

Inferring the Subsurface Geometry, Stress, and Strength of Slow-moving Landslides using 3D Velocity Measurements from the NASA/JPL UAVSAR

A. L. Handwerger^{1,2}, A. M. Booth³, M. Huang⁴, and E. J. Fielding²

¹Joint Institute for Regional Earth System Science and Engineering, University of California, Los Angeles, CA, USA

²Jet Propulsion Laboratory, California Institute of Technology, Pasadena, CA, USA

³Department of Geology, Portland State University, Portland, OR, USA

⁴Department of Geology, University of Maryland, College Park, MD, USA

Corresponding author: Alexander L. Handwerger (alexander.handwerger@jpl.nasa.gov)

Key Points:

- Landslide thickness can vary by tens of meters within a single landslide
- Slow-moving landslides have geometric scaling relations that span scaling values that are typical for soil and bedrock landslides
- The largest landslide complexes get larger by increasing area rather than increasing depth
- Landslide strength is scale-dependent, such that large landslides tend to be weaker than small landslides

Abstract

The hazardous impact and erosive potential of slow-moving landslides depends on landslide properties including velocity, surface and subsurface geometry, and frequency of occurrence. However, constraints on subsurface geometry are lacking because these types of landslides rarely fully evacuate material to create measurable hillslope scars. Here we use pixel offset tracking with data from the NASA/JPL Uninhabited Aerial Vehicle Synthetic Aperture Radar (UAVSAR) to measure the three-dimensional surface deformation of 134 slow-moving landslides in the northern California Coast Ranges. We apply volume conservation to infer the actively deforming thickness, volume, geometric scaling, stress, and friction angle of each landslide. These landslides move at average rates between $\sim 0.1\text{--}3$ m/yr and have areas of $\sim 7.8 \times 10^3\text{--}2.63 \times 10^6$ m², inferred mean thicknesses of $\sim 0.4\text{--}22$ m, and volumes of $\sim 7.08 \times 10^3\text{--}9.75 \times 10^6$ m³. The best-fit volume-area geometric scaling exponent is $\gamma = 1.2\text{--}1.5$, indicating that these landslides fall between typical soil and bedrock landslide scaling. A rollover in the scaling relationship suggests that the largest landslide complexes in our dataset become large primarily by increasing in area rather than thickness. In addition, the slow-moving landslides display scale-dependent frictional strength, such that large landslides tend to be weaker than small landslides. This decrease in frictional strength with landslide size is likely because larger landslides are composed of higher proportions of weak material. Our work shows how state-of-the-art remote sensing techniques can be used to better understand landslide processes and quantify their contribution to landscape evolution and hazards to human safety.

1 Introduction

Landslides are a major natural hazard and are often the dominant process that erodes mountainous landscapes (Korup et al., 2007; Larsen et al., 2010; Mackey & Roering, 2011;

Simoni et al., 2013). Both their hazardous impact and erosive potential depend on landslide properties including the velocity, surface and subsurface geometry, and frequency of occurrence. Measuring these landslide properties is challenging because landslides exhibit a wide range of velocities (mm/yr to m/s), spatial areas (10^0 - 10^8 m²), and volumes (10^{-1} - 10^{10} m³), and can occur in large numbers (hundreds to tens of thousands) over broad spatiotemporal scales (Cruden & Varnes, 1996; Hungr et al., 2014; Lacroix, Handwerger, et al., 2020; Larsen et al., 2010). Importantly, the landslide failure style also impacts our ability to measure landslide properties, such as thickness and volume, which can strongly influence runout and erosion rate (e.g., Korup et al., 2007; Larsen et al., 2010; Legros, 2002). Some landslides create clear and identifiable scars and deposits by evacuating material from the hillslope, making it possible to directly measure landslide properties from field data, digital elevation models (DEMs), and remote sensing observations (Bessette-Kirton et al., 2018; Warrick et al., 2019; Wartman et al., 2016). However, for landslides that move slowly for years or centuries (Mackey et al., 2009; Rutter & Green, 2011), i.e., slow-moving landslides (Lacroix et al., 2020), and do not create hillslope scars, it is difficult to constrain their thickness and volume because data are usually limited to isolated point measurements from boreholes (Schulz et al., 2018; Simoni et al., 2013; Travelletti & Malet, 2012), which do not capture the spatial variability exhibited by these landslides. It is therefore advantageous to develop and apply tools and methods that can be used to construct large inventories of slow-moving landslides and quantify their surface and subsurface properties.

Modern remote sensing tools, such as synthetic aperture radar (SAR), optical imagery, and lidar, provide high-resolution measurements of topography and ground surface deformation that can be used to identify and monitor landslides with millimeter- to centimeter-scale accuracy at spatial resolutions of 5 to 100 meters. Recent work using pixel offset tracking and SAR

interferometry with these data has quantified the two-dimensional (2D) and three-dimensional (3D) surface deformation of slow-moving landslides (Aryal et al., 2015; Booth et al., 2020; Hu et al., 2020; Lacroix et al., 2020; Stumpf et al., 2017; Travelletti et al., 2014). These studies, along with numerous ground-based investigations (e.g., Iverson & Major, 1987; Malet et al., 2002; Schulz et al., 2017), have shown that slow-moving landslides exhibit non-uniform spatial and temporal kinematic patterns. In addition, these high-resolution 3D surface deformation measurements can be used to infer the thickness and subsurface geometry of the actively moving part of the landslide (Aryal et al., 2015; Booth et al., 2020; Booth, Lamb, et al., 2013; Delbridge et al., 2016; Hu et al., 2020). These studies suggested that active landslide thickness can vary by tens of meters within a single landslide, and the slip surfaces have an irregular and bumpy morphology that differs considerably from commonly assumed, idealized geometric forms, such as semicircles, ellipsoids, and log spirals (e.g., Michel et al., 2020 and references therein). These large changes in thickness within a single landslide mass have important implications for estimating volume and sediment flux, designing field instrumentation and landslide mitigation strategies, and determining the stresses that control landslide kinematics. Although techniques that invert surface observations for subsurface characteristics are becoming more common, most studies have focused on individual landslides occurring under different and site-specific environmental conditions, making it difficult to identify more generic geometric scaling relations for slow-moving landslides.

In this study, we use data from the NASA/JPL Uninhabited Aerial Vehicle Synthetic Aperture Radar (UAVSAR) to construct an inventory of 134 active slow-moving landslides in a ~4700 km² area of the northern California Coast Ranges between 2016 and 2019 (Figure 1). These landslides occur in the Eel River catchment, a region well known for its slow-moving

landslides, and are driven by high seasonal rainfall (Bennett, Roering, et al., 2016; Booth, Roering, et al., 2013; Handwerger et al., 2013, 2015; Handwerger, Fielding, et al., 2019; Handwerger, Huang, et al., 2019; Kelsey, 1978; Mackey et al., 2009; Mackey & Roering, 2011; Roering et al., 2009, 2015; Schulz et al., 2018). The landslides are underlain by the Central Belt Franciscan mélangé, a mechanically weak and pervasively sheared bedrock with an argillaceous matrix that surrounds blocks of stronger rock types, including sandstone, chert and greenstone (Jayko et al., 1989; Jennings et al., 1977; McLaughlin et al., 1982, 2000). We measure the 3D surface deformation and geometry of each landslide, and use these data in a volume conservation framework to invert for their active thickness, volume, stress, and strength. We derive new geometric scaling relations for slow-moving landslides and make comparisons with a worldwide inventory of soil and bedrock landslides. Our work is the first to perform thickness inversions for a large inventory of landslides, and this approach could be applied to other groups of slow-moving landslides around the world. Our work also shows how state-of-the-art remote sensing techniques can be used to better understand landslide processes and quantify their contribution to landscape evolution.

2 Materials and Methods

2.1 UAVSAR Data and Processing

We use SAR data acquired by the NASA/JPL UAVSAR airborne system for our landslide investigation. UAVSAR has a left-looking radar attached to a NASA Gulfstream III airplane that operates with a L-band wavelength (~ 23.8 cm) and a swath width of ~ 20 km. UAVSAR data have a pixel spacing of 1.67 m in the range direction (measured along the line-of-sight, LOS) and 0.6 m in the azimuth direction (measured along the UAVSAR flight direction). We designed the UAVSAR data collection for the northern California Coast Ranges site

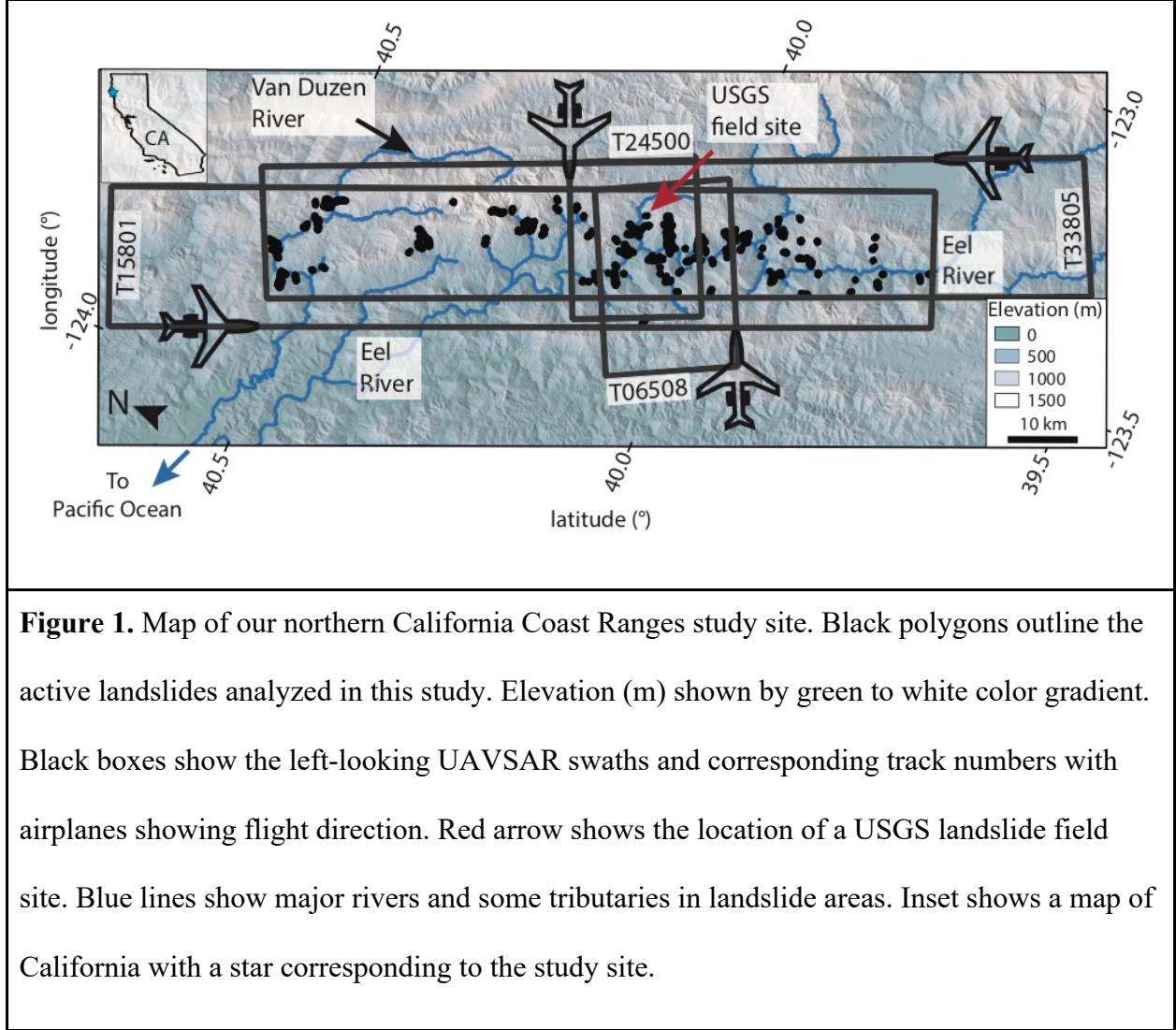
specifically to monitor a large quantity of slow-moving landslides that were initially identified by several previous studies (Bennett, Miller, et al., 2016; Handwerger et al., 2015; Kelsey, 1978; Mackey & Roering, 2011; Roering et al., 2009). Some of these UAVSAR data were used in a recent study by Handwerger, Fielding, et al. (2019) to analyze changes in landslide activity due to extreme rainfall. We collected data on 4 partially overlapping flight paths to increase data redundancy and to provide between 4 and 8 independent deformation measurements (Figure 1). There were 12 data acquisitions at our field site between April 2016 and May 2019. The time between data acquisitions ranges between 47 and 237 days, with a mean of 104 days (Table S1). UAVSAR Single-Look Complex (SLC) data are freely available at <https://uavsar.jpl.nasa.gov/>.

