

1 **Inferring the Subsurface Geometry, Stress, and Strength of Slow-moving**
2 **Landslides using 3D Velocity Measurements from the NASA/JPL UAVSAR**

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14 **Key Points:**

- 15 • Landslide thickness can vary by tens of meters within a single landslide
16 • Slow-moving landslides have geometric scaling relations that span scaling values that are
17 typical for soil and bedrock landslides
18 • The largest landslide complexes get larger by increasing area rather than increasing depth
19 • Landslide strength is scale-dependent, such that large landslides tend to be weaker than
20 small landslides

27 **Abstract**

28 The hazardous impact and erosive potential of slow-moving landslides depends on
29 landslide properties including velocity, surface and subsurface geometry, and frequency of
30 occurrence. However, constraints on subsurface geometry are lacking because these types of
31 landslides rarely fully evacuate material to create measurable hillslope scars. Here we use pixel
32 offset tracking with data from the NASA/JPL Uninhabited Aerial Vehicle Synthetic Aperture
33 Radar (UAVSAR) to measure the three-dimensional surface deformation of 134 slow-moving
34 landslides in the northern California Coast Ranges. We apply volume conservation to infer the
35 actively deforming thickness, volume, geometric scaling, stress, and friction angle of each
36 landslide. These landslides move at average rates between $\sim 0.1\text{--}3$ m/yr and have areas of $\sim 7.8 \times$
37 $10^3\text{--}2.63 \times 10^6$ m², inferred mean thicknesses of $\sim 0.4\text{--}22$ m, and volumes of $\sim 7.08 \times 10^3\text{--}9.75 \times$
38 10^6 m³. The best-fit volume-area geometric scaling exponent is $\gamma = 1.2\text{--}1.5$, indicating that these
39 landslides fall between typical soil and bedrock landslide scaling. A rollover in the scaling
40 relationship suggests that the largest landslide complexes in our dataset become large primarily
41 by increasing in area rather than thickness. In addition, the slow-moving landslides display scale-
42 dependent frictional strength, such that large landslides tend to be weaker than small landslides.
43 This decrease in frictional strength with landslide size is likely because larger landslides are
44 composed of higher proportions of weak material. Our work shows how state-of-the-art remote
45 sensing techniques can be used to better understand landslide processes and quantify their
46 contribution to landscape evolution and hazards to human safety.

47 **1 Introduction**

48 Landslides are a major natural hazard and are often the dominant process that erodes
49 mountainous landscapes (Korup et al., 2007; Larsen et al., 2010; Mackey & Roering, 2011;

50 Simoni et al., 2013). Both their hazardous impact and erosive potential depend on landslide
51 properties including the velocity, surface and subsurface geometry, and frequency of occurrence.
52 Measuring these landslide properties is challenging because landslides exhibit a wide range of
53 velocities (mm/yr to m/s), spatial areas ($10^0 - 10^8 \text{ m}^2$), and volumes ($10^{-1} - 10^{10} \text{ m}^3$), and can
54 occur in large numbers (hundreds to tens of thousands) over broad spatiotemporal scales (Cruden
55 & Varnes, 1996; Hungr et al., 2014; Lacroix, Handwerger, et al., 2020; Larsen et al., 2010).
56 Importantly, the landslide failure style also impacts our ability to measure landslide properties,
57 such as thickness and volume, which can strongly influence runout and erosion rate (e.g., Korup
58 et al., 2007; Larsen et al., 2010; Legros, 2002). Some landslides create clear and identifiable
59 scars and deposits by evacuating material from the hillslope, making it possible to directly
60 measure landslide properties from field data, digital elevation models (DEMs), and remote
61 sensing observations (Bessette-Kirton et al., 2018; Warrick et al., 2019; Wartman et al., 2016).
62 However, for landslides that move slowly for years or centuries (Mackey et al., 2009; Rutter &
63 Green, 2011), i.e., slow-moving landslides (Lacroix et al., 2020), and do not create hillslope
64 scars, it is difficult to constrain their thickness and volume because data are usually limited to
65 isolated point measurements from boreholes (Schulz et al., 2018; Simoni et al., 2013; Travelletti
66 & Malet, 2012), which do not capture the spatial variability exhibited by these landslides. It is
67 therefore advantageous to develop and apply tools and methods that can be used to construct
68 large inventories of slow-moving landslides and quantify their surface and subsurface properties.

69 Modern remote sensing tools, such as synthetic aperture radar (SAR), optical imagery,
70 and lidar, provide high-resolution measurements of topography and ground surface deformation
71 that can be used to identify and monitor landslides with millimeter- to centimeter-scale accuracy
72 at spatial resolutions of 5 to 100 meters. Recent work using pixel offset tracking and SAR

73 interferometry with these data has quantified the two-dimensional (2D) and three-dimensional
74 (3D) surface deformation of slow-moving landslides (Aryal et al., 2015; Booth et al., 2020; Hu et
75 al., 2020; Lacroix et al., 2020; Stumpf et al., 2017; Travelletti et al., 2014). These studies, along
76 with numerous ground-based investigations (e.g., Iverson & Major, 1987; Malet et al., 2002;
77 Schulz et al., 2017), have shown that slow-moving landslides exhibit non-uniform spatial and
78 temporal kinematic patterns. In addition, these high-resolution 3D surface deformation
79 measurements can be used to infer the thickness and subsurface geometry of the actively moving
80 part of the landslide (Aryal et al., 2015; Booth et al., 2020; Booth, Lamb, et al., 2013; Delbridge
81 et al., 2016; Hu et al., 2020). These studies suggested that active landslide thickness can vary by
82 tens of meters within a single landslide, and the slip surfaces have an irregular and bumpy
83 morphology that differs considerably from commonly assumed, idealized geometric forms, such
84 as semicircles, ellipsoids, and log spirals (e.g., Michel et al., 2020 and references therein). These
85 large changes in thickness within a single landslide mass have important implications for
86 estimating volume and sediment flux, designing field instrumentation and landslide mitigation
87 strategies, and determining the stresses that control landslide kinematics. Although techniques
88 that invert surface observations for subsurface characteristics are becoming more common, most
89 studies have focused on individual landslides occurring under different and site-specific
90 environmental conditions, making it difficult to identify more generic geometric scaling relations
91 for slow-moving landslides.

92 In this study, we use data from the NASA/JPL Uninhabited Aerial Vehicle Synthetic
93 Aperture Radar (UAVSAR) to construct an inventory of 134 active slow-moving landslides in a
94 ~4700 km² area of the northern California Coast Ranges between 2016 and 2019 (Figure 1).
95 These landslides occur in the Eel River catchment, a region well known for its slow-moving

96 landslides, and are driven by high seasonal rainfall (Bennett, Roering, et al., 2016; Booth,
97 Roering, et al., 2013; Handwerger et al., 2013, 2015; Handwerger, Fielding, et al., 2019;
98 Handwerger, Huang, et al., 2019; Kelsey, 1978; Mackey et al., 2009; Mackey & Roering, 2011;
99 Roering et al., 2009, 2015; Schulz et al., 2018). The landslides are underlain by the Central Belt
100 Franciscan mélangé, a mechanically weak and pervasively sheared bedrock with an argillaceous
101 matrix that surrounds blocks of stronger rock types, including sandstone, chert and greenstone
102 (Jayko et al., 1989; Jennings et al., 1977; McLaughlin et al., 1982, 2000). We measure the 3D
103 surface deformation and geometry of each landslide, and use these data in a volume conservation
104 framework to invert for their active thickness, volume, stress, and strength. We derive new
105 geometric scaling relations for slow-moving landslides and make comparisons with a worldwide
106 inventory of soil and bedrock landslides. Our work is the first to perform thickness inversions for
107 a large inventory of landslides, and this approach could be applied to other groups of slow-
108 moving landslides around the world. Our work also shows how state-of-the-art remote sensing
109 techniques can be used to better understand landslide processes and quantify their contribution to
110 landscape evolution.

111 **2 Materials and Methods**

112 **2.1 UAVSAR Data and Processing**

113 We use SAR data acquired by the NASA/JPL UAVSAR airborne system for our
114 landslide investigation. UAVSAR has a left-looking radar attached to a NASA Gulfstream III
115 airplane that operates with a L-band wavelength (~ 23.8 cm) and a swath width of ~ 20 km.
116 UAVSAR data have a pixel spacing of 1.67 m in the range direction (measured along the line-of-
117 sight, LOS) and 0.6 m in the azimuth direction (measured along the UAVSAR flight direction).
118 We designed the UAVSAR data collection for the northern California Coast Ranges site

119 specifically to monitor a large quantity of slow-moving landslides that were initially identified
120 by several previous studies (Bennett, Miller, et al., 2016; Handwerger et al., 2015; Kelsey, 1978;
121 Mackey & Roering, 2011; Roering et al., 2009). Some of these UAVSAR data were used in a
122 recent study by Handwerger, Fielding, et al. (2019) to analyze changes in landslide activity due
123 to extreme rainfall. We collected data on 4 partially overlapping flight paths to increase data
124 redundancy and to provide between 4 and 8 independent deformation measurements (Figure 1).
125 There were 12 data acquisitions at our field site between April 2016 and May 2019. The time
126 between data acquisitions ranges between 47 and 237 days, with a mean of 104 days (Table S1).
127 UAVSAR Single-Look Complex (SLC) data are freely available at <https://uavsar.jpl.nasa.gov/>.

128 We perform pixel offset tracking on the coregistered UAVSAR stack SLC data using the
129 *Ampcor* application, which is part of the JPL InSAR Scientific Computing Environment (ISCE)
130 version 2 software package (Rosen et al., 2012). Pixel offset tracking (sometimes referred to as
131 subpixel correlation) uses cross-correlation between SAR amplitude images to quantify image
132 offsets (i.e., displacement) due to ground surface motion in two dimensions; 1) the range or look
133 direction, and 2) the azimuth or along-track direction (Fialko et al., 2001; Fielding et al., 2020;
134 Pathier et al., 2006). We use the terms range/look direction and azimuth/along-track direction,
135 interchangeably. Pixel tracking has a precision up to $\sim 1/10$ of the pixel size, which corresponds
136 to ~ 6 cm in the along-track direction and ~ 17 cm in the range direction for a pair of UAVSAR
137 images. Although this technique is less precise than conventional InSAR, it does not involve
138 phase unwrapping and thus is better suited for measuring the decimeter- to meter-scale
139 displacements commonly displayed by many slow-moving landslides (Lacroix, Handwerger, et
140 al., 2020). To account for the differences in the range and along-track pixel size, we use a cross-
141 correlation window length of 128 pixels with a skip size of 32 pixels (distance between matching

142 window calculations) in the along-track direction and a cross-correlation window width of 64
143 pixels with a skip size of 16 pixels in the range direction, resulting in a window size of 77 m by
144 107 m. We geocode the pixel offset measurements to a 0.4 arcsecond (~ 12 m) pixel using the
145 TanDEM-X DEM provided by the German Aerospace Center (DLR). We process all possible
146 combinations of pixel offset tracking pairs, which results in 66 pixel offset tracking maps on
147 each track (264 in total) with single pair time spans ranging from 47 to 1148 days (Table S1).
148 We excluded 35 poor-quality pixel offset tracking maps from our analysis. We found these poor-
149 quality data tend to result from long duration pairs that exceed ~ 2 years, which are subject to
150 numerous changes in the ground surface (e.g., vegetation changes, anthropogenic changes) that
151 can deteriorate the cross-correlation result. We convert all of the offset maps to mean velocities
152 and then take the temporal average of the 31 remaining pixel offset velocity maps to make a
153 mean velocity map for our thickness inversions.

154

155 2.2 Three-dimensional Ground Surface Deformation

156 To solve for 3D deformation from SAR requires at least three independent measurements
157 of surface deformation. Each UAVSAR flight path provides two independent measurements of
158 surface motion from pixel offset tracking (i.e., along-track and range). Therefore, using pixel
159 offset tracking velocity maps, data from at least two flights is required for 3D inversions.
160 Because UAVSAR acquires data on 4 different flight paths in our field area (Figure 1), we have
161 a maximum of 8 deformation measurements in the central region of our field area where all 4
162 flight paths overlap and a maximum of 4 deformation measurements in the northern and southern
163 extents where only 2 flight paths overlap. Thus, we are always able to achieve an overdetermined
164 3D inversion.

