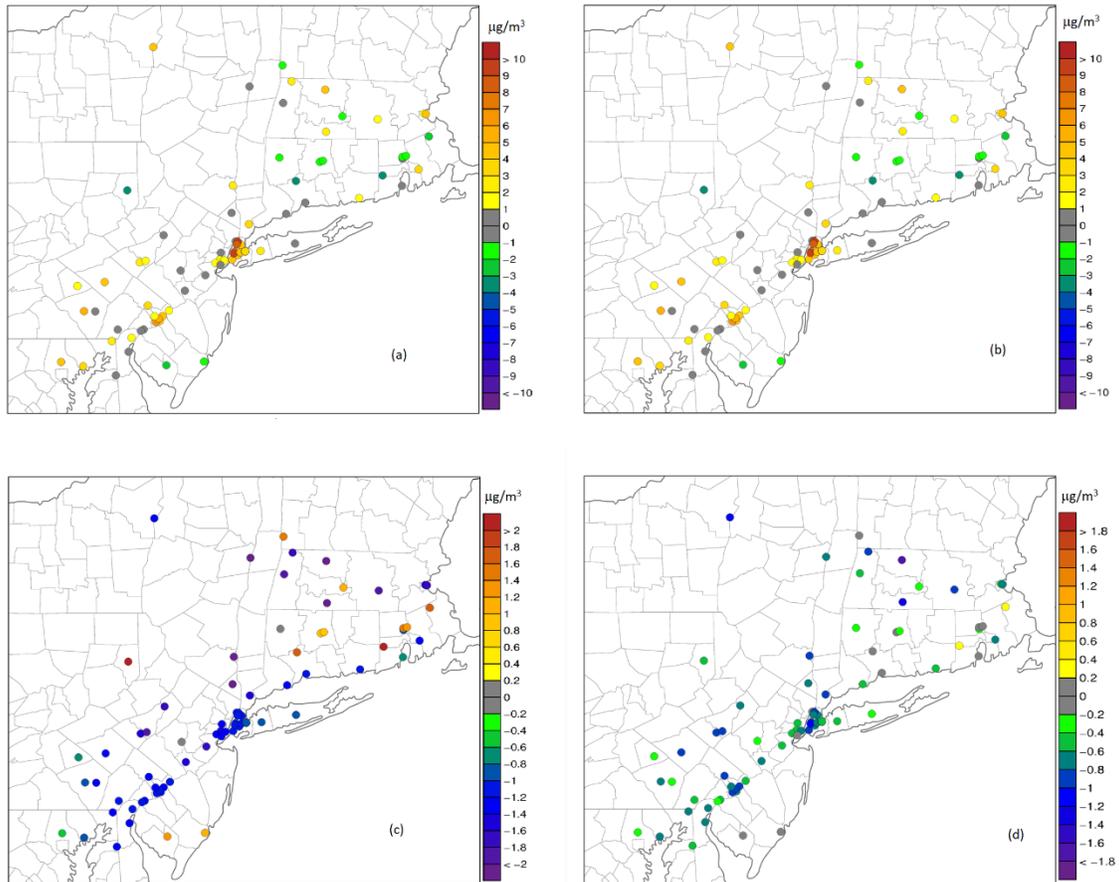
611 Figure 14. Hourly bar plots of PM_{2.5} aggregated over July and August 2018 for the
 612 northeast 1.33 km domain on left and over July 2018 for the 12 km CONUS on right. The
 613 colored bars represent the 25th and 75th percentiles of the PM_{2.5} concentration
 614 distributions for NEW (red), BASE (blue), and AQS observations (grey).

615

616 Spatial evaluation of PM_{2.5} in the 1.33 km domain for NEW and BASE are shown in
 617 Figure 15. For both NEW and BASE the average bias is greatest in NYC where high
 618 emissions are concentrated in small grid cells. The BASE (15b) also has high bias at
 619 most other sites. The NEW (15a) has less high bias and is roughly even between slightly
 620 high and slightly low biases outside of NYC. The lower plots show reduced bias (15d)
 621 and mean error (15c) for NEW compared to BASE at 83% and 93% of the AQS sites,

622 respectively. The urban effects of building area reduce the high bias in the cities but only
623 by a very small amount (less than 0.3%; not shown).