We perform pixel offset tracking on the coregistered UAVSAR stack SLC data using the *Ampcor* application, which is part of the JPL InSAR Scientific Computing Environment (ISCE) version 2 software package (Rosen et al., 2012). Pixel offset tracking (sometimes referred to as subpixel correlation) uses cross-correlation between SAR amplitude images to quantify image offsets (i.e., displacement) due to ground surface motion in two dimensions; 1) the range or look direction, and 2) the azimuth or along-track direction (Fialko et al., 2001; Fielding et al., 2020; Pathier et al., 2006). We use the terms range/look direction and azimuth/along-track direction, interchangeably. Pixel tracking has a precision up to $\sim 1/10$ of the pixel size, which corresponds to ~ 6 cm in the along-track direction and ~ 17 cm in the range direction for a pair of UAVSAR images. Although this technique is less precise than conventional InSAR, it does not involve phase unwrapping and thus is better suited for measuring the decimeter- to meter-scale displacements commonly displayed by many slow-moving landslides (Lacroix, Handwerger, et al., 2020). To account for the differences in the range and along-track pixel size, we use a cross-correlation window length of 128 pixels with a skip size of 32 pixels (distance between matching

window calculations) in the along-track direction and a cross-correlation window width of 64 pixels with a skip size of 16 pixels in the range direction, resulting in a window size of 77 m by 107 m. We geocode the pixel offset measurements to a 0.4 arcsecond (~ 12 m) pixel using the TanDEM-X DEM provided by the German Aerospace Center (DLR). We process all possible combinations of pixel offset tracking pairs, which results in 66 pixel offset tracking maps on each track (264 in total) with single pair time spans ranging from 47 to 1148 days (Table S1). We excluded 35 poor-quality pixel offset tracking maps from our analysis. We found these poor-quality data tend to result from long duration pairs that exceed ~ 2 years, which are subject to numerous changes in the ground surface (e.g., vegetation changes, anthropogenic changes) that can deteriorate the cross-correlation result. We convert all of the offset maps to mean velocities and then take the temporal average of the 31 remaining pixel offset velocity maps to make a mean velocity map for our thickness inversions.

2.2 Three-dimensional Ground Surface Deformation

To solve for 3D deformation from SAR requires at least three independent measurements of surface deformation. Each UAVSAR flight path provides two independent measurements of surface motion from pixel offset tracking (i.e., along-track and range). Therefore, using pixel offset tracking velocity maps, data from at least two flights is required for 3D inversions. Because UAVSAR acquires data on 4 different flight paths in our field area (Figure 1), we have a maximum of 8 deformation measurements in the central region of our field area where all 4 flight paths overlap and a maximum of 4 deformation measurements in the northern and southern extents where only 2 flight paths overlap. Thus, we are always able to achieve an overdetermined 3D inversion.



Each deformation measurement from pixel tracking is composed of the true displacement vector projected onto the along-track or LOS direction of the UAVSAR. We use a least-squares inversion to isolate the east, north, and vertical components of deformation defined in the form

$$\mathbf{d} = \mathbf{Gm},$$

$$\begin{bmatrix} v_{rng1} \\ v_{azi1} \\ \vdots \\ v_{rng,M} \\ v_{azi,M} \end{bmatrix} = \begin{bmatrix} \cos \xi_1 \sin \theta_1 & \sin \xi_1 \sin \theta_1 & -\cos \theta_1 \\ \cos \xi_1 & \sin \xi_1 & 0 \\ \vdots & \vdots & \vdots \\ \cos \xi_M \sin \theta_M & \sin \xi_M \sin \theta_M & -\cos \theta_M \\ \cos \xi_M & \sin \xi_M & 0 \end{bmatrix} \begin{bmatrix} v_{ew} \\ v_{ns} \\ v_{ud} \end{bmatrix}, \quad (1)$$

where $v_{rng,M}$ is the range (or look direction) velocity, $v_{azi,M}$ is the azimuth (or along-track direction) velocity, M is the flight path number (minimum of 2 needed for pixel offset tracking), ξ is the UAVSAR heading direction (i.e., along track direction) with counterclockwise as positive, θ is the UAVSAR look angle, and v_{ew} , v_{ns} , v_{ud} are the east-west, north-south, and vertical components of velocity, respectively.

The overdetermination of the 3D inversion allows us to reduce inversion error and constrain the uncertainty from the inversion. To constrain the uncertainty, we repeat the 3D inversion multiple times using different combinations of v_{rng} and v_{azi} . For instance, for landslides with 8 deformation measurements (i.e., 4 range and 4 azimuth measurements), we perform the 3D inversion 198 times using between 3 and 8 deformation measurements. We then take the mean and standard deviation of all of the inversions and use these values as the 3D velocities and inversion uncertainty, respectively. We further constrain the uncertainty in our velocity measurements by examining the apparent deformation rate of stable hillslopes. To reduce noise and error (i.e., unrealistically large displacements), we apply velocity thresholds and mask out pixels with apparent velocities > 50 m/yr, which is much faster than the typical velocity range displayed by the northern California Coast Ranges landslides (Bennett, Roering, et al., 2016; Handwerger, Fielding, et al., 2019; Roering et al., 2015). We also mask out pixels that have mean velocities less than their inversion uncertainty and use nearest neighbor interpolation to fill in these masked pixels.

2.3 Landslide Thickness Inversion

We use 3D surface velocity measurements from pixel offset tracking to infer the thickness, volume, and shear zone geometry of the active parts of each landslide using a

conservation of volume approach. We apply the method originally described by Booth, Lamb, et al., (2013) and more recently by Booth et al. (2020), which assumes that during our ~3 year study period, the measured surface velocity is representative of the depth-averaged velocity, the sliding surface does not change in time, there is minimal direct erosion or deposition of the landslide surface, and the landslide material density is uniform and constant. While landslides may violate these assumptions in general, they are reasonable for our study area for the following reasons: (1) at the Two Towers landslide, a U.S. Geological Survey (USGS) instrumented landslide in our study site (Schulz et al., 2018), the measured surface velocity was approximately equal to the depth-averaged velocity, and a narrow shear zone was identified (Figure S1); (2) the landslides were continuously active over the time periods that 3D displacements were measured, suggesting movement on the same slip surface; (3) minor amounts of direct surface erosion or deposition were likely confined to gully systems on the landslides' surfaces, which occupy a small percentage of the landslides' surface area (~1%) and therefore have a minimal effect on the inversion; and (4) dilation/compaction or shrinking/swelling that would cause changes in density is likely on the order of centimeters or less, which is small compared to surface velocity gradients (Booth et al., 2020; Delbridge et al., 2016; Iverson, 2005; Schulz et al., 2018), thus having limited influence of the measured 3D surface velocity. Therefore, for a landslide of constant density with no erosion or deposition, conservation of volume implies that

$$v_{ud} = \nabla \cdot (\bar{u}h) + u_{surf} \cdot \nabla z_{surf}, \quad (2)$$

where v_{ud} is the vertical component of the 3D landslide surface displacement vector, h is the active landslide thickness, u_{surf} is the vector of horizontal components of landslide surface velocity, \bar{u} is the depth-averaged vector of horizontal components of landslide velocity, and z_{surf} is the surface elevation measured from the 12 m TanDEM-X DEM. The first term on the right-

hand side of equation 2 is the contribution of flux divergence to the vertical component of the surface velocity, and the second term is the contribution due to advection of the sloped land surface. Because UAVSAR measures the velocity of the ground surface, u_{surf} , we assume that $\bar{u} = fu_{surf}$, where f is a constant that characterizes the thickness of the shear zone at the base of the landslide relative to the total landslide thickness. We constrain f using borehole inclinometer data from two boreholes at the USGS field station on the Two Towers landslide (supporting information and Figure S1). Unfortunately, the Two Towers landslide is not detectable with pixel tracking from UAVSAR data because the landslide is small (250 m long and 40 m wide) and moving too slowly (maximum speed ~ 6 cm/yr). Therefore, we are not able to directly test our thickness inversion method on the Two Towers landslide. Nonetheless, the ground-based data provide key information to constrain the value of f . Nearly all shear strain occurs in a zone that is < 0.3 m in one borehole and < 0.6 m in the second borehole. Using these data, we find that $f \sim 0.96$, which indicates that the landslide moves along a narrow shear zone with the material above translating essentially as a rigid block. For simplicity, we assume that $f = 1$ and that the landslides move as a rigid block. Other studies in California have also found that landslides move as a rigid plug above a narrow shear zone such that $f \sim 1$ is a reasonable approximation (Keefer & Johnson, 1983). Although f generically represents the ratio of depth-averaged to surface velocity, it can be related to specific rheologies if desired (Booth, Lamb, et al., 2013; Delbridge et al., 2016).

Incorporating f into equation 2 gives

$$v_{ud} = \nabla \cdot (fu_{surf}h) + u_{surf} \cdot \nabla z_{surf}, \quad (3)$$

which is a statement of conservation of volume in a Lagrangian reference frame (Booth et al., 2020; Delbridge et al., 2016). We discretize equation 3 using centered finite differences,

rearrange it as a system of linear equations, and then solve for thickness by minimizing, subject to non-negative constraints,

$$|Xh - b|^2 + \alpha^2 |\nabla^2 h|^2, \quad (4)$$

where X is a diagonally dominant matrix that contains the depth-averaged horizontal velocity data, b is a vector containing the topographic gradient and surface horizontal and vertical velocity data, and α is a damping parameter to regularize the ill-posed inverse problem. We explore a wide range of α from 10^{-3} to 10^1 and determine the best level of regularization using the Generalized Cross-Validation method (supporting information and Figure S2). We resample our ~ 12 m pixel spacing grid to square 10 m x 10 m pixel and perform the thickness inversion in the MATLAB software package using the CVX program, a package for specifying and solving convex programs (Grant & Boyd, 2014). The inferred thickness values represent the best solution that does not violate conservation of volume and assumes that the surface velocity is equal to the depth-averaged velocity. It is important to further emphasize that the thickness inversions are only relevant to the active parts of landslides such that there needs to be detectable surface deformation to invert for the landslide thickness. Specifically, the values of b (equation 4) need to differ from background values on known stable ground to infer non-zero thicknesses. Landslides or areas and kinematic zones within landslides that are not moving are therefore considered to have zero depth. Thickness in this study therefore means the “active thickness” during our study period.

Since both the matrix X and the vector b contain data with uncertainties, and the damping parameter necessarily introduces bias, estimating total uncertainty of the resulting thickness model is not straightforward. However, we make a minimum estimate following standard techniques from inverse theory, which reflects uncertainty in b only (supporting information).

Thickness uncertainty estimated in this way increases with landslide size and ranges from approximately ± 1.5 m to ± 4 m from the smallest to largest landslides in the study area (Figure S3).

2.4 Landslide Inventory and Geometric Scaling

To select landslides for 3D surface velocity and thickness inversions, we assemble a new inventory of active landslides in our ~ 4700 km² study area in the northern California Coast Ranges that includes only those landslides that show a significant deformation signal using the pixel offset tracking method. This limits our analysis to the faster-moving landslides that exhibit rates of decimeters to meters per year. Our landslide inventory was guided by a number of pre-existing landslide inventories for the northern California Coast Ranges (Bennett, Miller, et al., 2016; Handwerger, Fielding, et al., 2019; Kelsey, 1978; Mackey & Roering, 2011). We map the landslide boundaries in QGIS using the 3D velocity maps, hillshade maps constructed from 1 m pixel spacing lidar provided by OpenTopography (Roering, 2012), the ~ 12 m pixel spacing TanDEM-X DEM, and Google Earth imagery. Because slow-moving landslides display non-uniform spatial kinematic zones and complex kinematic histories (e.g., Nereson & Finnegan, 2018; Schulz et al., 2017; Stumpf et al., 2017), there are often differences between the landslide boundaries mapped with kinematic data and those mapped based on geomorphic interpretation of hillshades or aerial photos. These differences in mapping are especially important for our thickness inversions because including the parts of landslides that are not currently moving can cause the thickness inversion to produce unreliable results. Therefore, we use the temporally averaged landslide velocity and only map areas of each landslide that are moving during our study period. We then use QGIS to quantify the spatial metrics of each landslide, including

length, average width (defined as area divided by length), area, and slope angle. We also report the mean, median, 75th percentile, and maximum horizontal velocity, 3D velocity magnitude, and 3D inversion velocity errors for each landslide.

We then derive empirical geometric scaling relations for landslide thickness (h) and volume V from the measured landslide area A . Geometric scaling relations are commonly used to quantify erosion rates of large inventories of landslides and are important for understanding landslide mechanics (e.g., Guzzetti et al., 2009; Larsen et al., 2010; Milledge et al., 2014). These relations take the form of a power function where

$$V = c_V A^\gamma \text{ and } h = c_h A^\zeta, \quad (5a \text{ and } 5b)$$

where γ and ζ are scaling exponents and c_V and c_h are the intercepts. We constrain the coefficients of these power functions by log-transforming our data and finding the best-fit parameters with 95% confidence intervals using a linear least square inversion in MATLAB. Larsen et al. (2010) showed that these scaling relations hold over 9 orders of magnitude in area and 12 orders of magnitude in volume.