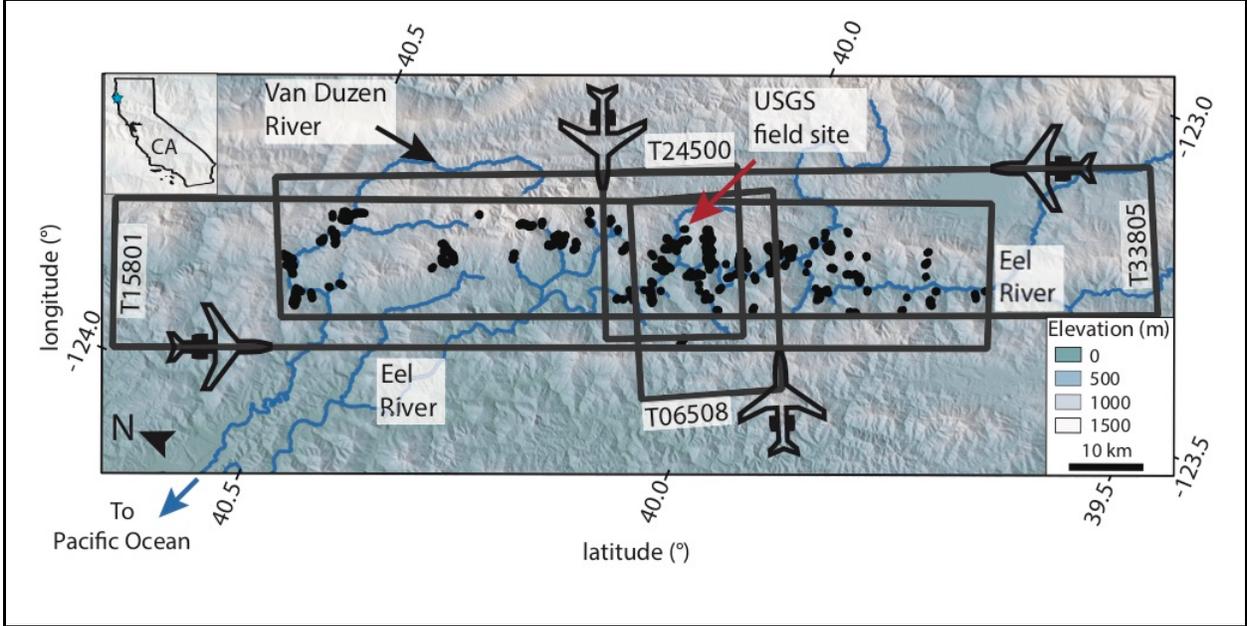


Figure 1. Map of our northern California Coast Ranges study site. Black polygons outline the active landslides analyzed in this study. Elevation (m) shown by green to white color gradient. Black boxes show the left-looking UAVSAR swaths and corresponding track numbers with airplanes showing flight direction. Red arrow shows the location of a USGS landslide field site. Blue lines show major rivers and some tributaries in landslide areas. Inset shows a map of California with a star corresponding to the study site.

165

166 Each deformation measurement from pixel tracking is composed of the true displacement

167 vector projected onto the along-track or LOS direction of the UAVSAR. We use a least-squares

168 inversion to isolate the east, north, and vertical components of deformation defined in the form

169 $\mathbf{d} = \mathbf{Gm}$,

170

$$\begin{bmatrix} v_{rng1} \\ v_{azi1} \\ \vdots \\ v_{rng,M} \\ v_{azi,M} \end{bmatrix} = \begin{bmatrix} \cos \xi_1 \sin \theta_1 & \sin \xi_1 \sin \theta_1 & -\cos \theta_1 \\ \cos \xi_1 & \sin \xi_1 & 0 \\ \vdots & \vdots & \vdots \\ \cos \xi_M \sin \theta_M & \sin \xi_M \sin \theta_M & -\cos \theta_M \\ \cos \xi_M & \sin \xi_M & 0 \end{bmatrix} \begin{bmatrix} v_{ew} \\ v_{ns} \\ v_{ud} \end{bmatrix}, \quad (1)$$

171 where $v_{rng,M}$ is the range (or look direction) velocity, $v_{azi,M}$ is the azimuth (or along-track
172 direction) velocity, M is the flight path number (minimum of 2 needed for pixel offset tracking),
173 ξ is the UAVSAR heading direction (i.e., along track direction) with counterclockwise as
174 positive, θ is the UAVSAR look angle, and v_{ew} , v_{ns} , v_{ud} are the east-west, north-south, and
175 vertical components of velocity, respectively.

176 The overdetermination of the 3D inversion allows us to reduce inversion error and
177 constrain the uncertainty from the inversion. To constrain the uncertainty, we repeat the 3D
178 inversion multiple times using different combinations of v_{rng} and v_{azi} . For instance, for landslides
179 with 8 deformation measurements (i.e., 4 range and 4 azimuth measurements), we perform the
180 3D inversion 198 times using between 3 and 8 deformation measurements. We then take the
181 mean and standard deviation of all of the inversions and use these values as the 3D velocities and
182 inversion uncertainty, respectively. We further constrain the uncertainty in our velocity
183 measurements by examining the apparent deformation rate of stable hillslopes. To reduce noise
184 and error (i.e., unrealistically large displacements), we apply velocity thresholds and mask out
185 pixels with apparent velocities > 50 m/yr, which is much faster than the typical velocity range
186 displayed by the northern California Coast Ranges landslides (Bennett, Roering, et al., 2016;
187 Handwerker, Fielding, et al., 2019; Roering et al., 2015). We also mask out pixels that have
188 mean velocities less than their inversion uncertainty and use nearest neighbor interpolation to fill
189 in these masked pixels.

190

191 2.3 Landslide Thickness Inversion

192 We use 3D surface velocity measurements from pixel offset tracking to infer the
193 thickness, volume, and shear zone geometry of the active parts of each landslide using a

194 conservation of volume approach. We apply the method originally described by Booth, Lamb, et
 195 al., (2013) and more recently by Booth et al. (2020), which assumes that during our ~3 year
 196 study period, the measured surface velocity is representative of the depth-averaged velocity, the
 197 sliding surface does not change in time, there is minimal direct erosion or deposition of the
 198 landslide surface, and the landslide material density is uniform and constant. While landslides
 199 may violate these assumptions in general, they are reasonable for our study area for the following
 200 reasons: (1) at the Two Towers landslide, a U.S. Geological Survey (USGS) instrumented
 201 landslide in our study site (Schulz et al., 2018), the measured surface velocity was approximately
 202 equal to the depth-averaged velocity, and a narrow shear zone was identified (Figure S1); (2) the
 203 landslides were continuously active over the time periods that 3D displacements were measured,
 204 suggesting movement on the same slip surface; (3) minor amounts of direct surface erosion or
 205 deposition were likely confined to gully systems on the landslides' surfaces, which occupy a
 206 small percentage of the landslides' surface area (~1%) and therefore have a minimal effect on the
 207 inversion; and (4) dilation/compaction or shrinking/swelling that would cause changes in density
 208 is likely on the order of centimeters or less, which is small compared to surface velocity
 209 gradients (Booth et al., 2020; Delbridge et al., 2016; Iverson, 2005; Schulz et al., 2018), thus
 210 having limited influence of the measured 3D surface velocity. Therefore, for a landslide of
 211 constant density with no erosion or deposition, conservation of volume implies that

$$212 \quad v_{ud} = \nabla \cdot (\bar{u}h) + u_{surf} \cdot \nabla z_{surf}, \quad (2)$$

213 where v_{ud} is the vertical component of the 3D landslide surface displacement vector, h is the
 214 active landslide thickness, u_{surf} is the vector of horizontal components of landslide surface
 215 velocity, \bar{u} is the depth-averaged vector of horizontal components of landslide velocity, and z_{surf}
 216 is the surface elevation measured from the 12 m TanDEM-X DEM. The first term on the right-

217 hand side of equation 2 is the contribution of flux divergence to the vertical component of the
 218 surface velocity, and the second term is the contribution due to advection of the sloped land
 219 surface. Because UAVSAR measures the velocity of the ground surface, u_{surf} , we assume that
 220 $\bar{u} = fu_{surf}$, where f is a constant that characterizes the thickness of the shear zone at the base of
 221 the landslide relative to the total landslide thickness. We constrain f using borehole inclinometer
 222 data from two boreholes at the USGS field station on the Two Towers landslide (supporting
 223 information and Figure S1). Unfortunately, the Two Towers landslide is not detectable with pixel
 224 tracking from UAVSAR data because the landslide is small (250 m long and 40 m wide) and
 225 moving too slowly (maximum speed ~ 6 cm/yr). Therefore, we are not able to directly test our
 226 thickness inversion method on the Two Towers landslide. Nonetheless, the ground-based data
 227 provide key information to constrain the value of f . Nearly all shear strain occurs in a zone that is
 228 < 0.3 m in one borehole and < 0.6 m in the second borehole. Using these data, we find that f
 229 ~ 0.96 , which indicates that the landslide moves along a narrow shear zone with the material
 230 above translating essentially as a rigid block. For simplicity, we assume that $f = 1$ and that the
 231 landslides move as a rigid block. Other studies in California have also found that landslides
 232 move as a rigid plug above a narrow shear zone such that $f \sim 1$ is a reasonable approximation
 233 (Keefer & Johnson, 1983). Although f generically represents the ratio of depth-averaged to
 234 surface velocity, it can be related to specific rheologies if desired (Booth, Lamb, et al., 2013;
 235 Delbridge et al., 2016).

236 Incorporating f into equation 2 gives

$$237 \quad v_{ud} = \nabla \cdot (fu_{surf}h) + u_{surf} \cdot \nabla z_{surf}, \quad (3)$$

238 which is a statement of conservation of volume in a Lagrangian reference frame (Booth et al.,
 239 2020; Delbridge et al., 2016). We discretize equation 3 using centered finite differences,

240 rearrange it as a system of linear equations, and then solve for thickness by minimizing, subject
 241 to non-negative constraints,

$$242 \quad |Xh - b|^2 + \alpha^2 |\nabla^2 h|^2, \quad (4)$$

243 where X is a diagonally dominant matrix that contains the depth-averaged horizontal velocity
 244 data, b is a vector containing the topographic gradient and surface horizontal and vertical
 245 velocity data, and α is a damping parameter to regularize the ill-posed inverse problem. We
 246 explore a wide range of α from 10^{-3} to 10^1 and determine the best level of regularization using
 247 the Generalized Cross-Validation method (supporting information and Figure S2). We resample
 248 our ~ 12 m pixel spacing grid to square 10 m x 10 m pixel and perform the thickness inversion in
 249 the MATLAB software package using the CVX program, a package for specifying and solving
 250 convex programs (Grant & Boyd, 2014). The inferred thickness values represent the best
 251 solution that does not violate conservation of volume and assumes that the surface velocity is
 252 equal to the depth-averaged velocity. It is important to further emphasize that the thickness
 253 inversions are only relevant to the active parts of landslides such that there needs to be detectable
 254 surface deformation to invert for the landslide thickness. Specifically, the values of b (equation
 255 4) need to differ from background values on known stable ground to infer non-zero thicknesses.
 256 Landslides or areas and kinematic zones within landslides that are not moving are therefore
 257 considered to have zero depth. Thickness in this study therefore means the “active thickness”
 258 during our study period.

259 Since both the matrix X and the vector b contain data with uncertainties, and the damping
 260 parameter necessarily introduces bias, estimating total uncertainty of the resulting thickness
 261 model is not straightforward. However, we make a minimum estimate following standard
 262 techniques from inverse theory, which reflects uncertainty in b only (supporting information).

263 Thickness uncertainty estimated in this way increases with landslide size and ranges from
264 approximately ± 1.5 m to ± 4 m from the smallest to largest landslides in the study area (Figure
265 S3).

266

267 2.4 Landslide Inventory and Geometric Scaling

268 To select landslides for 3D surface velocity and thickness inversions, we assemble a new
269 inventory of active landslides in our ~ 4700 km² study area in the northern California Coast
270 Ranges that includes only those landslides that show a significant deformation signal using the
271 pixel offset tracking method. This limits our analysis to the faster-moving landslides that exhibit
272 rates of decimeters to meters per year. Our landslide inventory was guided by a number of pre-
273 existing landslide inventories for the northern California Coast Ranges (Bennett, Miller, et al.,
274 2016; Handwerger, Fielding, et al., 2019; Kelsey, 1978; Mackey & Roering, 2011). We map the
275 landslide boundaries in QGIS using the 3D velocity maps, hillshade maps constructed from 1 m
276 pixel spacing lidar provided by OpenTopography (Roering, 2012), the ~ 12 m pixel spacing
277 TanDEM-X DEM, and Google Earth imagery. Because slow-moving landslides display non-
278 uniform spatial kinematic zones and complex kinematic histories (e.g., Nereson & Finnegan,
279 2018; Schulz et al., 2017; Stumpf et al., 2017), there are often differences between the landslide
280 boundaries mapped with kinematic data and those mapped based on geomorphic interpretation of
281 hillshades or aerial photos. These differences in mapping are especially important for our
282 thickness inversions because including the parts of landslides that are not currently moving can
283 cause the thickness inversion to produce unreliable results. Therefore, we use the temporally
284 averaged landslide velocity and only map areas of each landslide that are moving during our
285 study period. We then use QGIS to quantify the spatial metrics of each landslide, including

286 length, average width (defined as area divided by length), area, and slope angle. We also report
 287 the mean, median, 75th percentile, and maximum horizontal velocity, 3D velocity magnitude,
 288 and 3D inversion velocity errors for each landslide.

289 We then derive empirical geometric scaling relations for landslide thickness (h) and
 290 volume V from the measured landslide area A . Geometric scaling relations are commonly used to
 291 quantify erosion rates of large inventories of landslides and are important for understanding
 292 landslide mechanics (e.g., Guzzetti et al., 2009; Larsen et al., 2010; Milledge et al., 2014). These
 293 relations take the form of a power function where

$$294 \quad V = c_V A^\gamma \text{ and } h = c_h A^\zeta, \quad (5a \text{ and } 5b)$$

295 where γ and ζ are scaling exponents and c_V and c_h are the intercepts. We constrain the
 296 coefficients of these power functions by log-transforming our data and finding the best-fit
 297 parameters with 95% confidence intervals using a linear least square inversion in MATLAB.