624



625

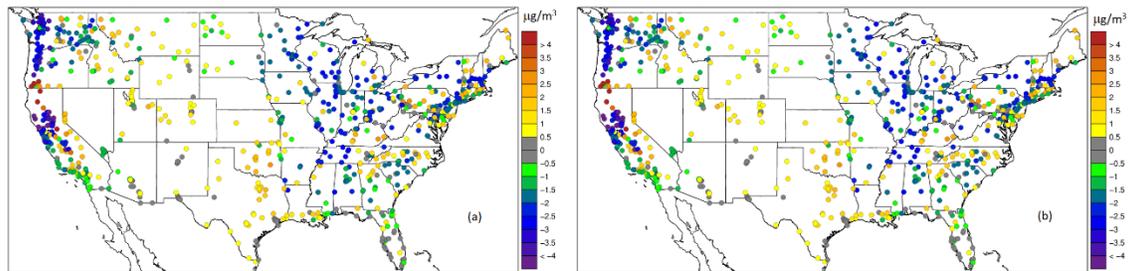
626 Figure 15. Evaluation of modeled hourly $PM_{2.5}$ compared to AQS measurements
627 averaged for July and August 2018: (a) NEW model bias, (b) BASE model bias, (c)
628 NEW – BASE absolute bias difference, and (d) NEW – BASE mean absolute error
629 difference.

630

631 Similar spatial evaluation for the 12 km CONUS simulations is shown in Figure 16.

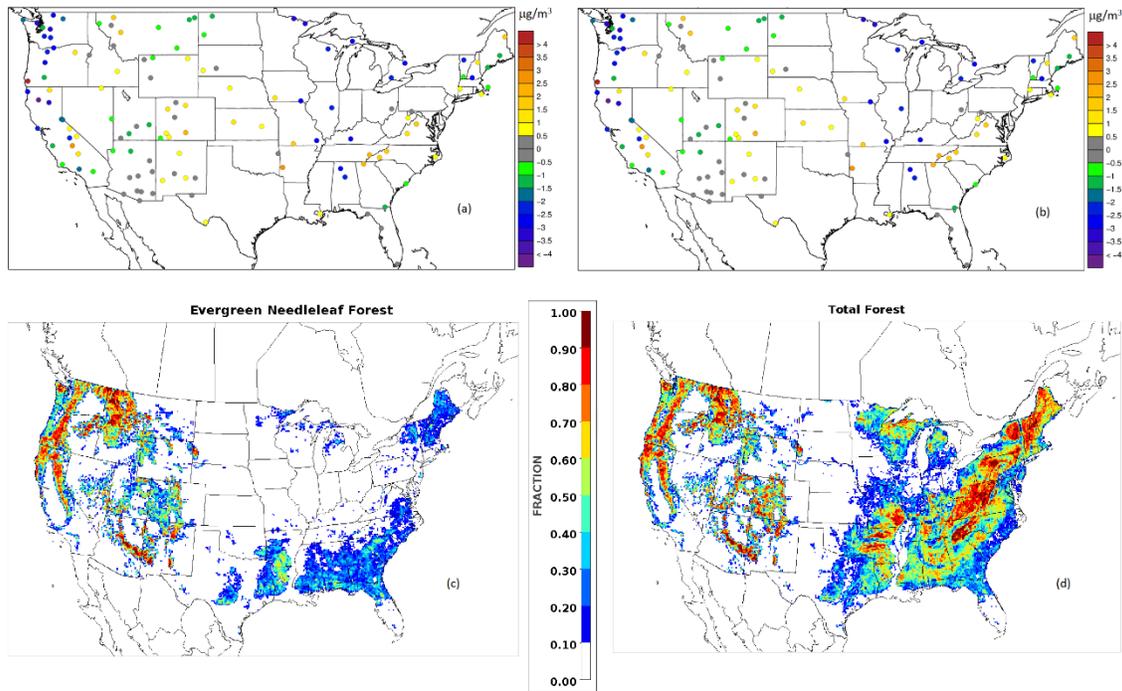
632 Overall, the $PM_{2.5}$ bias averaged over July 2018 at AQS sites is reduced at 64% of the

633 sites and the mean error is reduced at 77% of the sites. In areas where the BASE PM_{2.5}
634 concentrations are biased high such as the Great Lakes region, most of the sites in the
635 east, and the west coast, the NEW model reduces bias and error. In areas where the
636 BASE model was low, such as Texas and the southern plains, the bias is slightly
637 increased although mean absolute error is less affected. In these areas the difference in
638 dry deposition velocity is relatively small because there is less vegetation and not much
639 forest. Note that the increased bias and error at some sites in SW Oregon and northern
640 CA are related to very high PM_{2.5} concentrations caused by wildfires.



641

642 Figure 16. Evaluation of modeled hourly PM_{2.5} compared to AQS measurements
643 averaged for July 2018: (a) NEW – BASE absolute bias difference, and (b) NEW –
644 BASE mean absolute error difference.



645

646 Figure 17. Evaluation of modeled hourly $PM_{2.5}$ compared to IMPROVE measurements
 647 averaged for July 2018: (a) NEW – BASE absolute bias difference, (b) NEW – BASE
 648 mean absolute error difference, (c) gridded evergreen needleleaf forest fraction, (d) total
 649 gridded forest fraction.