2.5 Stress and Frictional Strength

We estimate the basal shear stress and frictional strength of each landslide by following the 3D Simplified Janbu method (Bunn et al., 2020; Hungr, 1987; Hungr et al., 1989; Leshchinsky, 2019). This method assumes that the vertical intercolumn shear forces are negligible. Each landslide is discretized into 3D columns with a surface area S_{basal} and total weight W . The basal surface area is defined by

$$S_{basal} = \Delta x \Delta y \frac{1 - \sin^2 \beta_y \sin^2 \beta_x}{\cos \beta_y \cos \beta_x}, \quad (6)$$

where Δx and Δy are the grid spacing in the x and y direction, respectively, β_x is the local dip angle perpendicular to the direction of motion and β_y is the local dip in the direction of motion.

The normal force N at the base of each column is defined by

$$N = \frac{W + pS_{basal} \tan \phi \sin \beta_x / F}{\cos \Delta_z \left(1 + \frac{\sin \beta_x \tan \phi}{F \cos \Delta_z} \right)}, \quad (7)$$

where p is the mean pore pressure acting at the base of each column, ϕ is the residual friction angle, F is the factor of safety, and Δ_z is the local dip angle defined in terms of the motion-parallel and motion-perpendicular dips by

$$\cos \Delta_z = \left(\sqrt{\frac{1}{1 + \tan^2 \beta_y + \tan^2 \beta_x}} \right). \quad (8)$$

Finally, F is defined by

$$F = \frac{\sum (N - pS_{basal}) \tan \phi \cos \beta_x}{\sum N \cos \Delta_z \tan \beta_x}, \quad (9)$$

where the summation is over all columns. The numerator is the resisting force, with the term in the parentheses defining the effective normal force, and $\tan \phi$ is the friction coefficient, and the denominator is the shear force. Cohesion is assumed to be negligible since our landslides are already moving. We set $F = 1$ (i.e. balanced forces at failure) and solve for shear stress, effective normal stress (shear force and normal force divided by S_{basal} , respectively), and friction angle under both dry and fully saturated (hydrostatic conditions) end members to produce a minimum and maximum estimate. Table S2 shows the dry and wet landslide density values used for our calculations.

3 Results

3.1 Landslide Inventory and 3D Velocity

We identified 134 active landslides in our northern California Coast Ranges field site (Figure 1), 19 of which were unmapped by previous studies (Bennett, Miller, et al., 2016; Handwerger, Fielding, et al., 2019; Mackey & Roering, 2011). These landslides have average widths from 66 to 556 m, lengths from 68 to 4727 m, areas from 7.8×10^3 to 2.63×10^6 m², and mean slope angles from 10 to 29 degrees (Table S3). Each landslide exhibited a non-uniform spatial velocity pattern (see examples in Figure 2). The spatial kinematic patterns remain fixed during our study period and are similar to those mapped in previous studies (see Bennett, Roering, et al., 2016; Handwerger, Fielding, et al., 2019; Mackey & Roering, 2011). The characteristic 3D velocity magnitude, defined as the 75th percentile value for each landslide, was calculated as $v_{3D} = (v_{ns}^2 + v_{ew}^2 + v_{ud}^2)^{1/2}$. The 3D velocity magnitude ranged from 0.162 to 2.92 m/yr. The average 3D velocity for the entire inventory was 0.622 ± 0.513 m/yr (± 1 standard deviation). The landslide motion was always primarily in the downslope direction (see example in Figures 2e and 2f), but at different locations we do measure areas of both uplift and subsidence within a single landslide (see example in Figure 2d). We note that local surface uplift occurs when the vertical component of the velocity vector dips less steeply than the topographic surface at a given point. As a result, the vertical velocity is often still negative even in areas where the topographic surface is locally being uplifted, and only when the vertical motion is upwards relative to horizontal do we observe positive vertical velocities. The 3D velocity uncertainty from the 3D inversion (equation 1) averaged over all 134 landslides was 0.1432 m/yr. We report the uncertainty for each individual landslide in Table S3. The 3D velocity magnitude uncertainty from examining the apparent velocity of stable hillslopes was ≤ 0.1 m/yr.

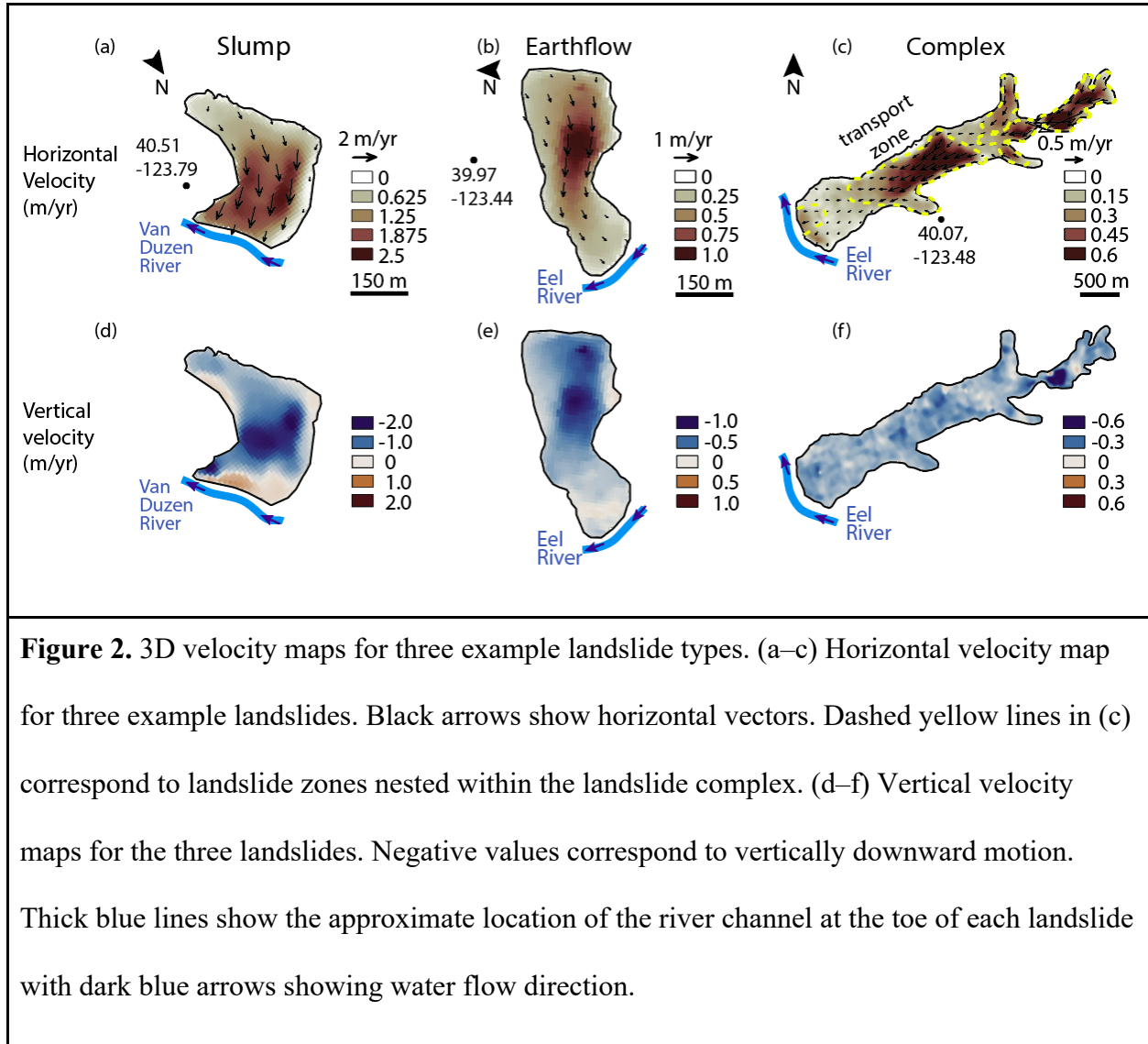
We classified the slow-moving landslides into three subgroups based on their geometry and kinematic patterns (Table S3). Figure 2 shows three example landslides which we define as

slumps, earthflows, and landslide complexes. The landslide complex shown in Figure 2 is the largest landslide in our dataset and is also known as the Boulder Creek landslide in several other studies (Bennett, Miller, et al., 2016; Bennett, Roering, et al., 2016; Handwerger et al., 2015, 2015; Handwerger, Fielding, et al., 2019; Handwerger, Huang, et al., 2019; Mackey & Roering, 2011; Roering et al., 2009). We defined slumps as landslides with lower length/width aspect ratios (median = 1.57 ± 1.00 , ± 1 standard deviation), a strong signal of positive vertical velocity components in the toe and negative vertical velocity components in the source area, and one primary kinematic zone (Figure 2a). We defined earthflows landslides as those with medium aspect ratios (median = 3.56 ± 1.88 , ± 1 standard deviation), one primary kinematic zone, and small magnitude, but mostly negative, vertical velocity components (Figure 2b). And we defined landslide complexes as those with higher aspect ratios (median = 5.13 ± 2.34 , ± 1 standard deviation), that are composed of multiple kinematic zones or even multiple landslides that coalesce into a single landslide mass (Figure 2c). Landslide complexes are relatively common in areas with slow-moving landslides (Cerovski-Darriau & Roering, 2016; Keefer & Johnson, 1983; Simoni et al., 2013). The mean 3D velocity magnitude was 0.585, 0.606, and 0.670 m/yr for slumps, earthflows, and landslide complexes, respectively.

3.2 Thickness, Volume, and Geometric Scaling Relations

The non-uniform kinematic patterns exhibited by these landslides are also reflected in their inferred subsurface geometry (Figure 3). We find that the thickness of each landslide varies spatially and can range from zero to tens of meters within the landslide boundaries. The slip surfaces are generally concave-up, but are rough and irregular in places, especially for landslide complexes. The mean active thickness of the individual landslides ranged from 0.4 to 22.4 m,

and the maximum active thickness ranged from 2.25 to 89.6 m. The mean, median, min, max, and standard deviation active thickness for each landslide are reported in Table S3. Below we describe our findings for the three example types of landslides shown in Figure 2. We note again that these landslides represent their subgroups to first order.



The example slump has one primary deep zone with a mean active thickness of ~ 7.5 m, maximum active thickness of ~ 31 , and standard deviation of ~ 8.7 m. The slip surface has a

concave-up profile. The slope of the slip surface deviates from the ground surface and is steeper near the headscarp and gentler near the toe. Some areas within the head of the landslide are inferred to have no active depth because the values of b (equation 4) are similar in magnitude to background values there. The example earthflow generally has a concave-up slip surface with some irregular bumps. The slip surface more closely mimics the ground surface in the main transport zone and the landslide has a mean active thickness of ~ 15 m, max active thickness of ~ 62 m, and a standard deviation of ~ 17 m. Lastly, the example landslide complex (i.e., Boulder Creek landslide complex) has several different active zones, each with an alternating concave-up and convex-up slip surface profile. The landslide slip surface is rough and irregular over the length of the entire landslide, but each deep zone corresponds to the different kinematic units that comprise the landslide complex (Figure 2c). The landslide complex has a mean active thickness of ~ 3.7 m, maximum active thickness of ~ 29 m, and a standard deviation of ~ 4.6 m. These thicknesses are small compared to the two previous examples because many areas of the landslide do not have a resolvable active thickness.

The Boulder Creek landslide complex shows that landslide complexes may correspond to smaller, faster, and possibly shallower features that are superimposed on a larger, slower, and possibly deeper-seated failure. If multiple failure planes are indeed present at these landslides, that would violate the assumption of a constant f throughout the landslide and cause unreliable thickness estimates. To minimize those potential effects, we performed separate thickness inversions for any isolated, faster-moving areas of a larger landslide complex, as well as for the entire landslide complex as a whole. If results were substantially different, we adopted the more reliable results for the smaller isolated landslides.