298 Larsen et al. (2010) showed that these scaling relations hold over 9 orders of magnitude in area
 299 and 12 orders of magnitude in volume.

300

301 2.5 Stress and Frictional Strength

302 We estimate the basal shear stress and frictional strength of each landslide by following
 303 the 3D Simplified Janbu method (Bunn et al., 2020; Hungr, 1987; Hungr et al., 1989;
 304 Leshchinsky, 2019). This method assumes that the vertical intercolumn shear forces are
 305 negligible. Each landslide is discretized into 3D columns with a surface area S_{basal} and total
 306 weight W . The basal surface area is defined by

$$307 \quad S_{basal} = \Delta x \Delta y \frac{1 - \sin^2 \beta_y \sin^2 \beta_x}{\cos \beta_y \cos \beta_x}, \quad (6)$$

308 where Δx and Δy are the grid spacing in the x and y direction, respectively, β_x is the local dip
 309 angle perpendicular to the direction of motion and β_y is the local dip in the direction of motion.

310 The normal force N at the base of each column is defined by

$$311 \quad N = \frac{W + pS_{basal} \tan \phi \sin \beta_x / F}{\cos \Delta_z \left(1 + \frac{\sin \beta_x \tan \phi}{F \cos \Delta_z} \right)}, \quad (7)$$

312 where p is the mean pore pressure acting at the base of each column, ϕ is the residual friction
 313 angle, F is the factor of safety, and Δ_z is the local dip angle defined in terms of the motion-
 314 parallel and motion-perpendicular dips by

$$315 \quad \cos \Delta_z = \left(\sqrt{\frac{1}{1 + \tan^2 \beta_y + \tan^2 \beta_x}} \right). \quad (8)$$

316 Finally, F is defined by

$$317 \quad F = \frac{\sum (N - pS_{basal}) \tan \phi \cos \beta_x}{\sum N \cos \Delta_z \tan \beta_x}, \quad (9)$$

318 where the summation is over all columns. The numerator is the resisting force, with the term in
 319 the parentheses defining the effective normal force, and $\tan \phi$ is the friction coefficient, and the
 320 denominator is the shear force. Cohesion is assumed to be negligible since our landslides are
 321 already moving. We set $F = 1$ (i.e. balanced forces at failure) and solve for shear stress, effective
 322 normal stress (shear force and normal force divided by S_{basal} , respectively), and friction angle
 323 under both dry and fully saturated (hydrostatic conditions) end members to produce a minimum
 324 and maximum estimate. Table S2 shows the dry and wet landslide density values used for our
 325 calculations.

326

327 **3 Results**

328 **3.1 Landslide Inventory and 3D Velocity**

329 We identified 134 active landslides in our northern California Coast Ranges field site
330 (Figure 1), 19 of which were unmapped by previous studies (Bennett, Miller, et al., 2016;
331 Handwerger, Fielding, et al., 2019; Mackey & Roering, 2011). These landslides have average
332 widths from 66 to 556 m, lengths from 68 to 4727 m, areas from 7.8×10^3 to 2.63×10^6 m², and
333 mean slope angles from 10 to 29 degrees (Table S3). Each landslide exhibited a non-uniform
334 spatial velocity pattern (see examples in Figure 2). The spatial kinematic patterns remain fixed
335 during our study period and are similar to those mapped in previous studies (see Bennett,
336 Roering, et al., 2016; Handwerger, Fielding, et al., 2019; Mackey & Roering, 2011). The
337 characteristic 3D velocity magnitude, defined as the 75th percentile value for each landslide, was
338 calculated as $v_{3D} = (v_{ns}^2 + v_{ew}^2 + v_{ud}^2)^{1/2}$. The 3D velocity magnitude ranged from 0.162 to 2.92
339 m/yr. The average 3D velocity for the entire inventory was 0.622 ± 0.513 m/yr (± 1 standard
340 deviation). The landslide motion was always primarily in the downslope direction (see example
341 in Figures 2e and 2f), but at different locations we do measure areas of both uplift and
342 subsidence within a single landslide (see example in Figure 2d). We note that local surface uplift
343 occurs when the vertical component of the velocity vector dips less steeply than the topographic
344 surface at a given point. As a result, the vertical velocity is often still negative even in areas
345 where the topographic surface is locally being uplifted, and only when the vertical motion is
346 upwards relative to horizontal do we observe positive vertical velocities. The 3D velocity
347 uncertainty from the 3D inversion (equation 1) averaged over all 134 landslides was 0.1432
348 m/yr. We report the uncertainty for each individual landslide in Table S3. The 3D velocity
349 magnitude uncertainty from examining the apparent velocity of stable hillslopes was ≤ 0.1 m/yr.

350 We classified the slow-moving landslides into three subgroups based on their geometry
351 and kinematic patterns (Table S3). Figure 2 shows three example landslides which we define as

352 slumps, earthflows, and landslide complexes. The landslide complex shown in Figure 2 is the
353 largest landslide in our dataset and is also known as the Boulder Creek landslide in several other
354 studies (Bennett, Miller, et al., 2016; Bennett, Roering, et al., 2016; Handwerger et al., 2015,
355 2015; Handwerger, Fielding, et al., 2019; Handwerger, Huang, et al., 2019; Mackey & Roering,
356 2011; Roering et al., 2009). We defined slumps as landslides with lower length/width aspect
357 ratios (median = 1.57 ± 1.00 , ± 1 standard deviation), a strong signal of positive vertical velocity
358 components in the toe and negative vertical velocity components in the source area, and one
359 primary kinematic zone (Figure 2a). We defined earthflows landslides as those with medium
360 aspect ratios (median = 3.56 ± 1.88 , ± 1 standard deviation), one primary kinematic zone, and
361 small magnitude, but mostly negative, vertical velocity components (Figure 2b). And we defined
362 landslide complexes as those with higher aspect ratios (median = 5.13 ± 2.34 , ± 1 standard
363 deviation), that are composed of multiple kinematic zones or even multiple landslides that
364 coalesce into a single landslide mass (Figure 2c). Landslide complexes are relatively common in
365 areas with slow-moving landslides (Cerovski-Darriau & Roering, 2016; Keefer & Johnson,
366 1983; Simoni et al., 2013). The mean 3D velocity magnitude was 0.585, 0.606, and 0.670 m/yr
367 for slumps, earthflows, and landslide complexes, respectively.

368

369 3.2 Thickness, Volume, and Geometric Scaling Relations

370 The non-uniform kinematic patterns exhibited by these landslides are also reflected in
371 their inferred subsurface geometry (Figure 3). We find that the thickness of each landslide varies
372 spatially and can range from zero to tens of meters within the landslide boundaries. The slip
373 surfaces are generally concave-up, but are rough and irregular in places, especially for landslide
374 complexes. The mean active thickness of the individual landslides ranged from 0.4 to 22.4 m,

375 and the maximum active thickness ranged from 2.25 to 89.6 m. The mean, median, min, max,
 376 and standard deviation active thickness for each landslide are reported in Table S3. Below we
 377 describe our findings for the three example types of landslides shown in Figure 2. We note again
 378 that these landslides represent their subgroups to first order.
 379

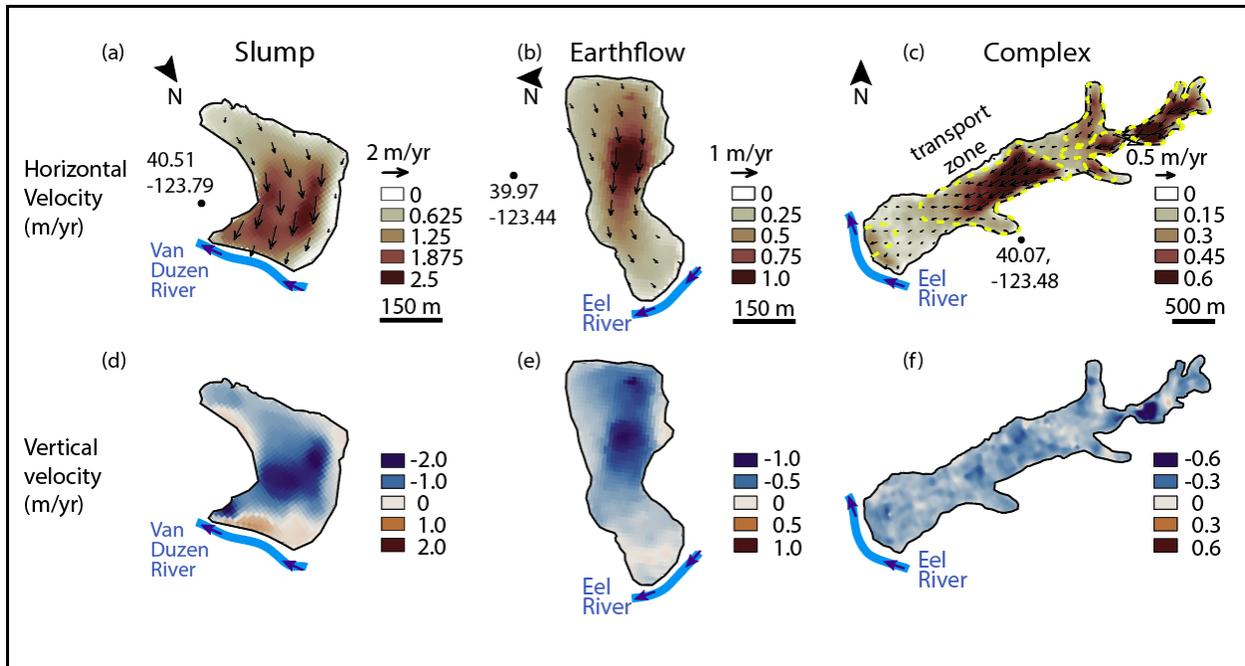


Figure 2. 3D velocity maps for three example landslide types. (a–c) Horizontal velocity map for three example landslides. Black arrows show horizontal vectors. Dashed yellow lines in (c) correspond to landslide zones nested within the landslide complex. (d–f) Vertical velocity maps for the three landslides. Negative values correspond to vertically downward motion. Thick blue lines show the approximate location of the river channel at the toe of each landslide with dark blue arrows showing water flow direction.

380
 381 The example slump has one primary deep zone with a mean active thickness of ~ 7.5 m,
 382 maximum active thickness of ~ 31 , and standard deviation of ~ 8.7 m. The slip surface has a

383 concave-up profile. The slope of the slip surface deviates from the ground surface and is steeper
384 near the headscarp and gentler near the toe. Some areas within the head of the landslide are
385 inferred to have no active depth because the values of b (equation 4) are similar in magnitude to
386 background values there. The example earthflow generally has a concave-up slip surface with
387 some irregular bumps. The slip surface more closely mimics the ground surface in the main
388 transport zone and the landslide has a mean active thickness of ~ 15 m, max active thickness of
389 ~ 62 m, and a standard deviation of ~ 17 m. Lastly, the example landslide complex (i.e., Boulder
390 Creek landslide complex) has several different active zones, each with an alternating concave-up
391 and convex-up slip surface profile. The landslide slip surface is rough and irregular over the
392 length of the entire landslide, but each deep zone corresponds to the different kinematic units that
393 comprise the landslide complex (Figure 2c). The landslide complex has a mean active thickness
394 of ~ 3.7 m, maximum active thickness of ~ 29 m, and a standard deviation of ~ 4.6 m. These
395 thicknesses are small compared to the two previous examples because many areas of the
396 landslide do not have a resolvable active thickness.

397 The Boulder Creek landslide complex shows that landslide complexes may correspond to
398 smaller, faster, and possibly shallower features that are superimposed on a larger, slower, and
399 possibly deeper-seated failure. If multiple failure planes are indeed present at these landslides,
400 that would violate the assumption of a constant f throughout the landslide and cause unreliable
401 thickness estimates. To minimize those potential effects, we performed separate thickness
402 inversions for any isolated, faster-moving areas of a larger landslide complex, as well as for the
403 entire landslide complex as a whole. If results were substantially different, we adopted the more
404 reliable results for the smaller isolated landslides.