650

651 Figure 17 shows $PM_{2.5}$ bias and error averaged over July 2018 at IMPROVE sites (the
 652 Interagency Monitoring of PROtected Visual Environments (IMPROVE;
 653 <http://vista.cira.colostate.edu/Improve/>). Since the IMPROVE network highlights Class I
 654 areas, many of the sites are in forested regions, particularly in the Pacific NW. Figure
 655 17c shows that many sites in WA, OR and northern CA are in evergreen needleleaf forest
 656 and that many of these sites have the greatest reduction in bias and error of $PM_{2.5}$.

657

658 **6. Conclusions and future work**

659 The modeling of dry deposition in general, and aerosol dry deposition in particular,
660 contribute large amounts of uncertainty to air quality and climate models (e.g., Solazzo et
661 al., 2012; Mahowald et al., 2017). Thus, improving the mechanistic underpinnings of dry
662 deposition calculations helps improve accuracy of the modeling systems. The new
663 development to the aerosol dry deposition model was driven by the aggregation of
664 experimental data showing that existing models are unable to accurately replicate the
665 observed relationship between particle size and dry deposition velocity especially for
666 forested areas. Since the quasi-laminar sublayer resistance R_b is usually the controlling
667 resistance during peak deposition conditions (daytime), the new developments focus on
668 revision of the R_b parameterizations, particularly the impaction efficiency. The
669 development process was to revise the impaction efficiency to get better agreement with
670 the aggregate measurements while maintaining physically plausible rationales. The key
671 innovation was to add a second term to the impaction efficiency (equation 7) that
672 represents impaction on microscale obstacles. For broadleaf vegetation the concept is
673 that many species have leaf hairs or other microscale roughness features. However,
674 needleleaf species generally do not have hairs but they may have ridges or other
675 microscale obstacles. Since experimental studies have found that needleleaf species have
676 high aerosol deposition rates, it is theorized that the needle shape itself may be a key
677 factor as discussed in section 5.

678

679 The main impact of the new model is to increase dry deposition velocity in the
680 accumulation size range. This has a large effect on $PM_{2.5}$ especially in forested areas
681 where the dry deposition velocity of accumulation mode mass can increase up to an order

682 of magnitude compared to the current model in CMAQv5.3. For the high-resolution
683 model application to the LISTOS the new model reduced PM_{2.5} concentration up to 40%
684 downwind of NYC in CT for the case shown in Figure 11. For these applications in
685 summer of 2018 where the BASE model generally overestimates PM_{2.5}, the NEW model
686 which has much greater accumulation mode mass dry deposition, results in mostly better
687 agreement with observations. However, aerosol concentrations are very difficult to
688 model accurately not just because of uncertainties in dry deposition but also uncertainties
689 in emissions, transport and diffusion, wet scavenging, and very complex chemistry which
690 involves semi-volatile organic and inorganic species (Appel et al., 2021).

691

692 Continued research on aerosol dry deposition is needed. More field studies, especially
693 for some of the lesser studied vegetation land use types such as croplands, grasslands, and
694 broadleaf deciduous forests, would help to confirm or refute the new model paradigm for
695 different vegetation types and help define parameters such as the macro and micro scale
696 characteristic obstacle dimensions and scaling factors. In addition to comparisons to
697 aggregates of measurements we also plan to model individual field studies in detail where
698 the aerosol dry deposition model is driven by observed micrometeorology and canopy
699 characteristics on an observational timestep basis (30 min or 1 hour).

700

701 Another next step in this research is to include the effects of brown vegetation in the
702 model. Currently, the vegetation fraction and LAI are specified in the WRF-CMAQ
703 system either by land-use category look-up table, where the parameters seasonally vary
704 between minimum and maximum values, or from MODIS satellite fraction of absorbed

705 photosynthetically active radiation (FPAR) and LAI retrievals. In either case, the
706 vegetation parameters are meant to represent live vegetation for evapotranspiration and
707 stomatal uptake of gaseous pollutants. However, for aerosol deposition, brown
708 vegetation can also provide surfaces for deposition. Therefore, we are planning to
709 include MODIS non-photosynthetic vegetation (NPV) and photosynthetic vegetation
710 (PV) fractions in the aerosol dry deposition model. Preliminary tests show that the
711 inclusion of NPV increases aerosol deposition over large areas of the western US.
712 Implementation of NPV has already been included for windblown dust emissions where
713 NPV reduces dust emissions by shielding the soil surface from the wind (Huang and
714 Foroutan 2021).

715

716 **Disclaimer**

717 Although this work was reviewed by EPA and approved for publication, it may not
718 necessarily reflect official Agency policy. Mention of commercial products does not
719 constitute endorsement by the Agency.

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725 **Data Availability Statement**

726 Data used to generate figures and table are available at <https://doi.org/10.23719/1524715>.

727

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