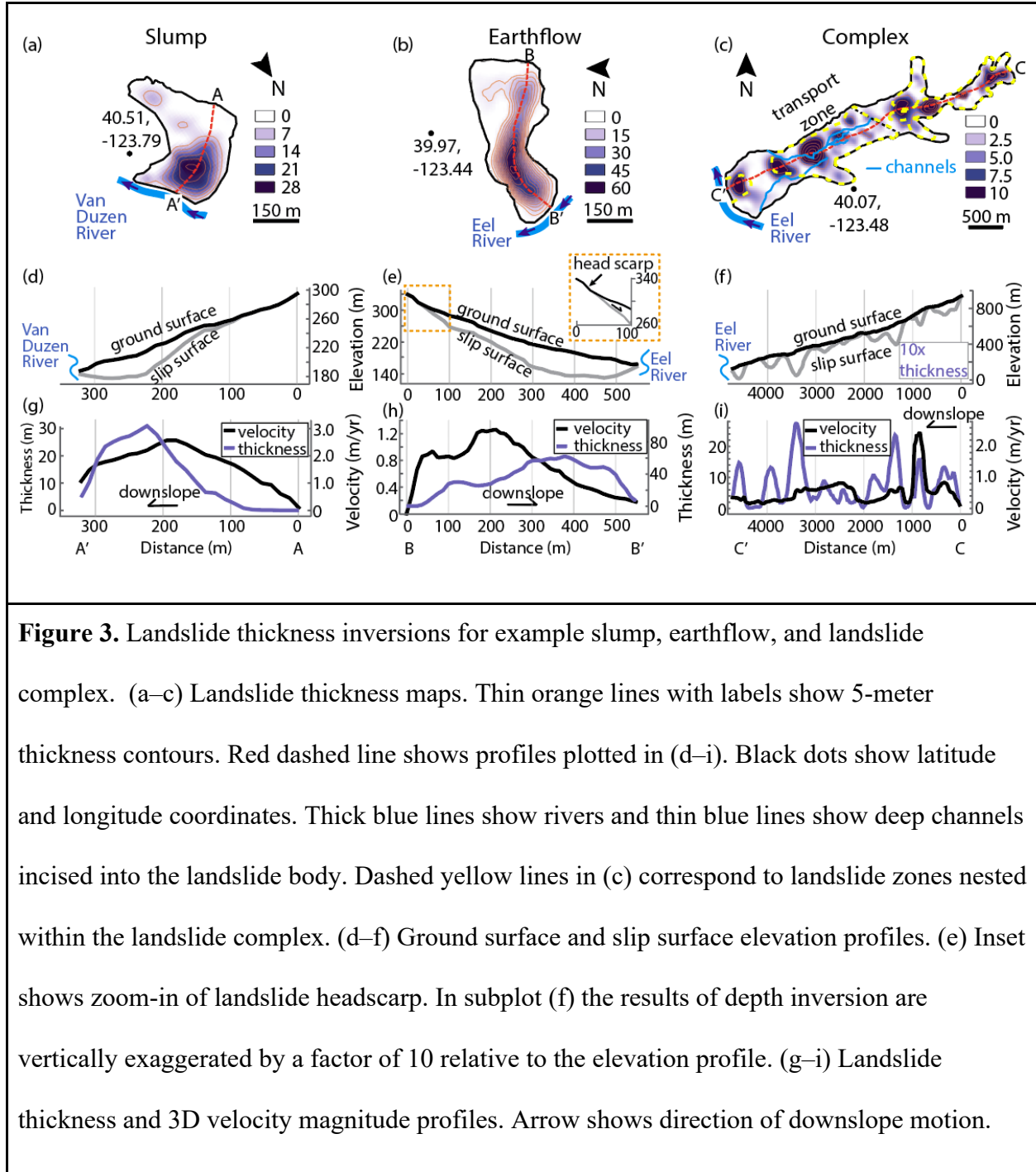


Figure 3. Landslide thickness inversions for example slump, earthflow, and landslide complex. (a–c) Landslide thickness maps. Thin orange lines with labels show 5-meter thickness contours. Red dashed line shows profiles plotted in (d–i). Black dots show latitude and longitude coordinates. Thick blue lines show rivers and thin blue lines show deep channels incised into the landslide body. Dashed yellow lines in (c) correspond to landslide zones nested within the landslide complex. (d–f) Ground surface and slip surface elevation profiles. (e) Inset shows zoom-in of landslide headscarp. In subplot (f) the results of depth inversion are vertically exaggerated by a factor of 10 relative to the elevation profile. (g–i) Landslide thickness and 3D velocity magnitude profiles. Arrow shows direction of downslope motion.

405

406 For example, we delineated the Boulder Creek landslide complex into 5 smaller landslides based

407 on the separate kinematic zones and performed thickness inversions for each kinematic zone or

408 sub-landslides (Figure 2c; Figure S4). We found that the spatial patterns of inferred thickness for

the sub-landslides were similar to the full landslide with some differences in the thickness magnitude suggesting it is appropriate to consider Boulder Creek either one large landslide complex that has 5 active areas or 5 smaller landslides.

We also found that the inferred active thickness for the Boulder Creek landslide was particularly irregular. While we expect areas that are not currently active to thin and even have zero thickness in places, the active transport zone on Boulder Creek also contains thin and thick patches (Figures 3c, 3f, and 3i). One explanation for this variability in the transport zone is that there is a large channel network incised into the Boulder Creek landslide (Figure 3). In some places the channel reaches depths of 15-20 meters (Figure S5). Since the thickness is measured as the vertical distance from the ground surface to the inferred basal sliding surface, the predicted thickness is expected to be low in places surrounding the channel if the channel depth is similar to the landslide thickness. Our findings indicate that the channel has incised to depths that approach the predicted sliding surface in several places (Figure S5).

Although we do not have borehole data to confirm our thickness estimates, we used the topography to verify the inferred slip surface elevation in several cases. Figure 3e inset and Figure S6 shows the landslide that has a clear headscarp that can be used to trace the sliding surface underneath the ground surface. The extension of the headscarp slip surface under the landslide provides confirmation that the inversion is approximating the slip surface elevation correctly. Figure S7 shows another slow-moving landslide that has filled into a pre-existing valley. Transects across this landslide show the ground surface of the filled-in valley and that the slip surface has the shape of the pre-existing valley, providing additional confirmation that our inversions are approximating the slip surfaces correctly. In addition, we compared our thickness inversions to thickness estimates from lidar. Mackey and Roering (2011) used lidar to measure

the toe height at the channel interface for dozens of landslides in the Eel River catchment. Of those landslides, 10 can be used to make direct comparisons with our dataset. We found overall good agreement between the landslides toe thickness estimated from lidar and from our inversions (Figure 4).

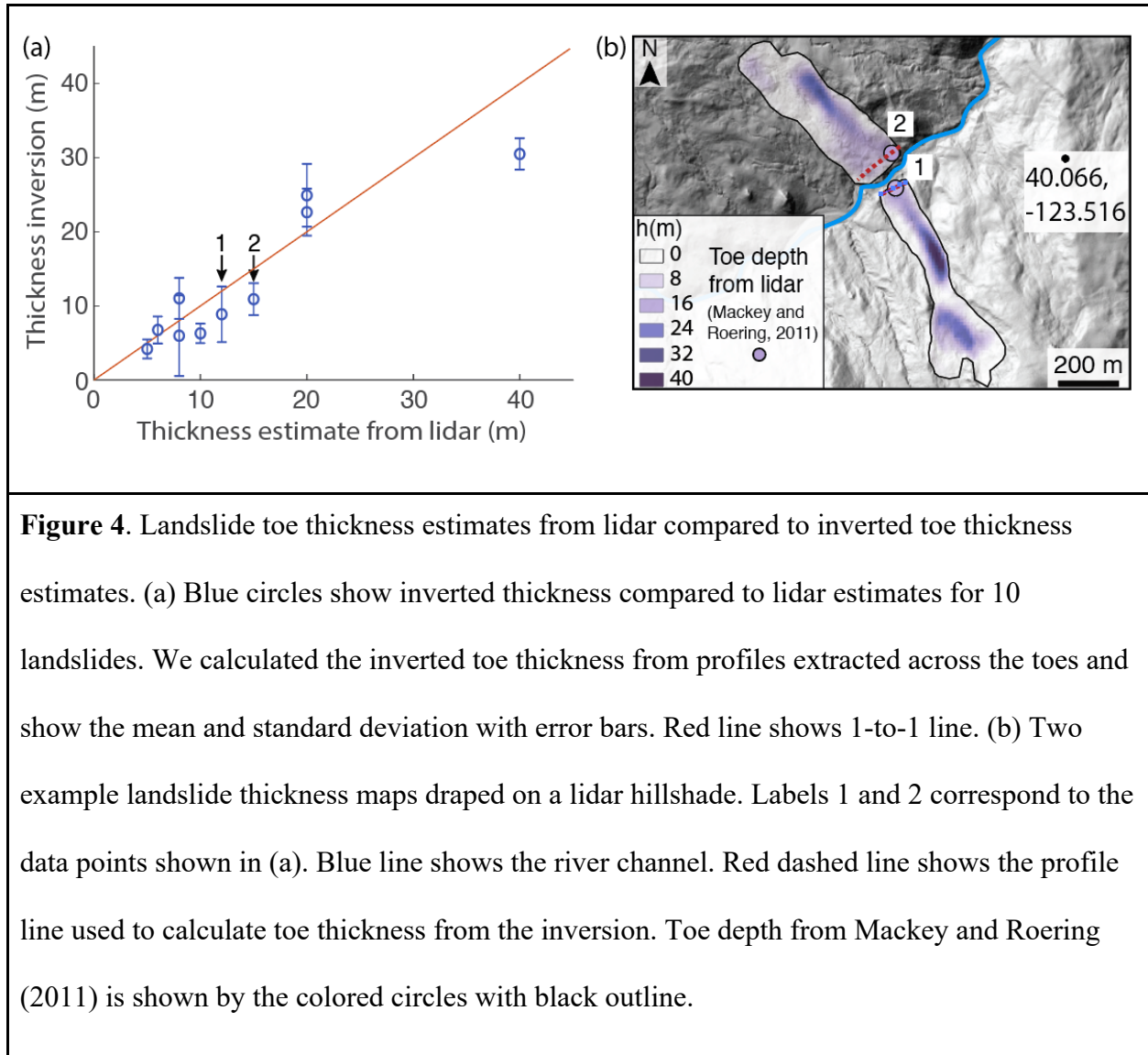


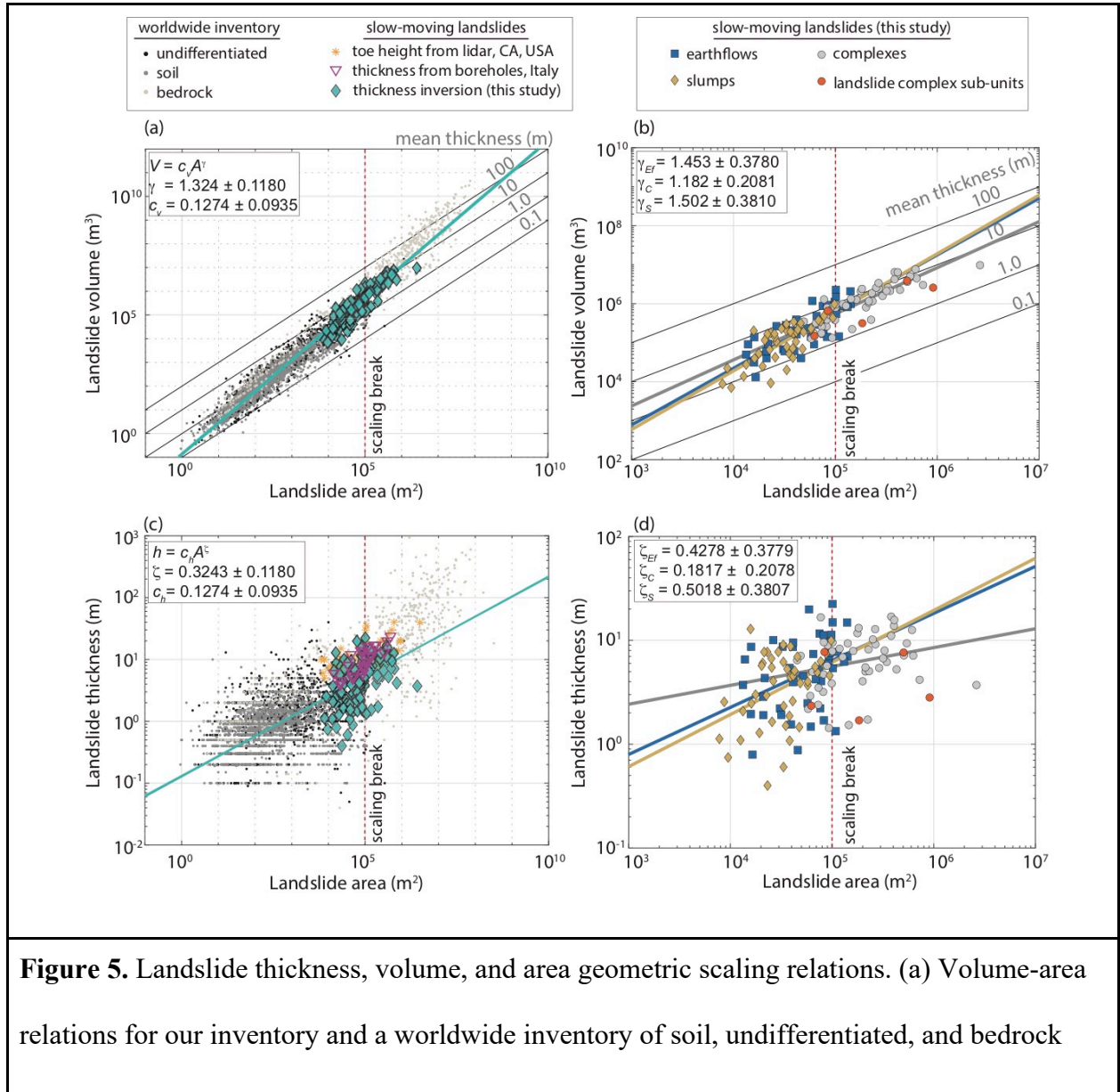
Figure 4. Landslide toe thickness estimates from lidar compared to inverted toe thickness estimates. (a) Blue circles show inverted thickness compared to lidar estimates for 10 landslides. We calculated the inverted toe thickness from profiles extracted across the toes and show the mean and standard deviation with error bars. Red line shows 1-to-1 line. (b) Two example landslide thickness maps draped on a lidar hillshade. Labels 1 and 2 correspond to the data points shown in (a). Blue line shows the river channel. Red dashed line shows the profile line used to calculate toe thickness from the inversion. Toe depth from Mackey and Roering (2011) is shown by the colored circles with black outline.