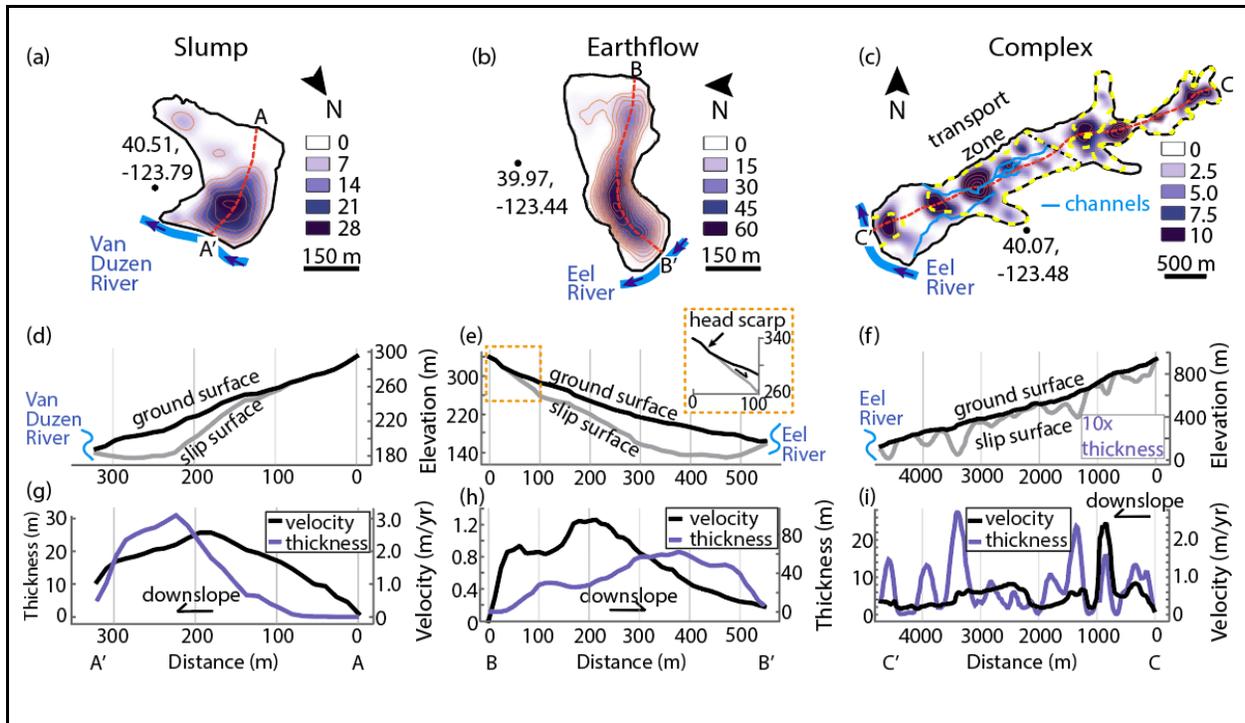


Figure 3. Landslide thickness inversions for example slump, earthflow, and landslide complex. (a–c) Landslide thickness maps. Thin orange lines with labels show 5-meter thickness contours. Red dashed line shows profiles plotted in (d–i). Black dots show latitude and longitude coordinates. Thick blue lines show rivers and thin blue lines show deep channels incised into the landslide body. Dashed yellow lines in (c) correspond to landslide zones nested within the landslide complex. (d–f) Ground surface and slip surface elevation profiles. (e) Inset shows zoom-in of landslide headscarp. In subplot (f) the results of depth inversion are vertically exaggerated by a factor of 10 relative to the elevation profile. (g–i) Landslide thickness and 3D velocity magnitude profiles. Arrow shows direction of downslope motion.

405

406 For example, we delineated the Boulder Creek landslide complex into 5 smaller landslides based

407 on the separate kinematic zones and performed thickness inversions for each kinematic zone or

408 sub-landslides (Figure 2c; Figure S4). We found that the spatial patterns of inferred thickness for

409 the sub-landslides were similar to the full landslide with some differences in the thickness
410 magnitude suggesting it is appropriate to consider Boulder Creek either one large landslide
411 complex that has 5 active areas or 5 smaller landslides.

412 We also found that the inferred active thickness for the Boulder Creek landslide was
413 particularly irregular. While we expect areas that are not currently active to thin and even have
414 zero thickness in places, the active transport zone on Boulder Creek also contains thin and thick
415 patches (Figures 3c, 3f, and 3i). One explanation for this variability in the transport zone is that
416 there is a large channel network incised into the Boulder Creek landslide (Figure 3). In some
417 places the channel reaches depths of 15-20 meters (Figure S5). Since the thickness is measured
418 as the vertical distance from the ground surface to the inferred basal sliding surface, the predicted
419 thickness is expected to be low in places surrounding the channel if the channel depth is similar
420 to the landslide thickness. Our findings indicate that the channel has incised to depths that
421 approach the predicted sliding surface in several places (Figure S5).

422 Although we do not have borehole data to confirm our thickness estimates, we used the
423 topography to verify the inferred slip surface elevation in several cases. Figure 3e inset and
424 Figure S6 shows the landslide that has a clear headscarp that can be used to trace the sliding
425 surface underneath the ground surface. The extension of the headscarp slip surface under the
426 landslide provides confirmation that the inversion is approximating the slip surface elevation
427 correctly. Figure S7 shows another slow-moving landslide that has filled into a pre-existing
428 valley. Transects across this landslide show the ground surface of the filled-in valley and that the
429 slip surface has the shape of the pre-existing valley, providing additional confirmation that our
430 inversions are approximating the slip surfaces correctly. In addition, we compared our thickness
431 inversions to thickness estimates from lidar. Mackey and Roering (2011) used lidar to measure

432 the toe height at the channel interface for dozens of landslides in the Eel River catchment. Of
 433 those landslides, 10 can be used to make direct comparisons with our dataset. We found overall
 434 good agreement between the landslides toe thickness estimated from lidar and from our
 435 inversions (Figure 4).

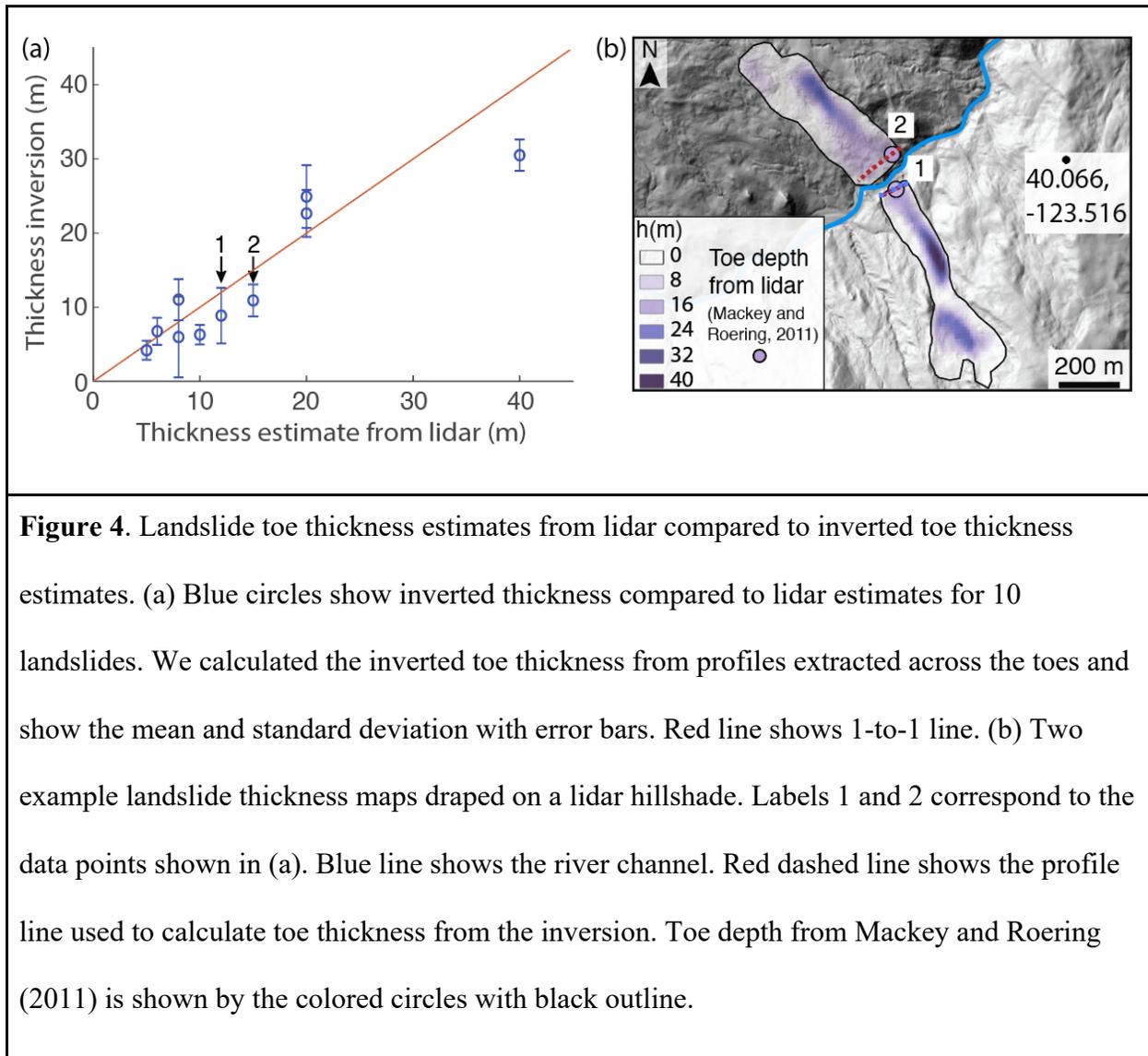


Figure 4. Landslide toe thickness estimates from lidar compared to inverted toe thickness estimates. (a) Blue circles show inverted thickness compared to lidar estimates for 10 landslides. We calculated the inverted toe thickness from profiles extracted across the toes and show the mean and standard deviation with error bars. Red line shows 1-to-1 line. (b) Two example landslide thickness maps draped on a lidar hillshade. Labels 1 and 2 correspond to the data points shown in (a). Blue line shows the river channel. Red dashed line shows the profile line used to calculate toe thickness from the inversion. Toe depth from Mackey and Roering (2011) is shown by the colored circles with black outline.

436

437 Using our thickness inversions for each landslide, we estimate that the individual
 438 landslide volumes range from 7.08×10^3 to 9.75×10^6 m³. We fit a power function to the
 439 volume-area to characterize the geometric scaling relations for these slow-moving landslides. We

440 also compared our inventory to a worldwide inventory of catastrophic soil, undifferentiated, and
441 bedrock landslides compiled by Larsen et al. (2010). We find that the slow-moving landslides in
442 the northern California Coast Ranges are larger in both area and volume than most soil
443 landslides, but smaller than the largest bedrock landslides around the world. The best fit volume-
444 area power function exponent (with 95% confidence) for the inventory was $\gamma = 1.324$ (1.206,
445 1.442) and has a R-square of 0.7869 (Figure 5). We report the scaling constants c_V and c_h in
446 Table S4. We observed an apparent break in the slope of the volume-area relation for the largest
447 landslides in our inventory with area $> 10^5$ m². To further investigate this break in slope, we also
448 fit volume-area scaling as a function of landslide type (slumps, earthflows, or complexes) and
449 find that the break in slope is primarily associated with the landslide complexes. By fitting a
450 power function to each landslide type, we find slumps $\gamma_S = 1.502$ (1.121, 1.883) with R-square:
451 0.6014, earthflows $\gamma_{Ef} = 1.453$ (1.075, 1.831) with R-square: 0.6016, and complexes $\gamma_C = 1.182$
452 (0.9739, 1.389) with R-square: 0.7402. Although these parameters are not statistically distinct at
453 the 95% confidence level, the fact that γ_S and γ_{Ef} overlap more with each other than with γ_C
454 supports the argument that the break in slope is likely related to landslide type. This likely
455 change in volume-area scaling for the landslide complexes shows that the largest landslides in
456 our inventory get larger in area but do not continue to get significantly deeper with increasing
457 area. In addition, we calculated the thickness-area scaling relations using the mean thickness to
458 represent each landslide. We compared these scaling relations to point based estimates (lidar)
459 and measurements (boreholes) of landslide thickness for slow-moving landslides in the northern
460 California Coast Ranges (Mackey and Roering, 2011) and the Reno River catchment, Apennines,
461 Italy (Simoni et al., 2013) (Figure 5). The best fit depth-area power function exponent (with 95%
462 confidence) for the inventory $\zeta = 0.3243$ (0.2063, 0.4424) with R-square: 0.1828, indicating a

463 weak increase in depth with area for the inventory as a whole (Figure 5). We also fit depth-area
 464 scaling as a function of landslide type and find slumps $\zeta_S = 0.5018$ (0.1211, 0.8826) with R-
 465 square: 0.1442, earthflows $\zeta_{Ef} = 0.453$ (0.07515, 0.8308) with R-square: 0.128, and for landslide
 466 complexes $\zeta_C = 0.1817$ (-0.02609, 0.3895) with R-square: 0.0631. Therefore, landslide thickness
 467 significantly increases with area for slumps and earthflows, but does not significantly vary with
 468 area for landslide complexes.