436

Using our thickness inversions for each landslide, we estimate that the individual landslide volumes range from 7.08×10^3 to $9.75 \times 10^6 \text{ m}^3$. We fit a power function to the volume-area to characterize the geometric scaling relations for these slow-moving landslides. We

also compared our inventory to a worldwide inventory of catastrophic soil, undifferentiated, and
 bedrock landslides compiled by Larsen et al. (2010). We find that the slow-moving landslides in
 the northern California Coast Ranges are larger in both area and volume than most soil
 landslides, but smaller than the largest bedrock landslides around the world. The best fit volume-
 area power function exponent (with 95% confidence) for the inventory was $\gamma = 1.324$ (1.206,
 1.442) and has a R-square of 0.7869 (Figure 5). We report the scaling constants c_V and c_h in
 Table S4. We observed an apparent break in the slope of the volume-area relation for the largest
 landslides in our inventory with area $> 10^5 \text{ m}^2$. To further investigate this break in slope, we also
 fit volume-area scaling as a function of landslide type (slumps, earthflows, or complexes) and
 find that the break in slope is primarily associated with the landslide complexes. By fitting a
 power function to each landslide type, we find slumps $\gamma_S = 1.502$ (1.121, 1.883) with R-square:
 0.6014, earthflows $\gamma_{Ef} = 1.453$ (1.075, 1.831) with R-square: 0.6016, and complexes $\gamma_C = 1.182$
 (0.9739, 1.389) with R-square: 0.7402. Although these parameters are not statistically distinct at
 the 95% confidence level, the fact that γ_S and γ_{Ef} overlap more with each other than with γ_C
 supports the argument that the break in slope is likely related to landslide type. This likely
 change in volume-area scaling for the landslide complexes shows that the largest landslides in
 our inventory get larger in area but do not continue to get significantly deeper with increasing
 area. In addition, we calculated the thickness-area scaling relations using the mean thickness to
 represent each landslide. We compared these scaling relations to point based estimates (lidar)
 and measurements (boreholes) of landslide thickness for slow-moving landslides in the northern
 California Coast Ranges (Mackey and Roering, 2011) and the Reno River catchment, Apennines,
 Italy (Simoni et al., 2013) (Figure 5). The best fit depth-area power function exponent (with 95%
 confidence) for the inventory $\zeta = 0.3243$ (0.2063, 0.4424) with R-square: 0.1828, indicating a

weak increase in depth with area for the inventory as a whole (Figure 5). We also fit depth-area scaling as a function of landslide type and find slumps $\zeta_S = 0.5018$ (0.1211, 0.8826) with R-square: 0.1442, earthflows $\zeta_{Ef} = 0.453$ (0.07515, 0.8308) with R-square: 0.128, and for landslide complexes $\zeta_C = 0.1817$ (-0.02609, 0.3895) with R-square: 0.0631. Therefore, landslide thickness significantly increases with area for slumps and earthflows, but does not significantly vary with area for landslide complexes.



landslides (Larsen et al. 2010). (b) Volume-area relation for slumps, earthflows, and landslide complexes. (a, b) Thin diagonal black lines show volume-area for various constant mean thicknesses. (c) Thickness-area relations for our inventory (mean thickness), the worldwide inventory (Larsen et al., 2010), and slow-moving landslides in the northern California Coast Ranges (Mackey and Roering, 2011) and the Apennine mountains, Italy (Simoni et al., 2013). (d) Landslide thickness-area relations by category. Orange circles in (b, d) correspond to the Boulder Creek landslide complex split into 5 smaller landslides (Figure S4). Red dashed vertical line shows an apparent break in scaling for the largest landslide complexes in our dataset.

469

470 3.3 Stress and Frictional Strength

471 We calculated the basal shear stress and effective normal stress under dry and saturated
 472 conditions end members assuming nil cohesion (Figure 6; Table S2). We find that the average
 473 shear stress for each landslide ranged from ~ 1.30 to ~ 102 kPa assuming dry conditions and ~ 1.51
 474 to ~ 122 kPa assuming fully saturated conditions. The average effective normal stress ranged
 475 from ~ 7.00 to ~ 373 kPa for dry conditions and ~ 3.98 to ~ 215 kPa for saturated conditions. The
 476 inferred landslide friction angle ϕ ranged from ~ 7 to ~ 28 degrees for dry conditions and ~ 13 to
 477 ~ 54 degrees for saturated conditions. We note that the inferred friction angle is larger when
 478 saturated conditions are assumed due to higher pore-water pressure. i.e., the pore-water pressure
 479 reduces the effective normal stress, so the inferred friction coefficient must be higher to maintain
 480 a factor of safety equal to one (equation 9). Our inferred friction angles encompass friction angle
 481 values measured in the laboratory for Franciscan mélange rocks and landslide material (Figure
 482 6). We also analyzed the friction angle as a function of landslide size and found a weak

decreasing power-function relationship with increasing size, defined as $\phi = kL^r$, where k is the intercept, L is the landslide length, and r is the exponent. For dry conditions (with 95% confidence interval), $k = 56.55$ (39.37, 73.73) and $r = -0.2052$ (-0.2562, -0.1541) with R-square: 0.3257. For saturated conditions (with 95% confidence interval), $k = 108.2$ (74.74, 141.8) and $r = -0.2073$ (-0.2593, -0.1553) with R-square: 0.3225. Figure 6c shows the mean friction angle for each landslide as a function of landslide length. The negative trend indicates that the largest landslides are weaker, on average, than smaller landslides.

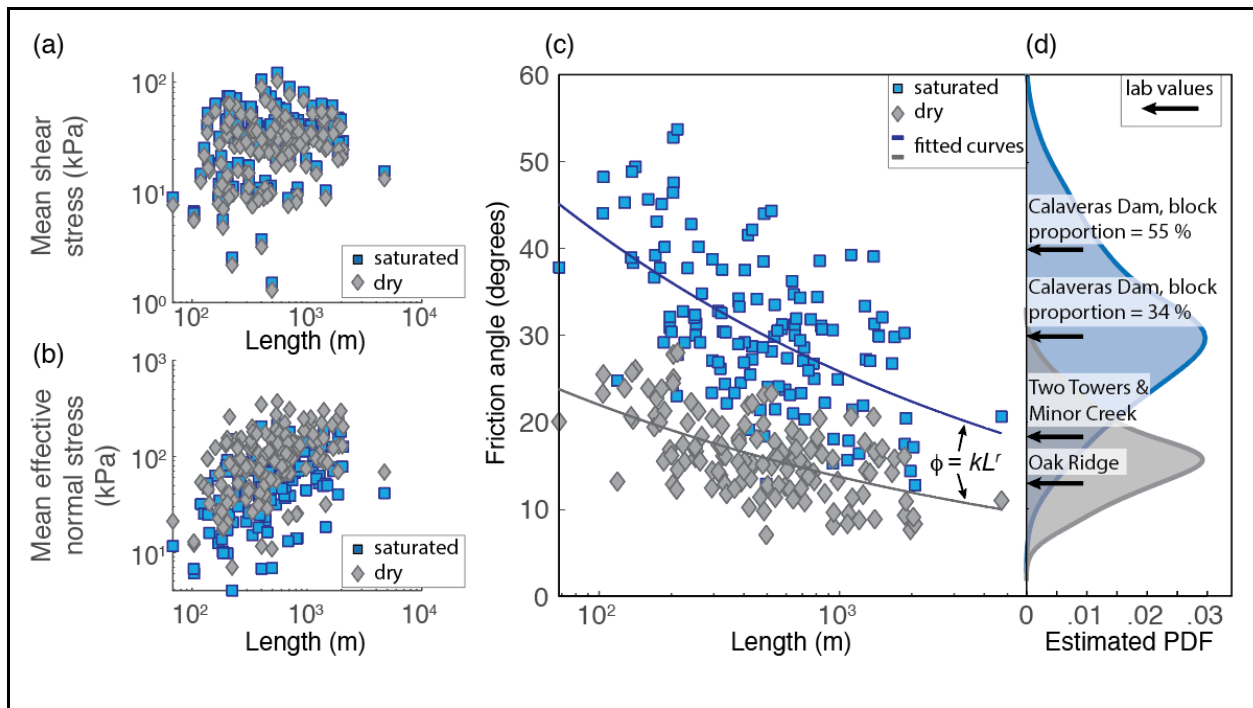


Figure 6. Inferred basal stresses and friction angle for dry and saturated end-members. (a) Mean shear stress, (b) mean effective normal stress, and (c) friction angle for each landslide. Solid line lines in (c) correspond to best-fit power function curves $\phi = kL^r$. For dry conditions, $k = 56.55$ and $r = -0.2052$ and for saturated conditions $k = 108.2$ and $r = -0.2073$. (d) Estimated probability density function for the full inventory. Black arrows show lab-based friction angle values for the Franciscan mélangé hosted Oak Ridge, Two Towers, and Minor Creek

landslides (Iverson & Major, 1987; Nereson & Finnegan, 2018; Schulz et al., 2018) and the Calaveras Dam, which is founded on Franciscan mélange (Roadifer et al., 2009). The Calaveras Dam samples are plotted for two different block-in-matrix proportions.

4 Discussion

4.1 Landslide Kinematics

Our 3D UAVSAR velocity measurements reveal 134 active slow-moving landslides in the northern California Coast Ranges moving at average rates from cm/yr to m/yr between 2016 and 2019. The 3D velocity data confirm that the motion of these landslides is generally in the downslope direction and that there are segments with vertical uplift. Uplift tends to be at the landslide toe due to the concave-up slip surface geometry, and the tendency for longitudinal shortening in the direction of motion to occur at the toe. However, it is possible that a component of uplift of landslide surfaces could also result from dilation or swelling (volumetric expansion), but the magnitude is small, likely on the order of a few centimeters at most (Booth et al., 2020; Delbridge et al., 2016; Iverson, 2005; Schulz et al., 2018). Including volume changes such as this in the thickness inversion would refine its results, but the amount of dilation or compaction occurring throughout an entire landslide and its temporal variation is generally unknown.

Our findings agree with previous work in this region that shows that these landslides exhibit slow, spatially non-uniform downslope motion. Several of the landslides in our study area (e.g., Boulder Creek) have been moving in this manner since at least 1944 (Bennett, Roering, et al., 2016; Mackey & Roering, 2011). Our findings also show that pixel offset tracking with very high resolution UAVSAR data is well-suited for monitoring landslides moving at rates > 10 cm/yr. Some satellites acquire very high resolution SAR with Spotlight

modes, including German TerraSAR-X and Italian COSMO-SkyMed that could provide similar measurements (e.g., Madson et al., 2019). We note that there are likely active landslides moving < 10 cm/yr that were not observed by our pixel offset tracking measurements. In addition, catastrophic landslides are not detected by our pixel offset tracking approach and would require different multi-temporal analysis of SAR, optical, or DEMs to identify these types of landslides (e.g., Intrieri et al., 2017; Jung & Yun, 2020). Our future work will revisit the time-dependent motion and strain rates of each landslide in relation to ongoing climate shifts using InSAR and pixel offset tracking techniques as we continue to collect UAVSAR data over many landslides in California.

4.2 Landslide Geometry

Our study is the first (to our knowledge) to apply landslide thickness inversions to more than a single landslide in a given region. Previous work (Booth et al., 2020; Booth, Lamb, et al., 2013; Delbridge et al., 2016) has used the same approach to analyze individual landslides, but these landslides occur in different regions and environmental conditions. Like these previous studies, however, we found that the active landslide thickness is variable and that the slip surfaces are rough and irregular in places. The non-uniform thickness and velocity of each landslide results in a non-uniform sediment flux, which has implications for understanding sediment motion along hillslopes (Booth et al., 2020; Guerriero et al., 2017). The shape of the slip surface likely also impacts the landslide kinematics and groundwater flow (Coe et al., 2009; Guerriero et al., 2014; Iverson & Major, 1987; Keefer & Johnson, 1983). Slip surfaces that are bumpy and rough may create additional resisting stresses that act to prevent runaway acceleration and permit long periods of slow landslide motion (Baum & Johnson, 1993; Booth et

al., 2018; Leshchinsky, 2019). Investigation of tectonic faults and glaciers also shows that slip surface roughness is an important parameter that controls frictional strength (Brodsky et al., 2016; Fang & Dunham, 2013; Meyer et al., 2018).

We found that the mean and maximum active thickness of the individual landslides ranged from ~ 0.40 to ~ 22.4 m and ~ 2.2 to ~ 90 m, respectively. For simplicity we set $f = 1$ indicating that these landslides are effectively blocks sliding on a slope. Yet the borehole data from the Two Towers landslide shows that $f \sim 0.96$. While changing f does not alter the spatial pattern of thickness it does impact the magnitude of the thickness and therefore the volume. Setting $f = 0.96$ would cause a 4% increase in the inferred thickness and volume of each landslide ($h \sim 1/f$) (Table S3). Nonetheless, our findings indicate that most of the sliding surfaces are deep-seated (median thickness for inventory = 5.5 m) and thus are expected to lie within the unweathered Central Belt Franciscan *mélange* bedrock. Recent work by Hahm et al., (2019) found that the Central Belt Franciscan *mélange* has a thin critical zone that is approximately ~ 3 m thick at the ridgetop. Therefore, the slow-moving landslides in the northern California Coast Ranges can be classified as bedrock landslides. However, we note that the Central Belt Franciscan *mélange* is a weak bedrock that has a strength more comparable to weathered rock or soil (Iverson, 2000; Keefer & Johnson, 1983; Nereson et al., 2018; Schulz et al., 2018).

Using our landslide inventory, we developed new volume-area and depth-area geometric scaling relations for slow-moving landslides. Geometric scaling relations are particularly useful for slow-moving landslides because these landslides rarely (if ever) evacuate hillslopes, or create clear scars or deposits that can be easily measured. As a result, most measurements of landslide thickness come from isolated boreholes, which are logistically challenging and expensive to install, and are difficult to extrapolate over an entire landslide. Our results provide best-fit power

function exponents ($\gamma = 1.324$, $\gamma_S = 1.502$, $\gamma_T = 1.453$, $\gamma_C = 1.182$) that are comparable to power function exponents for bedrock and soil landslides (Guzzetti et al., 2009; Larsen et al., 2010). Analysis of a worldwide landslide inventory by Larsen et al (2010) showed that soil landslides had a $\gamma_{soil} \sim 1.1$ -1.3, while bedrock landslides had $\gamma_{bedrock} \sim 1.3$ -1.6. These findings indicate that soil landslides tend to get larger by increasing their planform area rather than by increasing their thickness (and planform area). Larsen et al. (2010) proposed that soil landslides do not continue to get thicker with increasing area because they are limited by the maximum soil depth, which is set by the environmental conditions, while bedrock landslides can extend to depths that far exceed a typical soil cover. Our best-fit depth-area scaling power exponents ($\zeta = 0.3243$, $\zeta_S = 0.5018$, $\zeta_{Ef} = 0.453$, $\zeta_C = 0.1817$) are comparable (with a wide range) to previously published values for slow-moving landslides (Figure 4c). Simoni et al. (2013) reported depth-area scaling $\zeta = 0.44$ from borehole inclinometer data from 23 slow-moving landslides in the Apennine Mountains, Italy. Handwerger et al. (2013) reported depth-area scaling $\zeta = 0.29$ derived from lidar-based estimates of landslide toe thickness from 69 landslides in the Eel River catchment, several of which are also analyzed in this study. Importantly, neither of these studies used large inventories (> 100) or spatially extensive measurements of landslide thickness, which are especially important for slow-moving landslides with variable thicknesses. We note that the large range of scaling exponents suggests that scaling relations should be used with caution. Applying an incorrect scaling exponent to estimate volume for landslides with unknown thickness can lead to large errors in volume calculations (Larsen et al., 2010).

Our findings show that the slow-moving landslides located in the northern California Coast Ranges have geometric scaling exponents that lie in between the soil and bedrock type landslides, providing further evidence that these landslides are hosted in a weak rock that has

mechanical properties in between soil and competent rock (Nereson et al., 2018; Schulz et al., 2018). However, examining the best-fit power function exponents by landslide type suggests that slumps display close to self-similar scaling ($\gamma=1.5$), which is within the range of bedrock landslides, earthflows display scaling intermediate to soil and bedrock landslides, and landslide complexes display scaling most similar to soil landslides. Landslide complexes have similar scaling relations to soil landslides because our inversions show that landslide complexes do not continue to get deeper (on average) as they get larger in planform area. Figure 5 shows that the landslide complexes with the largest areas display a scaling that tends to follow a constant mean depth. This finding has two plausible explanations: 1) the mean landslide depth is limited by a strong layer in the *mélange*, or 2) that landslide complexes are an amalgamation of multiple smaller and shallower landslides. The second explanation provides a reason for why large landslide complexes tend to have multiple kinematic units (e.g., Aryal et al., 2012; Hu et al., 2020).

To further explore the hypothesis that landslide complexes are composed of multiple smaller landslides, we delineated the Boulder Creek landslide complex into 5 smaller sub-landslides and performed a thickness inversion for each sub-landslide (Figure S4). While the thickness patterns are nearly identical to the thickness inversion for the full landslide complex, the magnitude of the inferred thickness differs in some places, and the area of each landslide is smaller, which places them into the space populated by mostly earthflows on the depth-area and volume-area plots (orange circles in Figures 5b and 5d). It is also important to note that some of these differences in the magnitude of the thickness estimate are due to differences in the pixel resolution of the sub-landslides and the full landslide. For the full landslide complex we had to downsample the grid to a 20 m pixel due to computational limitation. The Boulder Creek

landslide complex is the largest landslide in our dataset and is the only landslide that required this downsampling. Nonetheless, mapping landslide complexes as one large landslide results in a low mean thickness relative to the landslide area which affects the geometric scaling relations. While more investigation is warranted, our thickness inversions have caused us to reevaluate how we think about large landslide complexes.

Lastly, the inferred thickness of many of the slow-moving landslides, in particular the landslide complexes, can be highly variable with deeper active zones and thinner or zero thickness areas that are not currently moving (e.g., Figure 3c). It is important to note that these irregular thickness patterns may not align with inferred thickness based on geomorphic or structural interpretations. This discrepancy is likely related to the long-lasting geomorphic imprint that slow-moving landslides have made on the landscape. Landslide surface morphology may last for decades or longer after a landslide completely stops moving (e.g., Booth et al., 2017), which can make it challenging to infer the active landslide thickness without kinematic data. Although our approach is useful for identifying the currently active portions of landslides and inferring their thickness based on volume conservation (with assumptions), it does not allow us to infer the subsurface geometry of the often larger inactive landslide body. As a result, we emphasize the need for more comparisons between ground- and remote sensing-based investigation of landslide geometry. In particular, direct comparison between numerous ground-based measurements from boreholes and structural mapping are needed to widely test the results of our remote sensing approach. Nonetheless, we find our thickness inversions are producing reasonable estimates of landslide thickness in the cases we were able to test (Figures 4, S6, and S7).

4.3 Landslide Basal Stress and Strength

We calculated the shear (or driving) stress and effective normal stress for each landslide under both dry and saturated (hydrostatic pore-water pressure) end member conditions. There is a weak increasing relationship between landslide size and mean effective normal stress. Since the mean shear stress must be balanced by resisting stress, defined here as the effective normal stress times the friction coefficient, our findings imply that the larger landslides are weaker (i.e., lower residual friction angles) than the smaller landslides. We hypothesize that larger landslides are weaker than smaller landslides because of strength heterogeneity in the Franciscan mélange bedrock and the increased likelihood of incorporating weak material within larger volumes. Laboratory measurements of the strength of the Franciscan mélange rocks have shown that the proportion of the blocks hosted in the argillaceous matrix controls the overall rock strength (Roadifer et al., 2009). This implies that larger landslides may have a decreased proportion of blocks and are therefore controlled by the weak argillaceous matrix. Scale-dependent strength has also been observed along other landslides and faults. Brodsky et al. (2016) suggested that faults are weaker at large spatial scales because they encompass larger weak zones. A recent study by Bunn et al., (2020) found that the inferred shear strength of landslides decreases with increasing landslide size. They proposed that smaller landslides were stronger because they occur in cemented cohesive materials and larger landslides were in a residual state. Although we assumed nil cohesion to back-calculate the residual frictional strength of the active landslides, it is likely that cohesion is important in controlling the initial landslide failure due the high clay content of the Central Belt Franciscan mélange (e.g., Milledge et al., 2014).

We also compared our inferred friction angles to lab-based measurements for landslide material and rocks in the Franciscan mélange (Figure 6). The Franciscan mélange rock has a

friction angle that depends on the block-in-matrix proportion (Roadifer et al., 2009). The friction angle of the rock was 30 degrees with 34% block proportion and 40 degrees with 55% block proportion. The Two Towers landslide (Schulz et al., 2018) and Minor Creek landslide (Iverson, 2000), both located in the northern California Coast Ranges have drained residual friction angles of 18.75 and 18 degrees, respectively. The Oak Ridge landslide in the central California Coast ranges has a residual friction angle between 12-14 degrees (Nereson et al., 2018). Many of our inferred friction angles are larger than the values measured in the laboratory for the landslides. These differences may be partially attributed to differences in laboratory- and field-scale measurements (Bunn et al., 2020; Marone, 1998; Van Asch et al., 2007).

Our inversion produces a widespread range of friction angle values from 7 to 53 degrees with a median of 30 degrees for saturated conditions and 16 degrees for dry conditions. Higher apparent friction angles are required for saturated conditions to offset the weakening induced by the elevated pore-water pressure in order to maintain equilibrium conditions (equation 9). Due to the high seasonal rainfall in the northern California Coast Ranges, the slow-moving landslides in California are typically saturated during the wet season and partially saturated during the dry season (Hahm et al., 2019; Iverson & Major, 1987; Schulz et al., 2018). Therefore, the true mean landslide-scale friction angle values likely lie somewhere between our inferred values for saturated and dry conditions. Lastly, although our inferred friction angle values encompass the lab-based values from similar landslides and rocks, the large spread makes it difficult to identify a single representative value for slow-moving landslides in the northern California Coast Ranges. This finding further highlights the heterogeneous nature of the Central Belt Franciscan mélange lithologic unit.

4.4 What Controls the Size of Slow-moving Landslides?

Landslide size is set by the landslide mechanical properties, slope geometry, and environmental conditions. For most landslides, the maximum size is typically limited to the maximum hillslope size, such that the landslide length does not exceed the hillslope length. The landslide thickness is typically set by the location of a weak layer beneath the ground surface, or at a depth where there are changes in strength and permeability, such as the soil to bedrock transition or critical zone to unweathered bedrock transition (Booth, Roering, et al., 2013; Larsen et al., 2010; Milledge et al., 2014). Using a 3D slope stability model for shallow soil landslides that accounts for the forces acting on the landslide basal slip surface, lateral margins, and passive/active wedges at the toe/head, Milledge et al. (2014) found that the critical area and depth that can fail as a landslide depends on the topography, pore-water pressure, and landslide material properties, including density, cohesion, and friction angle. In that model the pore-water pressure plays a fundamental role in determining the critical landslide size and failure depth, such that higher pore-water pressures decrease the critical size required for failure. Large landslides therefore occur when high pore pressures are reached over a correspondingly large spatial area. At our northern California Coast Range study site, the relatively thin, but laterally extensive critical zone that is often saturated during the wet season (Hahm et al., 2019), may promote laterally extensive landslides by elevating the water table height simultaneously over large areas. Milledge et al. (2014)'s model also predicts that landslide thickness should increase as the square root of the landslide area and that the failure depth sets the minimum landslide area. Our best-fit thickness-area scaling exponents for slumps and earthflows are close to a square root scaling (exponents ~ 0.5 with large 95th confidence intervals). Our results also suggest that the landslide thickness controls the minimum area, but does not bound its maximum size. Instead,

slow-moving landslides can continue to grow in area by becoming a landslide complex consisting of multiple, connected, sub-landslides without becoming significantly deeper on average. Large landslide complexes can occupy multiple hillslopes, and fill valleys and catchments such that their size may exceed the typical hillslope size. Thus, it seems that the catchment size sets the maximum area for slow-moving landslides. Our thickness inversion results also indicate that large landslides are weaker than small landslides. This finding may indicate that large landslides become large by incorporating weak material. It is possible that the largest landslides grow over time and take decades to develop (e.g., Mackey & Roering, 2011). As many of our landslide complexes seem to be composed of several smaller sub-landslides or kinematic zones, it is possible that these features have connected through time as slip surfaces propagate along the slope.