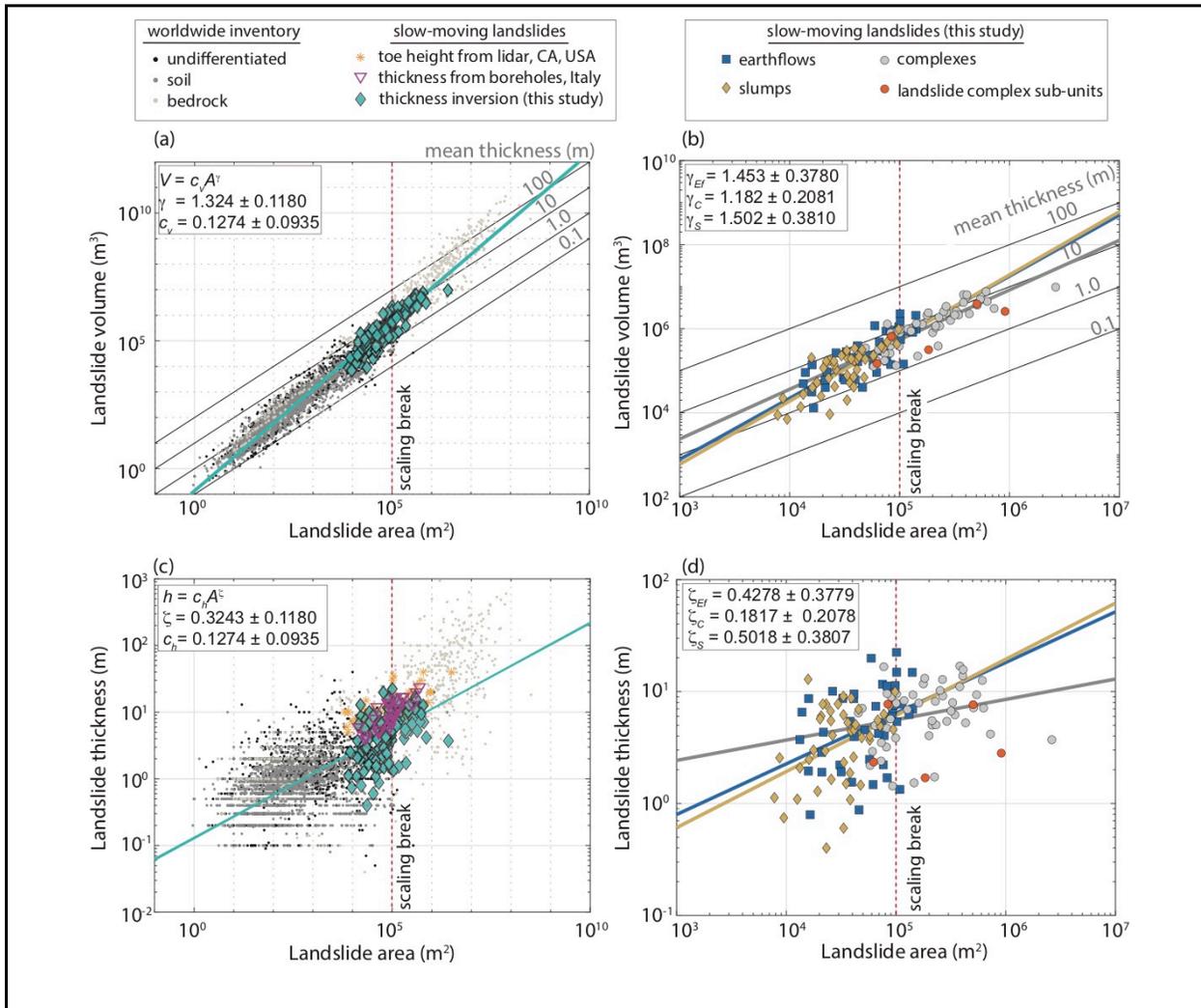


Figure 5. Landslide thickness, volume, and area geometric scaling relations. (a) Volume-area relations for our inventory and a worldwide inventory of soil, undifferentiated, and bedrock

landslides (Larsen et al. 2010). (b) Volume-area relation for slumps, earthflows, and landslide complexes. (a, b) Thin diagonal black lines show volume-area for various constant mean thicknesses. (c) Thickness-area relations for our inventory (mean thickness), the worldwide inventory (Larsen et al., 2010), and slow-moving landslides in the northern California Coast Ranges (Mackey and Roering, 2011) and the Apennine mountains, Italy (Simoni et al., 2013). (d) Landslide thickness-area relations by category. Orange circles in (b, d) correspond to the Boulder Creek landslide complex split into 5 smaller landslides (Figure S4). Red dashed vertical line shows an apparent break in scaling for the largest landslide complexes in our dataset.

469

470 3.3 Stress and Frictional Strength

471 We calculated the basal shear stress and effective normal stress under dry and saturated
472 conditions end members assuming nil cohesion (Figure 6; Table S2). We find that the average
473 shear stress for each landslide ranged from ~ 1.30 to ~ 102 kPa assuming dry conditions and ~ 1.51
474 to ~ 122 kPa assuming fully saturated conditions. The average effective normal stress ranged
475 from ~ 7.00 to ~ 373 kPa for dry conditions and ~ 3.98 to ~ 215 kPa for saturated conditions. The
476 inferred landslide friction angle ϕ ranged from ~ 7 to ~ 28 degrees for dry conditions and ~ 13 to
477 ~ 54 degrees for saturated conditions. We note that the inferred friction angle is larger when
478 saturated conditions are assumed due to higher pore-water pressure. i.e., the pore-water pressure
479 reduces the effective normal stress, so the inferred friction coefficient must be higher to maintain
480 a factor of safety equal to one (equation 9). Our inferred friction angles encompass friction angle
481 values measured in the laboratory for Franciscan *mélange* rocks and landslide material (Figure
482 6). We also analyzed the friction angle as a function of landslide size and found a weak

483 decreasing power-function relationship with increasing size, defined as $\phi = kL^r$, where k is the
 484 intercept, L is the landslide length, and r is the exponent. For dry conditions (with 95%
 485 confidence interval), $k = 56.55$ (39.37, 73.73) and $r = -0.2052$ (-0.2562, -0.1541) with R-square:
 486 0.3257. For saturated conditions (with 95% confidence interval), $k = 108.2$ (74.74, 141.8) and $r =$
 487 -0.2073 (-0.2593, -0.1553) with R-square: 0.3225. Figure 6c shows the mean friction angle for
 488 each landslide as a function of landslide length. The negative trend indicates that the largest
 489 landslides are weaker, on average, than smaller landslides.

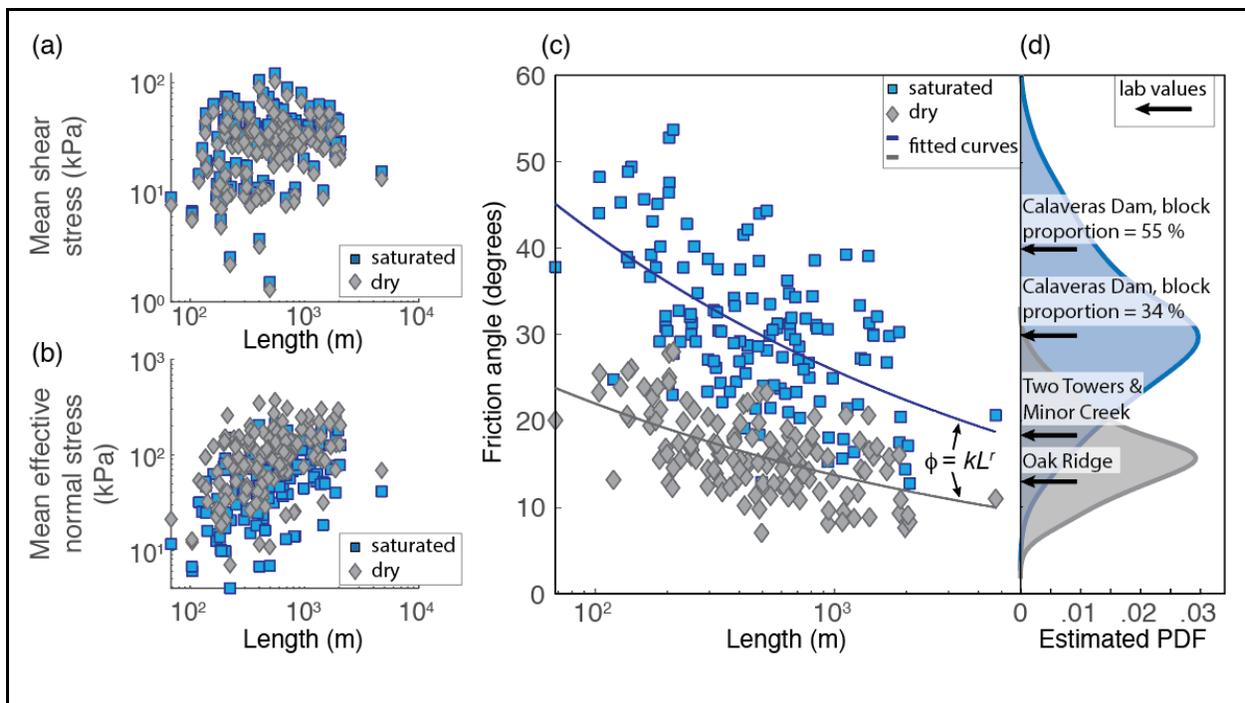


Figure 6. Inferred basal stresses and friction angle for dry and saturated end-members. (a) Mean shear stress, (b) mean effective normal stress, and (c) friction angle for each landslide. Solid line lines in (c) correspond to best-fit power function curves $\phi = kL^r$. For dry conditions, $k = 56.55$ and $r = -0.2052$ and for saturated conditions $k = 108.2$ and $r = -0.2073$. (d) Estimated probability density function for the full inventory. Black arrows show lab-based friction angle values for the Franciscan mélangé hosted Oak Ridge, Two Towers, and Minor Creek

landslides (Iverson & Major, 1987; Nereson & Finnegan, 2018; Schulz et al., 2018) and the Calaveras Dam, which is founded on Franciscan mélangé (Roadifer et al., 2009). The Calaveras Dam samples are plotted for two different block-in-matrix proportions.

490

491 **4 Discussion**

492 4.1 Landslide Kinematics

493 Our 3D UAVSAR velocity measurements reveal 134 active slow-moving landslides in
494 the northern California Coast Ranges moving at average rates from cm/yr to m/yr between 2016
495 and 2019. The 3D velocity data confirm that the motion of these landslides is generally in the
496 downslope direction and that there are segments with vertical uplift. Uplift tends to be at the
497 landslide toe due to the concave-up slip surface geometry, and the tendency for longitudinal
498 shortening in the direction of motion to occur at the toe. However, it is possible that a component
499 of uplift of landslide surfaces could also result from dilation or swelling (volumetric expansion),
500 but the magnitude is small, likely on the order of a few centimeters at most (Booth et al., 2020;
501 Delbridge et al., 2016; Iverson, 2005; Schulz et al., 2018). Including volume changes such as this
502 in the thickness inversion would refine its results, but the amount of dilation or compaction
503 occurring throughout an entire landslide and its temporal variation is generally unknown.

504 Our findings agree with previous work in this region that shows that these landslides
505 exhibit slow, spatially non-uniform downslope motion. Several of the landslides in our study
506 area (e.g., Boulder Creek) have been moving in this manner since at least 1944 (Bennett,
507 Roering, et al., 2016; Mackey & Roering, 2011). Our findings also show that pixel offset
508 tracking with very high resolution UAVSAR data is well-suited for monitoring landslides
509 moving at rates > 10 cm/yr. Some satellites acquire very high resolution SAR with Spotlight

510 modes, including German TerraSAR-X and Italian COSMO-SkyMed that could provide similar
511 measurements (e.g., Madson et al., 2019). We note that there are likely active landslides moving
512 < 10 cm/yr that were not observed by our pixel offset tracking measurements. In addition,
513 catastrophic landslides are not detected by our pixel offset tracking approach and would require
514 different multi-temporal analysis of SAR, optical, or DEMs to identify these types of landslides
515 (e.g., Intrieri et al., 2017; Jung & Yun, 2020). Our future work will revisit the time-dependent
516 motion and strain rates of each landslide in relation to ongoing climate shifts using InSAR and
517 pixel offset tracking techniques as we continue to collect UAVSAR data over many landslides in
518 California.

519

520 4.2 Landslide Geometry

521 Our study is the first (to our knowledge) to apply landslide thickness inversions to more
522 than a single landslide in a given region. Previous work (Booth et al., 2020; Booth, Lamb, et al.,
523 2013; Delbridge et al., 2016) has used the same approach to analyze individual landslides, but
524 these landslides occur in different regions and environmental conditions. Like these previous
525 studies, however, we found that the active landslide thickness is variable and that the slip
526 surfaces are rough and irregular in places. The non-uniform thickness and velocity of each
527 landslide results in a non-uniform sediment flux, which has implications for understanding
528 sediment motion along hillslopes (Booth et al., 2020; Guerriero et al., 2017). The shape of the
529 slip surface likely also impacts the landslide kinematics and groundwater flow (Coe et al., 2009;
530 Guerriero et al., 2014; Iverson & Major, 1987; Keefer & Johnson, 1983). Slip surfaces that are
531 bumpy and rough may create additional resisting stresses that act to prevent runaway
532 acceleration and permit long periods of slow landslide motion (Baum & Johnson, 1993; Booth et

533 al., 2018; Leshchinsky, 2019). Investigation of tectonic faults and glaciers also shows that slip
534 surface roughness is an important parameter that controls frictional strength (Brodsky et al.,
535 2016; Fang & Dunham, 2013; Meyer et al., 2018).

536 We found that the mean and maximum active thickness of the individual landslides
537 ranged from ~ 0.40 to ~ 22.4 m and ~ 2.2 to ~ 90 m, respectively. For simplicity we set $f = 1$
538 indicating that these landslides are effectively blocks sliding on a slope. Yet the borehole data
539 from the Two Towers landslide shows that $f \sim 0.96$. While changing f does not alter the spatial
540 pattern of thickness it does impact the magnitude of the thickness and therefore the volume.
541 Setting $f = 0.96$ would cause a 4% increase in the inferred thickness and volume of each
542 landslide ($h \sim 1/f$) (Table S3). Nonetheless, our findings indicate that most of the sliding surfaces
543 are deep-seated (median thickness for inventory = 5.5 m) and thus are expected to lie within the
544 unweathered Central Belt Franciscan *mélange* bedrock. Recent work by Hahm et al., (2019)
545 found that the Central Belt Franciscan *mélange* has a thin critical zone that is approximately ~ 3
546 m thick at the ridgetop. Therefore, the slow-moving landslides in the northern California Coast
547 Ranges can be classified as bedrock landslides. However, we note that the Central Belt
548 Franciscan *mélange* is a weak bedrock that has a strength more comparable to weathered rock or
549 soil (Iverson, 2000; Keefer & Johnson, 1983; Nereson et al., 2018; Schulz et al., 2018).

550 Using our landslide inventory, we developed new volume-area and depth-area geometric
551 scaling relations for slow-moving landslides. Geometric scaling relations are particularly useful
552 for slow-moving landslides because these landslides rarely (if ever) evacuate hillslopes, or create
553 clear scars or deposits that can be easily measured. As a result, most measurements of landslide
554 thickness come from isolated boreholes, which are logistically challenging and expensive to
555 install, and are difficult to extrapolate over an entire landslide. Our results provide best-fit power

556 function exponents ($\gamma = 1.324$, $\gamma_S = 1.502$, $\gamma_T = 1.453$, $\gamma_C = 1.182$) that are comparable to power
557 function exponents for bedrock and soil landslides (Guzzetti et al., 2009; Larsen et al., 2010).
558 Analysis of a worldwide landslide inventory by Larsen et al (2010) showed that soil landslides
559 had a $\gamma_{soil} \sim 1.1-1.3$, while bedrock landslides had $\gamma_{bedrock} \sim 1.3-1.6$. These findings indicate that
560 soil landslides tend to get larger by increasing their planform area rather than by increasing their
561 thickness (and planform area). Larsen et al. (2010) proposed that soil landslides do not continue
562 to get thicker with increasing area because they are limited by the maximum soil depth, which is
563 set by the environmental conditions, while bedrock landslides can extend to depths that far
564 exceed a typical soil cover. Our best-fit depth-area scaling power exponents ($\zeta = 0.3243$, $\zeta_S =$
565 0.5018 , $\zeta_{Ef} = 0.453$, $\zeta_C = 0.1817$) are comparable (with a wide range) to previously published
566 values for slow-moving landslides (Figure 4c). Simoni et al. (2013) reported depth-area scaling ζ
567 $= 0.44$ from borehole inclinometer data from 23 slow-moving landslides in the Apennine
568 Mountains, Italy. Handwerger et al. (2013) reported depth-area scaling $\zeta = 0.29$ derived from
569 lidar-based estimates of landslide toe thickness from 69 landslides in the Eel River catchment,
570 several of which are also analyzed in this study. Importantly, neither of these studies used large
571 inventories (> 100) or spatially extensive measurements of landslide thickness, which are
572 especially important for slow-moving landslides with variable thicknesses. We note that the large
573 range of scaling exponents suggests that scaling relations should be used with caution. Applying
574 an incorrect scaling exponent to estimate volume for landslides with unknown thickness can lead
575 to large errors in volume calculations (Larsen et al., 2010).

576 Our findings show that the slow-moving landslides located in the northern California
577 Coast Ranges have geometric scaling exponents that lie in between the soil and bedrock type
578 landslides, providing further evidence that these landslides are hosted in a weak rock that has

579 mechanical properties in between soil and competent rock (Nereson et al., 2018; Schulz et al.,
580 2018). However, examining the best-fit power function exponents by landslide type suggests that
581 slumps display close to self-similar scaling ($\gamma=1.5$), which is within the range of bedrock
582 landslides, earthflows display scaling intermediate to soil and bedrock landslides, and landslide
583 complexes display scaling most similar to soil landslides. Landslide complexes have similar
584 scaling relations to soil landslides because our inversions show that landslide complexes do not
585 continue to get deeper (on average) as they get larger in planform area. Figure 5 shows that the
586 landslide complexes with the largest areas display a scaling that tends to follow a constant mean
587 depth. This finding has two plausible explanations: 1) the mean landslide depth is limited by a
588 strong layer in the *mélange*, or 2) that landslide complexes are an amalgamation of multiple
589 smaller and shallower landslides. The second explanation provides a reason for why large
590 landslide complexes tend to have multiple kinematic units (e.g., Aryal et al., 2012; Hu et al.,
591 2020).

592 To further explore the hypothesis that landslide complexes are composed of multiple
593 smaller landslides, we delineated the Boulder Creek landslide complex into 5 smaller sub-
594 landslides and performed a thickness inversion for each sub-landslide (Figure S4). While the
595 thickness patterns are nearly identical to the thickness inversion for the full landslide complex,
596 the magnitude of the inferred thickness differs in some places, and the area of each landslide is
597 smaller, which places them into the space populated by mostly earthflows on the depth-area and
598 volume-area plots (orange circles in Figures 5b and 5d). It is also important to note that some of
599 these differences in the magnitude of the thickness estimate are due to differences in the pixel
600 resolution of the sub-landslides and the full landslide. For the full landslide complex we had to
601 downsample the grid to a 20 m pixel due to computational limitation. The Boulder Creek

602 landslide complex is the largest landslide in our dataset and is the only landslide that required
603 this downsampling. Nonetheless, mapping landslide complexes as one large landslide results in a
604 low mean thickness relative to the landslide area which affects the geometric scaling relations.
605 While more investigation is warranted, our thickness inversions have caused us to reevaluate
606 how we think about large landslide complexes.

607 Lastly, the inferred thickness of many of the slow-moving landslides, in particular the
608 landslide complexes, can be highly variable with deeper active zones and thinner or zero
609 thickness areas that are not currently moving (e.g., Figure 3c). It is important to note that these
610 irregular thickness patterns may not align with inferred thickness based on geomorphic or
611 structural interpretations. This discrepancy is likely related to the long-lasting geomorphic
612 imprint that slow-moving landslides have made on the landscape. Landslide surface morphology
613 may last for decades or longer after a landslide completely stops moving (e.g., Booth et al.,
614 2017), which can make it challenging to infer the active landslide thickness without kinematic
615 data. Although our approach is useful for identifying the currently active portions of landslides
616 and inferring their thickness based on volume conservation (with assumptions), it does not allow
617 us to infer the subsurface geometry of the often larger inactive landslide body. As a result, we
618 emphasize the need for more comparisons between ground- and remote sensing-based
619 investigation of landslide geometry. In particular, direct comparison between numerous ground-
620 based measurements from boreholes and structural mapping are needed to widely test the results
621 of our remote sensing approach. Nonetheless, we find our thickness inversions are producing
622 reasonable estimates of landslide thickness in the cases we were able to test (Figures 4, S6, and
623 S7).

624

625 4.3 Landslide Basal Stress and Strength

626 We calculated the shear (or driving) stress and effective normal stress for each landslide
627 under both dry and saturated (hydrostatic pore-water pressure) end member conditions. There is
628 a weak increasing relationship between landslide size and mean effective normal stress. Since the
629 mean shear stress must be balanced by resisting stress, defined here as the effective normal stress
630 times the friction coefficient, our findings imply that the larger landslides are weaker (i.e., lower
631 residual friction angles) than the smaller landslides. We hypothesize that larger landslides are
632 weaker than smaller landslides because of strength heterogeneity in the Franciscan *mélange*
633 bedrock and the increased likelihood of incorporating weak material within larger volumes.
634 Laboratory measurements of the strength of the Franciscan *mélange* rocks have shown that the
635 proportion of the blocks hosted in the argillaceous matrix controls the overall rock strength
636 (Roadifer et al., 2009). This implies that larger landslides may have a decreased proportion of
637 blocks and are therefore controlled by the weak argillaceous matrix. Scale-dependent strength
638 has also been observed along other landslides and faults. Brodsky et al. (2016) suggested that
639 faults are weaker at large spatial scales because they encompass larger weak zones. A recent
640 study by Bunn et al., (2020) found that the inferred shear strength of landslides decreases with
641 increasing landslide size. They proposed that smaller landslides were stronger because they
642 occur in cemented cohesive materials and larger landslides were in a residual state. Although we
643 assumed nil cohesion to back-calculate the residual frictional strength of the active landslides, it
644 is likely that cohesion is important in controlling the initial landslide failure due the high clay
645 content of the Central Belt Franciscan *mélange* (e.g., Milledge et al., 2014).

646 We also compared our inferred friction angles to lab-based measurements for landslide
647 material and rocks in the Franciscan *mélange* (Figure 6). The Franciscan *mélange* rock has a

648 friction angle that depends on the block-in-matrix proportion (Roadifer et al., 2009). The friction
649 angle of the rock was 30 degrees with 34% block proportion and 40 degrees with 55% block
650 proportion. The Two Towers landslide (Schulz et al., 2018) and Minor Creek landslide (Iverson,
651 2000), both located in the northern California Coast Ranges have drained residual friction angles
652 of 18.75 and 18 degrees, respectively. The Oak Ridge landslide in the central California Coast
653 ranges has a residual friction angle between 12-14 degrees (Nereson et al., 2018). Many of our
654 inferred friction angles are larger than the values measured in the laboratory for the landslides.
655 These differences may be partially attributed to differences in laboratory- and field-scale
656 measurements (Bunn et al., 2020; Marone, 1998; Van Asch et al., 2007).

657 Our inversion produces a widespread range of friction angle values from 7 to 53 degrees
658 with a median of 30 degrees for saturated conditions and 16 degrees for dry conditions. Higher
659 apparent friction angles are required for saturated conditions to offset the weakening induced by
660 the elevated pore-water pressure in order to maintain equilibrium conditions (equation 9). Due to
661 the high seasonal rainfall in the northern California Coast Ranges, the slow-moving landslides in
662 California are typically saturated during the wet season and partially saturated during the dry
663 season (Hahm et al., 2019; Iverson & Major, 1987; Schulz et al., 2018). Therefore, the true mean
664 landslide-scale friction angle values likely lie somewhere between our inferred values for
665 saturated and dry conditions. Lastly, although our inferred friction angle values encompass the
666 lab-based values from similar landslides and rocks, the large spread makes it difficult to identify
667 a single representative value for slow-moving landslides in the northern California Coast Ranges.
668 This finding further highlights the heterogeneous nature of the Central Belt Franciscan mélange
669 lithologic unit.

670

671 4.4 What Controls the Size of Slow-moving Landslides?

672 Landslide size is set by the landslide mechanical properties, slope geometry, and
673 environmental conditions. For most landslides, the maximum size is typically limited to the
674 maximum hillslope size, such that the landslide length does not exceed the hillslope length. The
675 landslide thickness is typically set by the location of a weak layer beneath the ground surface, or
676 at a depth where there are changes in strength and permeability, such as the soil to bedrock
677 transition or critical zone to unweathered bedrock transition (Booth, Roering, et al., 2013; Larsen
678 et al., 2010; Milledge et al., 2014). Using a 3D slope stability model for shallow soil landslides
679 that accounts for the forces acting on the landslide basal slip surface, lateral margins, and
680 passive/active wedges at the toe/head, Milledge et al. (2014) found that the critical area and
681 depth that can fail as a landslide depends on the topography, pore-water pressure, and landslide
682 material properties, including density, cohesion, and friction angle. In that model the pore-water
683 pressure plays a fundamental role in determining the critical landslide size and failure depth,
684 such that higher pore-water pressures decrease the critical size required for failure. Large
685 landslides therefore occur when high pore pressures are reached over a correspondingly large
686 spatial area. At our northern California Coast Range study site, the relatively thin, but laterally
687 extensive critical zone that is often saturated during the wet season (Hahm et al., 2019), may
688 promote laterally extensive landslides by elevating the water table height simultaneously over
689 large areas. Milledge et al. (2014)'s model also predicts that landslide thickness should increase
690 as the square root of the landslide area and that the failure depth sets the minimum landslide area.
691 Our best-fit thickness-area scaling exponents for slumps and earthflows are close to a square root
692 scaling (exponents ~ 0.5 with large 95th confidence intervals). Our results also suggest that the
693 landslide thickness controls the minimum area, but does not bound its maximum size. Instead,

694 slow-moving landslides can continue to grow in area by becoming a landslide complex
695 consisting of multiple, connected, sub-landslides without becoming significantly deeper on
696 average. Large landslide complexes can occupy multiple hillslopes, and fill valleys and
697 catchments such that their size may exceed the typical hillslope size. Thus, it seems that the
698 catchment size sets the maximum area for slow-moving landslides. Our thickness inversion
699 results also indicate that large landslides are weaker than small landslides. This finding may
700 indicate that large landslides become large by incorporating weak material. It is possible that the
701 largest landslides grow over time and take decades to develop (e.g., Mackey & Roering, 2011).
702 As many of our landslide complexes seem to be composed of several smaller sub-landslides or
703 kinematic zones, it is possible that these features have connected through time as slip surfaces
704 propagate along the slope.