5 Conclusions

We measured the 3D surface velocity of more than one hundred slow-moving landslides in the northern California Coast Ranges with data from the NASA/JPL UAVSAR. We used volume conservation techniques to infer the active thickness, volume, stress, and strength of each landslide. The thickness of each landslide is variable and can range from zero to tens of meters sometimes resulting in an irregular slip surface geometry. Volume-area geometric scaling relations suggest that these landslides have similarities to both soil and bedrock landslides around the world. Although their failure planes are likely hosted in unweathered bedrock, their depth seems to be limited, producing a scaling similar to soil landslides for the largest landslide complexes. The inferred residual friction angles are also scale-dependent, like faults, such that large landslides tend to be weaker than small landslides. This decrease in inferred friction angle

with landslide size is likely because larger landslides are composed of larger proportions of weak bedrock. Our study represents the first to use the conservation of volume approach for numerous landslides occurring under the same environmental conditions. Our results provide key insights into the subsurface geometry and stresses that control the behavior of slow-moving landslides. Our work shows how state-of-the-art remote sensing techniques can be used to better understand landslide processes for hazards and to quantify their contribution to landscape evolution.

Acknowledgements

We thank Bill Schulz for sharing data and providing a critical review of this manuscript. We thank Ben Mackey and Georgina Bennett for sharing mapped landslide polygons. Thanks to Isaac Larsen, Benedikt Bayer, and Alessandro Simoni for providing landslide size data. We thank Yang Zheng and the UAVSAR flight and data processing teams for their help with acquiring and processing the data. Part of this research was carried out at the Jet Propulsion Laboratory, California Institute of Technology, under a contract with the National Aeronautics and Space Administration (80NM0018D0004), and supported by the Earth Surface and Interior focus area.

Data Availability

Landslide geometry data used in this study are listed in the references: Larsen et al., (2010), Mackey and Roering, (2011), Simoni et al., (2013) and are included in the figures. Borehole thickness data at the Two Towers landslide is in reference: Schulz et al. (2018). Lidar digital elevation models are provided by OpenTopography and may be downloaded online (<http://www.opentopography.org>). OpenTopography lidar data acquisition and processing was

completed by the National Center for Airborne Laser Mapping (NCALM;
<http://www.ncalm.org>). NCALM funding was provided by NSF's Division of Earth Sciences,
 Instrumentation and Facilities Program EAR-1043051. Topographic data are also provided by
 the German Aerospace Center (DLR) under data proposal DEM GEOL1478 awarded to A. L. H.
 To acquire these data, proposals may be submitted to the DLR online ([https://tandemx-
 science.dlr.de/](https://tandemx-science.dlr.de/)). NASA/JPL UAVSAR data used in this study are freely available and may be
 downloaded through their website (<https://uavsar.jpl.nasa.gov/>).

References

- Aryal, A., Brooks, B. A., Reid, M. E., Bawden, G. W., & Pawlak, G. R. (2012). Displacement fields from
 point cloud data: Application of particle imaging velocimetry to landslide geodesy. *Journal of
 Geophysical Research: Earth Surface* (2003–2012), 117(F1).
<https://doi.org/10.1029/2011JF002161>
- Aryal, A., Brooks, B., & Reid, M. E. (2015). Landslide subsurface slip geometry inferred from 3D surface
 displacement fields. *Geophysical Research Letters*. 42(5), 1411-1417.
<https://doi.org/10.1002/2014GL062688>
- Aster, R. C., Borchers, B., & Thurber, C. H. (2013). Chapter Four - Tikhonov Regularization. In R. C.
 Aster, B. Borchers, & C. H. Thurber (Eds.), *Parameter Estimation and Inverse Problems (Second
 Edition)* (Second Edition, pp. 93–127). Boston: Academic Press. [https://doi.org/10.1016/B978-0-
 12-385048-5.00004-5](https://doi.org/10.1016/B978-0-12-385048-5.00004-5)
- Baum, R. L., & Johnson, A. M. (1993). Steady movement of landslides in fine-grained soils: A model for
 sliding over an irregular slip surface. *US Geological Survey Bulletin (USA)*.
- Bennett, G. L., Roering, J. J., Mackey, B. H., Handwerger, A. L., Schmidt, D. A., & Guillod, B. P.
 (2016). Historic drought puts the brakes on earthflows in Northern California. *Geophysical
 Research Letters*, 43(11), 5725–5731. <https://doi.org/10.1002/2016GL068378>

- Bennett, G. L., Miller, S. R., Roering, J. J., & Schmidt, D. A. (2016). Landslides, threshold slopes, and the survival of relict terrain in the wake of the Mendocino Triple Junction. *Geology*, 44(5), 363–366. <https://doi.org/10.1130/G37530.1>
- Bessette-Kirton, E. K., Coe, J. A., & Zhou, W. (2018). Using stereo satellite imagery to account for ablation, entrainment, and compaction in volume calculations for rock avalanches on glaciers: Application to the 2016 Lamplugh rock avalanche in Glacier Bay National Park, Alaska. *Journal of Geophysical Research: Earth Surface*, 123(4), 622–641. <https://doi.org/10.1002/2017JF004512>
- Booth, A. M., Lamb, M. P., Avouac, J.-P., & Delacourt, C. (2013). Landslide velocity, thickness, and rheology from remote sensing: La Clapière landslide, France. *Geophysical Research Letters*, 40(16), 4299–4304. <https://doi.org/10.1002/grl.50828>
- Booth, A. M., Roering, J. J., & Rempel, A. W. (2013). Topographic signatures and a general transport law for deep-seated landslides in a landscape evolution model. *Journal of Geophysical Research: Earth Surface*, 118(2), 603–624. <https://doi.org/10.1002/jgrf.20051>
- Booth, A. M., LaHusen, S. R., Duvall, A. R., & Montgomery, D. R. (2017). Holocene history of deep-seated landsliding in the North Fork Stillaguamish River valley from surface roughness analysis, radiocarbon dating, and numerical landscape evolution modeling. *Journal of Geophysical Research: Earth Surface*, 122(2), 456–472. <https://doi.org/10.1002/2016JF003934>
- Booth, A. M., McCarley, J., Hinkle, J., Shaw, S., Ampuero, J.-P., & Lamb, M. P. (2018). Transient Reactivation of a Deep-Seated Landslide by Undrained Loading Captured With Repeat Airborne and Terrestrial Lidar. *Geophysical Research Letters*, 45(10), 4841–4850. <https://doi.org/10.1029/2018GL077812>
- Booth, A. M., McCarley, J. C., & Nelson, J. (2020). Multi-year, three-dimensional landslide surface deformation from repeat lidar and response to precipitation: Mill Gulch earthflow, California. *Landslides*, 1–14. <https://doi.org/10.1007/s10346-020-01364-z>
- Brodsky, E. E., Kirkpatrick, J. D., & Candela, T. (2016). Constraints from fault roughness on the scale-dependent strength of rocks. *Geology*, 44(1), 19–22. <https://doi.org/10.1130/G37206.1>

- Bunn, M., Leshchinsky, B., & Olsen, M. J. (2020). Geologic Trends in Shear Strength Properties Inferred through Three-Dimensional Back-Analysis of Landslide Inventories. *Journal of Geophysical Research: Earth Surface*, e2019JF005461. <https://doi.org/10.1029/2019JF005461>
- Cerovski-Darriau, C., & Roering, J. J. (2016). Influence of anthropogenic land-use change on hillslope erosion in the Waipaoa River Basin, New Zealand. *Earth Surface Processes and Landforms*, 41(15), 2167–2176. <https://doi.org/10.1002/esp.3969>
- Coe, J. A., McKenna, J. P., Godt, J. W., & Baum, R. L. (2009). Basal-topographic control of stationary ponds on a continuously moving landslide. *Earth Surface Processes and Landforms*, 34(2), 264–279. <https://doi.org/10.1002/esp.1721>
- Cruden, D. M., & Varnes, D. J. (1996). Landslides: Investigation and Mitigation. Chapter 3-Landslide types and processes. *Transportation Research Board Special Report*, (247).
- Delbridge, B. G., Bürgmann, R., Fielding, E., Hensley, S., & Schulz, W. H. (2016). Three-dimensional surface deformation derived from airborne interferometric UAVSAR: Application to the Slumgullion Landslide. *Journal of Geophysical Research: Solid Earth*, 121(5), 3951–3977. <https://doi.org/10.1002/2015JB012559>
- Fang, Z., & Dunham, E. M. (2013). Additional shear resistance from fault roughness and stress levels on geometrically complex faults. *Journal of Geophysical Research: Solid Earth*, 118(7), 3642–3654. <https://doi.org/10.1002/jgrb.50262>
- Fialko, Y., Simons, M., & Agnew, D. (2001). The complete (3-D) surface displacement field in the epicentral area of the 1999 Mw7. 1 Hector Mine earthquake, California, from space geodetic observations. *Geophysical Research Letters*, 28(16), 3063–3066. <https://doi.org/10.1029/2001GL013174>
- Fielding, E. J., Liu, Z., Stephenson, O. L., Zhong, M., Liang, C., Moore, A., et al. (2020). Surface Deformation Related to the 2019 M w 7.1 and 6.4 Ridgecrest Earthquakes in California from GPS, SAR Interferometry, and SAR Pixel Offsets. *Seismological Research Letters*. <https://doi.org/10.1785/0220190302>

- Grant, M., & Boyd, S. (2014). *CVX: Matlab Software for Disciplined Convex Programming, version 2.1*. Retrieved from <http://cvxr.com/cvx>
- Guerriero, L., Coe, J. A., Revellino, P., Grelle, G., Pinto, F., & Guadagno, F. M. (2014). Influence of slip-surface geometry on earth-flow deformation, Montaguto earth flow, southern Italy. *Geomorphology*, 219(0), 285 – 305. <http://dx.doi.org/10.1016/j.geomorph.2014.04.039>
- Guerriero, L., Bertello, L., Cardozo, N., Berti, M., Grelle, G., & Revellino, P. (2017). Unsteady sediment discharge in earth flows: A case study from the Mount Pizzuto earth flow, southern Italy. *Geomorphology*, 295, 260–284. <https://doi.org/10.1016/j.geomorph.2017.07.011>
- Guzzetti, F., Ardizzone, F., Cardinali, M., Rossi, M., & Valigi, D. (2009). Landslide volumes and landslide mobilization rates in Umbria, central Italy. *Earth and Planetary Science Letters*, 279(3–4), 222–229. <https://doi.org/10.1016/j.epsl.2009.01.005>
- Hahm, W. J., Rempe, D. M., Dralle, D. N., Dawson, T. E., Lovill, S. M., Bryk, A. B., et al. (2019). Lithologically controlled subsurface critical zone thickness and water storage capacity determine regional plant community composition. *Water Resources Research*, 55(4), 3028–3055. <https://doi.org/10.1029/2018WR023760>
- Handwerger, A. L., Roering, J. J., & Schmidt, D. A. (2013). Controls on the seasonal deformation of slow-moving landslides. *Earth and Planetary Science Letters*, 377, 239–247. <https://doi.org/10.1016/j.epsl.2013.06.047>
- Handwerger, A. L., Roering, J. J., Schmidt, D. A., & Rempel, A. W. (2015). Kinematics of earthflows in the Northern California Coast Ranges using satellite interferometry. *Geomorphology*, 246, 321–333. <https://doi.org/10.1016/j.geomorph.2015.06.003>
- Handwerger, A. L., Huang, M.-H., Fielding, E. J., Booth, A. M., & Bürgmann, R. (2019). A shift from drought to extreme rainfall drives a stable landslide to catastrophic failure. *Scientific Reports*, 9(1), 1569. <https://doi.org/10.1038/s41598-018-38300-0>
- Handwerger, A. L., Fielding, E. J., Huang, M.-H., Bennett, G. L., Liang, C., & Schulz, W. H. (2019). Widespread initiation, reactivation, and acceleration of landslides in the northern California Coast