705

706 **5 Conclusions**

707 We measured the 3D surface velocity of more than one hundred slow-moving landslides
708 in the northern California Coast Ranges with data from the NASA/JPL UAVSAR. We used
709 volume conservation techniques to infer the active thickness, volume, stress, and strength of each
710 landslide. The thickness of each landslide is variable and can range from zero to tens of meters
711 sometimes resulting in an irregular slip surface geometry. Volume-area geometric scaling
712 relations suggest that these landslides have similarities to both soil and bedrock landslides
713 around the world. Although their failure planes are likely hosted in unweathered bedrock, their
714 depth seems to be limited, producing a scaling similar to soil landslides for the largest landslide
715 complexes. The inferred residual friction angles are also scale-dependent, like faults, such that
716 large landslides tend to be weaker than small landslides. This decrease in inferred friction angle

717 with landslide size is likely because larger landslides are composed of larger proportions of weak
718 bedrock. Our study represents the first to use the conservation of volume approach for numerous
719 landslides occurring under the same environmental conditions. Our results provide key insights
720 into the subsurface geometry and stresses that control the behavior of slow-moving landslides.
721 Our work shows how state-of-the-art remote sensing techniques can be used to better understand
722 landslide processes for hazards and to quantify their contribution to landscape evolution.

723

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733

734 **Data Availability**

735 Landslide geometry data used in this study are listed in the references: Larsen et al.,
736 (2010), Mackey and Roering, (2011), Simoni et al., (2013) and are included in the figures.
737 Borehole thickness data at the Two Towers landslide is in reference: Schulz et al. (2018). Lidar
738 digital elevation models are provided by OpenTopography and may be downloaded online
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744 To acquire these data, proposals may be submitted to the DLR online ([https://tandemx-](https://tandemx-science.dlr.de/)
745 [science.dlr.de/](https://tandemx-science.dlr.de/)). NASA/JPL UAVSAR data used in this study are freely available and may be
746 downloaded through their website (<https://uavsar.jpl.nasa.gov/>).

747

748 **References**

- 749 Aryal, A., Brooks, B. A., Reid, M. E., Bawden, G. W., & Pawlak, G. R. (2012). Displacement fields from
750 point cloud data: Application of particle imaging velocimetry to landslide geodesy. *Journal of*
751 *Geophysical Research: Earth Surface* (2003–2012), 117(F1).
752 <https://doi.org/10.1029/2011JF002161>
- 753 Aryal, A., Brooks, B., & Reid, M. E. (2015). Landslide subsurface slip geometry inferred from 3D surface
754 displacement fields. *Geophysical Research Letters*. 42(5), 1411-1417.
755 <https://doi.org/10.1002/2014GL062688>
- 756 Aster, R. C., Borchers, B., & Thurber, C. H. (2013). Chapter Four - Tikhonov Regularization. In R. C.
757 Aster, B. Borchers, & C. H. Thurber (Eds.), *Parameter Estimation and Inverse Problems (Second*
758 *Edition)* (Second Edition, pp. 93–127). Boston: Academic Press. [https://doi.org/10.1016/B978-0-](https://doi.org/10.1016/B978-0-12-385048-5.00004-5)
759 [12-385048-5.00004-5](https://doi.org/10.1016/B978-0-12-385048-5.00004-5)
- 760 Baum, R. L., & Johnson, A. M. (1993). Steady movement of landslides in fine-grained soils: A model for
761 sliding over an irregular slip surface. *US Geological Survey Bulletin (USA)*.
- 762 Bennett, G. L., Roering, J. J., Mackey, B. H., Handwerger, A. L., Schmidt, D. A., & Guillod, B. P.
763 (2016). Historic drought puts the brakes on earthflows in Northern California. *Geophysical*
764 *Research Letters*, 43(11), 5725–5731. <https://doi.org/10.1002/2016GL068378>

- 765 Bennett, G. L., Miller, S. R., Roering, J. J., & Schmidt, D. A. (2016). Landslides, threshold slopes, and
766 the survival of relict terrain in the wake of the Mendocino Triple Junction. *Geology*, *44*(5), 363–
767 366. <https://doi.org/10.1130/G37530.1>
- 768 Bessette-Kirton, E. K., Coe, J. A., & Zhou, W. (2018). Using stereo satellite imagery to account for
769 ablation, entrainment, and compaction in volume calculations for rock avalanches on glaciers:
770 Application to the 2016 Lamplugh rock avalanche in Glacier Bay National Park, Alaska. *Journal*
771 *of Geophysical Research: Earth Surface*, *123*(4), 622–641. <https://doi.org/10.1002/2017JF004512>
- 772 Booth, A. M., Lamb, M. P., Avouac, J.-P., & Delacourt, C. (2013). Landslide velocity, thickness, and
773 rheology from remote sensing: La Clapière landslide, France. *Geophysical Research Letters*,
774 *40*(16), 4299–4304. <https://doi.org/10.1002/grl.50828>
- 775 Booth, A. M., Roering, J. J., & Rempel, A. W. (2013). Topographic signatures and a general transport law
776 for deep-seated landslides in a landscape evolution model. *Journal of Geophysical Research:*
777 *Earth Surface*, *118*(2), 603–624. <https://doi.org/10.1002/jgrf.20051>
- 778 Booth, A. M., LaHusen, S. R., Duvall, A. R., & Montgomery, D. R. (2017). Holocene history of deep-
779 seated landsliding in the North Fork Stillaguamish River valley from surface roughness analysis,
780 radiocarbon dating, and numerical landscape evolution modeling. *Journal of Geophysical*
781 *Research: Earth Surface*, *122*(2), 456–472. <https://doi.org/10.1002/2016JF003934>
- 782 Booth, A. M., McCarley, J., Hinkle, J., Shaw, S., Ampuero, J.-P., & Lamb, M. P. (2018). Transient
783 Reactivation of a Deep-Seated Landslide by Undrained Loading Captured With Repeat Airborne
784 and Terrestrial Lidar. *Geophysical Research Letters*, *45*(10), 4841–4850.
785 <https://doi.org/10.1029/2018GL077812>
- 786 Booth, A. M., McCarley, J. C., & Nelson, J. (2020). Multi-year, three-dimensional landslide surface
787 deformation from repeat lidar and response to precipitation: Mill Gulch earthflow, California.
788 *Landslides*, 1–14. <https://doi.org/10.1007/s10346-020-01364-z>
- 789 Brodsky, E. E., Kirkpatrick, J. D., & Candela, T. (2016). Constraints from fault roughness on the scale-
790 dependent strength of rocks. *Geology*, *44*(1), 19–22. <https://doi.org/10.1130/G37206.1>

- 791 Bunn, M., Leshchinsky, B., & Olsen, M. J. (2020). Geologic Trends in Shear Strength Properties Inferred
792 through Three-Dimensional Back-Analysis of Landslide Inventories. *Journal of Geophysical*
793 *Research: Earth Surface*, e2019JF005461. <https://doi.org/10.1029/2019JF005461>
- 794 Cerovski-Darriau, C., & Roering, J. J. (2016). Influence of anthropogenic land-use change on hillslope
795 erosion in the Waipaoa River Basin, New Zealand. *Earth Surface Processes and Landforms*,
796 *41*(15), 2167–2176. <https://doi.org/10.1002/esp.3969>
- 797 Coe, J. A., McKenna, J. P., Godt, J. W., & Baum, R. L. (2009). Basal-topographic control of stationary
798 ponds on a continuously moving landslide. *Earth Surface Processes and Landforms*, *34*(2), 264–
799 279. <https://doi.org/10.1002/esp.1721>
- 800 Cruden, D. M., & Varnes, D. J. (1996). Landslides: Investigation and Mitigation. Chapter 3-Landslide
801 types and processes. *Transportation Research Board Special Report*, (247).
- 802 Delbridge, B. G., Bürgmann, R., Fielding, E., Hensley, S., & Schulz, W. H. (2016). Three-dimensional
803 surface deformation derived from airborne interferometric UAVSAR: Application to the
804 Slumgullion Landslide. *Journal of Geophysical Research: Solid Earth*, *121*(5), 3951–3977.
805 <https://doi.org/10.1002/2015JB012559>
- 806 Fang, Z., & Dunham, E. M. (2013). Additional shear resistance from fault roughness and stress levels on
807 geometrically complex faults. *Journal of Geophysical Research: Solid Earth*, *118*(7), 3642–3654.
808 <https://doi.org/10.1002/jgrb.50262>
- 809 Fialko, Y., Simons, M., & Agnew, D. (2001). The complete (3-D) surface displacement field in the
810 epicentral area of the 1999 Mw7. 1 Hector Mine earthquake, California, from space geodetic
811 observations. *Geophysical Research Letters*, *28*(16), 3063–3066.
812 <https://doi.org/10.1029/2001GL013174>
- 813 Fielding, E. J., Liu, Z., Stephenson, O. L., Zhong, M., Liang, C., Moore, A., et al. (2020). Surface
814 Deformation Related to the 2019 M w 7.1 and 6.4 Ridgecrest Earthquakes in California from
815 GPS, SAR Interferometry, and SAR Pixel Offsets. *Seismological Research Letters*.
816 <https://doi.org/10.1785/0220190302>

- 817 Grant, M., & Boyd, S. (2014). *CVX: Matlab Software for Disciplined Convex Programming, version 2.1*.
818 Retrieved from <http://cvxr.com/cvx>
- 819 Guerriero, L., Coe, J. A., Revellino, P., Grelle, G., Pinto, F., & Guadagno, F. M. (2014). Influence of slip-
820 surface geometry on earth-flow deformation, Montaguto earth flow, southern Italy.
821 *Geomorphology*, 219(0), 285 – 305. <http://dx.doi.org/10.1016/j.geomorph.2014.04.039>
- 822 Guerriero, L., Bertello, L., Cardozo, N., Berti, M., Grelle, G., & Revellino, P. (2017). Unsteady sediment
823 discharge in earth flows: A case study from the Mount Pizzuto earth flow, southern Italy.
824 *Geomorphology*, 295, 260–284. <https://doi.org/10.1016/j.geomorph.2017.07.011>
- 825 Guzzetti, F., Ardizzone, F., Cardinali, M., Rossi, M., & Valigi, D. (2009). Landslide volumes and
826 landslide mobilization rates in Umbria, central Italy. *Earth and Planetary Science Letters*, 279(3–
827 4), 222–229. <https://doi.org/10.1016/j.epsl.2009.01.005>
- 828 Hahm, W. J., Rempe, D. M., Dralle, D. N., Dawson, T. E., Lovill, S. M., Bryk, A. B., et al. (2019).
829 Lithologically controlled subsurface critical zone thickness and water storage capacity determine
830 regional plant community composition. *Water Resources Research*, 55(4), 3028–3055.
831 <https://doi.org/10.1029/2018WR023760>
- 832 Handwerger, A. L., Roering, J. J., & Schmidt, D. A. (2013). Controls on the seasonal deformation of
833 slow-moving landslides. *Earth and Planetary Science Letters*, 377, 239–247.
834 <https://doi.org/10.1016/j.epsl.2013.06.047>
- 835 Handwerger, A. L., Roering, J. J., Schmidt, D. A., & Rempel, A. W. (2015). Kinematics of earthflows in
836 the Northern California Coast Ranges using satellite interferometry. *Geomorphology*, 246, 321–
837 333. <https://doi.org/10.1016/j.geomorph.2015.06.003>
- 838 Handwerger, A. L., Huang, M.-H., Fielding, E. J., Booth, A. M., & Bürgmann, R. (2019). A shift from
839 drought to extreme rainfall drives a stable landslide to catastrophic failure. *Scientific Reports*,
840 9(1), 1569. <https://doi.org/10.1038/s41598-018-38300-0>
- 841 Handwerger, A. L., Fielding, E. J., Huang, M.-H., Bennett, G. L., Liang, C., & Schulz, W. H. (2019).
842 Widespread initiation, reactivation, and acceleration of landslides in the northern California Coast

- 843 Ranges due to extreme rainfall. *Journal of Geophysical Research: Earth Surface*, 124(7), 1782–
844 1797. <https://doi.org/10.1029/2019JF005035>
- 845 Hu, X., Bürgmann, R., Schulz, W. H., & Fielding, E. J. (2020). Four-dimensional surface motions of the
846 Slumgullion landslide and quantification of hydrometeorological forcing. *Nature*
847 *Communications*, 11(1), 1–9. <https://doi.org/10.1038/s41467-020-16617-7>
- 848 Hungr, O. (1987). An extension of Bishop’s simplified method of slope stability analysis to three
849 dimensions. *Geotechnique*, 37(1), 113–117. <https://doi.org/10.1680/geot.1987.37.1.113>
- 850 Hungr, O., Salgado, F., & Byrne, P. (1989). Evaluation of a three-dimensional method of slope stability
851 analysis. *Canadian Geotechnical Journal*, 26(4), 679–686. <https://doi.org/10.1139/t89-079>
- 852 Hungr, O., Leroueil, S., & Picarelli, L. (2014). The Varnes classification of landslide types, an update.
853 *Landslides*, 11(2), 167–194. <https://doi.org/10.1007/s10346-013-0436-y>
- 854 Intrieri, E., Raspini, F., Fumagalli, A., Lu, P., Del Conte, S., Farina, P., et al. (2017). The Maoxian
855 landslide as seen from space: detecting precursors of failure with Sentinel-1 data. *Landslides*, 1–
856 11. <https://doi.org/10.1007/s10346-017-0915-7>
- 857 Iverson, R. M. (2000). Landslide triggering by rain infiltration. *Water Resources Research*, 36(7), 1897–
858 1910. <https://doi.org/10.1029/2000WR900090>
- 859 Iverson, R. M. (2005). Regulation of landslide motion by dilatancy and pore pressure feedback. *Journal*
860 *of Geophysical Research: Earth Surface*, 110(F2). <https://doi.org/10.1029/2004JF000268>
- 861 Iverson, R. M., & Major, J. J. (1987). Rainfall, ground-water flow, and seasonal movement at Minor
862 Creek landslide, northwestern California: Physical interpretation of empirical relations.
863 *Geological Society of America Bulletin*, 99(4), 579–594. [https://doi.org/10.1130/0016-
864 7606\(1987\)99<579:RGFASM>2.0.CO;2](https://doi.org/10.1130/0016-7606(1987)99<579:RGFASM>2.0.CO;2)
- 865 Jayko, A., Blake, M., McLaughlin, R., Ohlin, H., Ellen, S., & Kelsey, H. (1989). Reconnaissance
866 Geologic Map of the Covelo 30-by 60-Minute Quadrangle. *Northern California: US Geological*
867 *Survey Miscellaneous Field Investigation Map MF-2001, Scale, 1(100), 000.*
868 <https://doi.org/10.3133/mf2001>