- Ranges due to extreme rainfall. *Journal of Geophysical Research: Earth Surface*, 124(7), 1782–1797. <https://doi.org/10.1029/2019JF005035>
- Hu, X., Bürgmann, R., Schulz, W. H., & Fielding, E. J. (2020). Four-dimensional surface motions of the Slumgullion landslide and quantification of hydrometeorological forcing. *Nature Communications*, 11(1), 1–9. <https://doi.org/10.1038/s41467-020-16617-7>
- Hungr, O. (1987). An extension of Bishop's simplified method of slope stability analysis to three dimensions. *Geotechnique*, 37(1), 113–117. <https://doi.org/10.1680/geot.1987.37.1.113>
- Hungr, O., Salgado, F., & Byrne, P. (1989). Evaluation of a three-dimensional method of slope stability analysis. *Canadian Geotechnical Journal*, 26(4), 679–686. <https://doi.org/10.1139/t89-079>
- Hungr, O., Leroueil, S., & Picarelli, L. (2014). The Varnes classification of landslide types, an update. *Landslides*, 11(2), 167–194. <https://doi.org/10.1007/s10346-013-0436-y>
- Intrieri, E., Raspini, F., Fumagalli, A., Lu, P., Del Conte, S., Farina, P., et al. (2017). The Maoxian landslide as seen from space: detecting precursors of failure with Sentinel-1 data. *Landslides*, 1–11. <https://doi.org/10.1007/s10346-017-0915-7>
- Iverson, R. M. (2000). Landslide triggering by rain infiltration. *Water Resources Research*, 36(7), 1897–1910. <https://doi.org/10.1029/2000WR900090>
- Iverson, R. M. (2005). Regulation of landslide motion by dilatancy and pore pressure feedback. *Journal of Geophysical Research: Earth Surface*, 110(F2). <https://doi.org/10.1029/2004JF000268>
- Iverson, R. M., & Major, J. J. (1987). Rainfall, ground-water flow, and seasonal movement at Minor Creek landslide, northwestern California: Physical interpretation of empirical relations. *Geological Society of America Bulletin*, 99(4), 579–594. [https://doi.org/10.1130/0016-7606\(1987\)99<579:RGFASM>2.0.CO;2](https://doi.org/10.1130/0016-7606(1987)99<579:RGFASM>2.0.CO;2)
- Jayko, A., Blake, M., McLaughlin, R., Ohlin, H., Ellen, S., & Kelsey, H. (1989). Reconnaissance Geologic Map of the Covelo 30-by 60-Minute Quadrangle. *Northern California: US Geological Survey Miscellaneous Field Investigation Map MF-2001*, Scale, 1(100), 000. <https://doi.org/10.3133/mf2001>

- Jennings, C. W., Strand, R. G., & Rogers, T. H. (1977). Geologic map of California: California Division of Mines and Geology, scale 1:750,000.
- Jung, J., & Yun, S.-H. (2020). Evaluation of Coherent and Incoherent Landslide Detection Methods Based on Synthetic Aperture Radar for Rapid Response: A Case Study for the 2018 Hokkaido Landslides. *Remote Sensing*, 12(2), 265. <https://doi.org/10.3390/rs12020265>
- Kasper, van W., John A, S., William, N., & Luis, T. (2002). Data and model uncertainty estimation for linear inversion. *Geophysical Journal International*, 149(3), 625–632. <https://doi.org/10.1046/j.1365-246X.2002.01660.x>
- Keefer, D. K., & Johnson, A. M. (1983). Earth flows: Morphology, mobilization, and movement. Washington: United States Government Printing Office. No 1264. <https://doi.org/10.3133/pp1264>
- Kelsey, H. M. (1978). Earthflows in Franciscan melange, Van Duzen River basin, California. *Geology*, 6(6), 361–364. [https://doi.org/10.1130/0091-7613\(1978\)6<361:EIFMVD>2.0.CO;2](https://doi.org/10.1130/0091-7613(1978)6<361:EIFMVD>2.0.CO;2)
- Korup, O., Clague, J. J., Hermanns, R. L., Hewitt, K., Strom, A. L., & Weidinger, J. T. (2007). Giant landslides, topography, and erosion. *Earth and Planetary Science Letters*, 261(3–4), 578–589. <https://doi.org/10.1016/j.epsl.2007.07.025>
- Lacroix, P., Dehecq, A., & Taipe, E. (2020). Irrigation-triggered landslides in a Peruvian desert caused by modern intensive farming. *Nature Geoscience*, 13(1), 56–60. <https://doi.org/10.1038/s41561-019-0500-x>
- Lacroix, P., Handwerger, A. L., & Bièvre, G. (2020). Life and death of slow-moving landslides. *Nature Reviews Earth & Environment*, 1–16. <https://doi.org/10.1038/s43017-020-0072-8>
- Larsen, I. J., Montgomery, D. R., & Korup, O. (2010). Landslide erosion controlled by hillslope material. *Nature Geoscience*, 3(4), 247–251. <https://doi.org/10.1038/ngeo776>
- Legros, F. (2002). The mobility of long-runout landslides. *Engineering Geology*, 63(3–4), 301–331. [https://doi.org/10.1016/S0013-7952\(01\)00090-4](https://doi.org/10.1016/S0013-7952(01)00090-4)
- Leshchinsky, B. (2019). Quantifying the influence of failure surface asperities on the basal shear resistance of translational landslides. *Landslides*, 16(7), 1375–1383.

<https://doi.org/10.1007/s10346-019-01185-9>

- Mackey, B. H., & Roering, J. J. (2011). Sediment yield, spatial characteristics, and the long-term evolution of active earthflows determined from airborne LiDAR and historical aerial photographs, Eel River, California. *Geological Society of America Bulletin*, 123(7–8), 1560–1576. <https://doi.org/10.1130/B30306.1>
- Mackey, B. H., Roering, J. J., & McKean, J. (2009). Long-term kinematics and sediment flux of an active earthflow, Eel River, California. *Geology*, 37(9), 803–806. <https://doi.org/10.1130/G30136A.1>
- Madson, A., Fielding, E., Sheng, Y., & Cavanaugh, K. (2019). High-resolution spaceborne, airborne and in situ landslide kinematic measurements of the slumgullion landslide in Southwest Colorado. *Remote Sensing*, 11(3), 265. <https://doi.org/10.3390/rs11030265>
- Malet, J.-P., Maquaire, O., & Calais, E. (2002). The use of Global Positioning System techniques for the continuous monitoring of landslides: application to the Super-Sauze earthflow (Alpes-de-Haute-Provence, France). *Geomorphology*, 43(1–2), 33–54. [https://doi.org/10.1016/S0169-555X\(01\)00098-8](https://doi.org/10.1016/S0169-555X(01)00098-8)
- Marone, C. (1998). Laboratory-derived friction laws and their application to seismic faulting. *Annual Review of Earth and Planetary Sciences*, 26(1), 643–696. <https://doi.org/10.1146/annurev.earth.26.1.643>
- McLaughlin, R. J., Kling, S. A., Poore, R. Z., McDougall, K., & Beutner, E. C. (1982). Post-middle Miocene accretion of Franciscan rocks, northwestern California. *Geological Society of America Bulletin*, 93(7), 595–605. [https://doi.org/10.1130/0016-7606\(1982\)93<595:PMAOFR>2.0.CO;2](https://doi.org/10.1130/0016-7606(1982)93<595:PMAOFR>2.0.CO;2)
- McLaughlin, R. J., Blake, S., Jayko, M., Irwin, A., Aalto, W., Carver, K., et al. (2000). Geologic map of the Cape Mendocino, Eureka, Garberville, and southwestern part of the Hayfork 30 X 60 Quadrangles and Adjacent Offshore Area, Northern California.
- Meyer, C. R., Downey, A. S., & Rempel, A. W. (2018). Freeze-on limits bed strength beneath sliding glaciers. *Nature Communications*, 9(1), 1–6. <https://doi.org/10.1038/s41467-018-05716-1>
- Michel, J., Dario, C., Marc-Henri, D., Thierry, O., Marina, P. I., & Benjamin, R. (2020). A review of

- methods used to estimate initial landslide failure surface depths and volumes. *Engineering Geology*, 267, 105478. <https://doi.org/10.1016/j.enggeo.2020.105478>
- Milledge, D. G., Bellugi, D., McKean, J. A., Densmore, A. L., & Dietrich, W. E. (2014). A multidimensional stability model for predicting shallow landslide size and shape across landscapes. *Journal of Geophysical Research: Earth Surface*, 119(11), 2481–2504. <https://doi.org/10.1002/2014JF003135>
- Nereson, A. L., Davila Olivera, S., & Finnegan, N. J. (2018). Field and Remote-Sensing Evidence for Hydro-mechanical Isolation of a Long-Lived Earthflow in Central California. *Geophysical Research Letters*, 45(18), 9672–9680. <https://doi.org/10.1029/2018GL079430>
- Nereson, A. L., & Finnegan, N. J. (2018). Drivers of earthflow motion revealed by an 80 yr record of displacement from Oak Ridge earthflow, Diablo Range, California, USA. *Geological Society of America Bulletin*. 131(3-4), 389-402. <https://doi.org/10.1130/B32020.1>
- Pathier, E., Fielding, E. J., Wright, T. J., Walker, R., Parsons, B. E., & Hensley, S. (2006). Displacement field and slip distribution of the 2005 Kashmir earthquake from SAR imagery. *Geophysical Research Letters*, 33(20). <https://doi.org/10.1029/2006GL027193>
- Roadifer, J. W., Forrest, M. P., & Lindquist, E. S. (2009). Evaluation of shear strength of melange foundation at Calaveras Dam. Proceedings of U. S. Society for Dams, Annual Meeting and Conference, 29th, on "Managing Our Water Retention Systems, ", 507–521.
- Roering, J.J (2012). Eel River, CA: Landsliding and the Evolution of Mountainous Landscapes in collaboration with National Center for Airborne Laser Mapping (NCALM), distributed by OpenTopography. <https://doi.org/10.5069/G9XS5S9P>
- Roering, J. J., Stimely, L. L., Mackey, B. H., & Schmidt, D. A. (2009). Using DInSAR, airborne LiDAR, and archival air photos to quantify landsliding and sediment transport. *Geophysical Research Letters*, 36(19). <https://doi.org/10.1029/2009GL040374>
- Roering, J. J., Mackey, B. H., Handwerger, A. L., Booth, A. M., Schmidt, D. A., Bennett, G. L., & Cerovski-Darriau, C. (2015). Beyond the angle of repose: A review and synthesis of landslide

- processes in response to rapid uplift, Eel River, Northern California. *Geomorphology*, 236, 109–131. <https://doi.org/10.1016/j.geomorph.2015.02.013>
- Rosen, P. A., Gurrola, E., Sacco, G. F., & Zebker, H. (2012). The InSAR scientific computing environment. In *Synthetic Aperture Radar, 2012. EUSAR. 9th European Conference on* (pp. 730–733). VDE.
- Rutter, E., & Green, S. (2011). Quantifying creep behaviour of clay-bearing rocks below the critical stress state for rapid failure: Mam Tor landslide, Derbyshire, England. *Journal of the Geological Society*, 168(2), 359–372. <https://doi.org/10.1144/0016-76492010-133>
- Schulz, W. H., Coe, J. A., Ricci, P. P., Smoczyk, G. M., Shurtleff, B. L., & Panosky, J. (2017). Landslide kinematics and their potential controls from hourly to decadal timescales: Insights from integrating ground-based InSAR measurements with structural maps and long-term monitoring data. *Geomorphology*, 285, 121–136. <https://doi.org/10.1016/j.geomorph.2017.02.011>
- Schulz, W. H., Smith, J. B., Wang, G., Jiang, Y., & Roering, J. J. (2018). Clayey landslide initiation and acceleration strongly modulated by soil swelling. *Geophysical Research Letters*, 45(4), 1888–1896. <https://doi.org/10.1002/2017GL076807>
- Simoni, A., Ponza, A., Picotti, V., Berti, M., & Dinelli, E. (2013). Earthflow sediment production and Holocene sediment record in a large Apennine catchment. *Geomorphology*, 188, 42–53. <https://doi.org/10.1016/j.geomorph.2012.12.006>
- Stumpf, A., Malet, J.-P., & Delacourt, C. (2017). Correlation of satellite image time-series for the detection and monitoring of slow-moving landslides. *Remote Sensing of Environment*, 189, 40–55. <https://doi.org/10.1016/j.rse.2016.11.007>
- Travelletti, J., & Malet, J.-P. (2012). Characterization of the 3D geometry of flow-like landslides: A methodology based on the integration of heterogeneous multi-source data. *Engineering Geology*, 128, 30–48. <https://doi.org/10.1016/j.enggeo.2011.05.003>
- Travelletti, J., Malet, J.-P., & Delacourt, C. (2014). Image-based correlation of Laser Scanning point cloud time series for landslide monitoring. *International Journal of Applied Earth Observation*

- 973 *and Geoinformation*, 32, 1–18. <https://doi.org/10.1016/j.jag.2014.03.022>
- 974 Van Asch, T. W., Van Beek, L., & Bogaard, T. (2007). Problems in predicting the mobility of slow-
- 975 moving landslides. *Engineering Geology*, 91(1), 46–55.
- 976 <https://doi.org/10.1016/j.enggeo.2006.12.012>
- 977 Warrick, J. A., Ritchie, A. C., Schmidt, K. M., Reid, M. E., & Logan, J. (2019). Characterizing the
- 978 catastrophic 2017 Mud Creek landslide, California, using repeat structure-from-motion (SfM)
- 979 photogrammetry. *Landslides*, 1–19. <https://doi.org/10.1007/s10346-019-01160-4>
- 980 Wartman, J., Montgomery, D. R., Anderson, S. A., Keaton, J. R., Benoît, J., dela Chapelle, J., & Gilbert,
- 981 R. (2016). The 22 March 2014 Oso landslide, Washington, USA. *Geomorphology*, 253, 275–288.
- 982 <https://doi.org/10.1016/j.geomorph.2015.10.022>