- 869 Jennings, C. W., Strand, R. G., & Rogers, T. H. (1977). Geologic map of California: California Division
870 of Mines and Geology, scale 1:750,000.
- 871 Jung, J., & Yun, S.-H. (2020). Evaluation of Coherent and Incoherent Landslide Detection Methods
872 Based on Synthetic Aperture Radar for Rapid Response: A Case Study for the 2018 Hokkaido
873 Landslides. *Remote Sensing*, *12*(2), 265. <https://doi.org/10.3390/rs12020265>
- 874 Kasper, van W., John A, S., William, N., & Luis, T. (2002). Data and model uncertainty estimation for
875 linear inversion. *Geophysical Journal International*, *149*(3), 625–632.
876 <https://doi.org/10.1046/j.1365-246X.2002.01660.x>
- 877 Keefer, D. K., & Johnson, A. M. (1983). Earth flows: Morphology, mobilization, and movement.
878 Washington: United States Government Printing Office. No 1264. <https://doi.org/10.3133/pp1264>
- 879 Kelsey, H. M. (1978). Earthflows in Franciscan melange, Van Duzen River basin, California. *Geology*,
880 *6*(6), 361–364. [https://doi.org/10.1130/0091-7613\(1978\)6<361:EIFMVD>2.0.CO;2](https://doi.org/10.1130/0091-7613(1978)6<361:EIFMVD>2.0.CO;2)
- 881 Korup, O., Clague, J. J., Hermanns, R. L., Hewitt, K., Strom, A. L., & Weidinger, J. T. (2007). Giant
882 landslides, topography, and erosion. *Earth and Planetary Science Letters*, *261*(3–4), 578–589.
883 <https://doi.org/10.1016/j.epsl.2007.07.025>
- 884 Lacroix, P., Dehecq, A., & Taipe, E. (2020). Irrigation-triggered landslides in a Peruvian desert caused by
885 modern intensive farming. *Nature Geoscience*, *13*(1), 56–60. [https://doi.org/10.1038/s41561-019-](https://doi.org/10.1038/s41561-019-0500-x)
886 [0500-x](https://doi.org/10.1038/s41561-019-0500-x)
- 887 Lacroix, P., Handwerger, A. L., & Bièvre, G. (2020). Life and death of slow-moving landslides. *Nature*
888 *Reviews Earth & Environment*, 1–16. <https://doi.org/10.1038/s43017-020-0072-8>
- 889 Larsen, I. J., Montgomery, D. R., & Korup, O. (2010). Landslide erosion controlled by hillslope material.
890 *Nature Geoscience*, *3*(4), 247–251. <https://doi.org/10.1038/ngeo776>
- 891 Legros, F. (2002). The mobility of long-runout landslides. *Engineering Geology*, *63*(3–4), 301–331.
892 [https://doi.org/10.1016/S0013-7952\(01\)00090-4](https://doi.org/10.1016/S0013-7952(01)00090-4)
- 893 Leshchinsky, B. (2019). Quantifying the influence of failure surface asperities on the basal shear
894 resistance of translational landslides. *Landslides*, *16*(7), 1375–1383.

- 895 <https://doi.org/10.1007/s10346-019-01185-9>
- 896 Mackey, B. H., & Roering, J. J. (2011). Sediment yield, spatial characteristics, and the long-term
897 evolution of active earthflows determined from airborne LiDAR and historical aerial
898 photographs, Eel River, California. *Geological Society of America Bulletin*, *123*(7–8), 1560–
899 1576. <https://doi.org/10.1130/B30306.1>
- 900 Mackey, B. H., Roering, J. J., & McKean, J. (2009). Long-term kinematics and sediment flux of an active
901 earthflow, Eel River, California. *Geology*, *37*(9), 803–806. <https://doi.org/10.1130/G30136A.1>
- 902 Madson, A., Fielding, E., Sheng, Y., & Cavanaugh, K. (2019). High-resolution spaceborne, airborne and
903 in situ landslide kinematic measurements of the slumgullion landslide in Southwest Colorado.
904 *Remote Sensing*, *11*(3), 265. <https://doi.org/10.3390/rs11030265>
- 905 Malet, J.-P., Maquaire, O., & Calais, E. (2002). The use of Global Positioning System techniques for the
906 continuous monitoring of landslides: application to the Super-Sauze earthflow (Alpes-de-Haute-
907 Provence, France). *Geomorphology*, *43*(1–2), 33–54. [https://doi.org/10.1016/S0169-](https://doi.org/10.1016/S0169-555X(01)00098-8)
908 [555X\(01\)00098-8](https://doi.org/10.1016/S0169-555X(01)00098-8)
- 909 Marone, C. (1998). Laboratory-derived friction laws and their application to seismic faulting. *Annual*
910 *Review of Earth and Planetary Sciences*, *26*(1), 643–696.
911 <https://doi.org/10.1146/annurev.earth.26.1.643>
- 912 McLaughlin, R. J., Kling, S. A., Poore, R. Z., McDougall, K., & Beutner, E. C. (1982). Post–middle
913 Miocene accretion of Franciscan rocks, northwestern California. *Geological Society of America*
914 *Bulletin*, *93*(7), 595–605. [https://doi.org/10.1130/0016-7606\(1982\)93<595:PMAOFR>2.0.CO;2](https://doi.org/10.1130/0016-7606(1982)93<595:PMAOFR>2.0.CO;2)
- 915 McLaughlin, R. J., Blake, S., Jayko, M., Irwin, A., Aalto, W., Carver, K., et al. (2000). Geologic map of
916 the Cape Mendocino, Eureka, Garberville, and southwestern part of the Hayfork 30 X 60
917 Quadrangles and Adjacent Offshore Area, Northern California.
- 918 Meyer, C. R., Downey, A. S., & Rempel, A. W. (2018). Freeze-on limits bed strength beneath sliding
919 glaciers. *Nature Communications*, *9*(1), 1–6. <https://doi.org/10.1038/s41467-018-05716-1>
- 920 Michel, J., Dario, C., Marc-Henri, D., Thierry, O., Marina, P. I., & Benjamin, R. (2020). A review of

- 921 methods used to estimate initial landslide failure surface depths and volumes. *Engineering*
922 *Geology*, 267, 105478. <https://doi.org/10.1016/j.enggeo.2020.105478>
- 923 Milledge, D. G., Bellugi, D., McKean, J. A., Densmore, A. L., & Dietrich, W. E. (2014). A
924 multidimensional stability model for predicting shallow landslide size and shape across
925 landscapes. *Journal of Geophysical Research: Earth Surface*, 119(11), 2481–2504.
926 <https://doi.org/10.1002/2014JF003135>
- 927 Nereson, A. L., Davila Olivera, S., & Finnegan, N. J. (2018). Field and Remote-Sensing Evidence for
928 Hydro-mechanical Isolation of a Long-Lived Earthflow in Central California. *Geophysical*
929 *Research Letters*, 45(18), 9672–9680. <https://doi.org/10.1029/2018GL079430>
- 930 Nereson, A. L., & Finnegan, N. J. (2018). Drivers of earthflow motion revealed by an 80 yr record of
931 displacement from Oak Ridge earthflow, Diablo Range, California, USA. *Geological Society of*
932 *America Bulletin*. 131(3-4), 389-402. <https://doi.org/10.1130/B32020.1>
- 933 Pathier, E., Fielding, E. J., Wright, T. J., Walker, R., Parsons, B. E., & Hensley, S. (2006). Displacement
934 field and slip distribution of the 2005 Kashmir earthquake from SAR imagery. *Geophysical*
935 *Research Letters*, 33(20). <https://doi.org/10.1029/2006GL027193>
- 936 Roadifer, J. W., Forrest, M. P., & Lindquist, E. S. (2009). Evaluation of shear strength of melange
937 foundation at Calaveras Dam. Proceedings of U. S. Society for Dams, Annual Meeting and
938 Conference, 29th, on "Managing Our Water Retention Systems, ", 507–521.
- 939 Roering, J.J (2012). Eel River, CA: Landsliding and the Evolution of Mountainous Landscapes in
940 collaboration with National Center for Airborne Laser Mapping (NCALM), distributed by
941 OpenTopography. <https://doi.org/10.5069/G9XS5S9P>
- 942 Roering, J. J., Stimely, L. L., Mackey, B. H., & Schmidt, D. A. (2009). Using DInSAR, airborne LiDAR,
943 and archival air photos to quantify landsliding and sediment transport. *Geophysical Research*
944 *Letters*, 36(19). <https://doi.org/10.1029/2009GL040374>
- 945 Roering, J. J., Mackey, B. H., Handwerger, A. L., Booth, A. M., Schmidt, D. A., Bennett, G. L., &
946 Cerovski-Darriau, C. (2015). Beyond the angle of repose: A review and synthesis of landslide

- 947 processes in response to rapid uplift, Eel River, Northern California. *Geomorphology*, 236, 109–
948 131. <https://doi.org/10.1016/j.geomorph.2015.02.013>
- 949 Rosen, P. A., Gurrola, E., Sacco, G. F., & Zebker, H. (2012). The InSAR scientific computing
950 environment. In *Synthetic Aperture Radar, 2012. EUSAR. 9th European Conference on* (pp. 730–
951 733). VDE.
- 952 Rutter, E., & Green, S. (2011). Quantifying creep behaviour of clay-bearing rocks below the critical stress
953 state for rapid failure: Mam Tor landslide, Derbyshire, England. *Journal of the Geological*
954 *Society*, 168(2), 359–372. <https://doi.org/10.1144/0016-76492010-133>
- 955 Schulz, W. H., Coe, J. A., Ricci, P. P., Smoczyk, G. M., Shurtleff, B. L., & Panosky, J. (2017). Landslide
956 kinematics and their potential controls from hourly to decadal timescales: Insights from
957 integrating ground-based InSAR measurements with structural maps and long-term monitoring
958 data. *Geomorphology*, 285, 121–136. <https://doi.org/10.1016/j.geomorph.2017.02.011>
- 959 Schulz, W. H., Smith, J. B., Wang, G., Jiang, Y., & Roering, J. J. (2018). Clayey landslide initiation and
960 acceleration strongly modulated by soil swelling. *Geophysical Research Letters*, 45(4), 1888–
961 1896. <https://doi.org/10.1002/2017GL076807>
- 962 Simoni, A., Ponza, A., Picotti, V., Berti, M., & Dinelli, E. (2013). Earthflow sediment production and
963 Holocene sediment record in a large Apennine catchment. *Geomorphology*, 188, 42–53.
964 <https://doi.org/10.1016/j.geomorph.2012.12.006>
- 965 Stumpf, A., Malet, J.-P., & Delacourt, C. (2017). Correlation of satellite image time-series for the
966 detection and monitoring of slow-moving landslides. *Remote Sensing of Environment*, 189, 40–
967 55. <https://doi.org/10.1016/j.rse.2016.11.007>
- 968 Travelletti, J., & Malet, J.-P. (2012). Characterization of the 3D geometry of flow-like landslides: A
969 methodology based on the integration of heterogeneous multi-source data. *Engineering Geology*,
970 128, 30–48. <https://doi.org/10.1016/j.enggeo.2011.05.003>
- 971 Travelletti, J., Malet, J.-P., & Delacourt, C. (2014). Image-based correlation of Laser Scanning point
972 cloud time series for landslide monitoring. *International Journal of Applied Earth Observation*

- 973 *and Geoinformation*, 32, 1–18. <https://doi.org/10.1016/j.jag.2014.03.022>
- 974 Van Asch, T. W., Van Beek, L., & Bogaard, T. (2007). Problems in predicting the mobility of slow-
975 moving landslides. *Engineering Geology*, 91(1), 46–55.
976 <https://doi.org/10.1016/j.enggeo.2006.12.012>
- 977 Warrick, J. A., Ritchie, A. C., Schmidt, K. M., Reid, M. E., & Logan, J. (2019). Characterizing the
978 catastrophic 2017 Mud Creek landslide, California, using repeat structure-from-motion (SfM)
979 photogrammetry. *Landslides*, 1–19. <https://doi.org/10.1007/s10346-019-01160-4>
- 980 Wartman, J., Montgomery, D. R., Anderson, S. A., Keaton, J. R., Benoît, J., dela Chapelle, J., & Gilbert,
981 R. (2016). The 22 March 2014 Oso landslide, Washington, USA. *Geomorphology*, 253, 275–288.
982 <https://doi.org/10.1016/j.geomorph.2015.10.022>