

1 2 **ENSO and SAM influence on the generation of long episodes of Rossby Wave** 3 **Packets during Southern Hemisphere Summer**

4 **Iago Pérez¹ , Marcelo Barreiro² , and Cristina Masoller³ **

5 ¹Departamento de Ciencias de la atmósfera, Facultad de Ciencias, Universidad de la República,
6 Montevideo, Uruguay, ²Physics department-Nonlinear, Universitat Politècnica de Catalunya, Rambla St.
7 Nebridi 22, Terrasa 08222, Barcelona, Spain

8 Corresponding author: Iago Pérez (iperez@fisica.edu.uy)

9 **Key Points:**

- 10 • Occurrence of long-lived Wave Packets shows large interannual variability in southern hemisphere
11 summer.
- 12 • El Niño and negative Southern Annual Mode phases create atmospheric conditions that favor long-lived
13 Rossby Wave Packets.
- 14 • Years with largest numbers of long-lived Rossby Wave Packets present a zonally symmetric and narrow
15 upper level jet shifted northward.

16 **Abstract**

17 This study aims at understanding the impact of low-frequency climate modes on Rossby Wave
18 Packets (RWPs) during the southern hemisphere summer. In particular, we focus on long-lived
19 RWPs (lifespan above 8 days) and determine how El Niño-Southern Oscillation (ENSO) and
20 Southern Annular Mode (SAM) modulate their frequency of occurrence plus the main areas of
21 detection and dissipation. We find that the occurrence of long lived RWPs is maximum during El
22 Niño years and negative SAM events. Years with largest numbers of long-lived RWPs are
23 characterized by a zonally symmetric and narrow upper level jet that is shifted northward from
24 its climatological position. Conversely, when the jet is shifted southward, as during positive
25 SAM phases, particularly in the southwestern Pacific basin, the number of long-lived RWPs
26 detected diminishes. El Niño sets atmospheric conditions that support the formation of long lived
27 RWPs whereas La Niña years presents high interannual variability in the frequency of
28 occurrence. Moreover, during El Niño events the main formation area is between 61-120°E and
29 its main dissipation area between 300-359°E. During La Niña events, the main formation area
30 moves to 241-300°E and no main dissipation area is identified. During positive SAM two main
31 formation areas appear at 61-120°E and 241-300°E and two main dissipation areas between 121-
32 180° and 301-359°, whereas in negative SAM only one formation area at 241-300° is detected
33 and no main dissipation area is found.

34 **Keywords:** long-lived RWPs, ENSO, SAM.

35 **1 Introduction**

36 Atmospheric circulation in mid-latitudes is dominated by the upper level jet and associated storm
37 track. Predictability on time scales longer than the synoptic is limited and on seasonal time scales
38 is directly dependent on tropical surface ocean conditions. On intermediate, subseasonal scales
39 predictability relies on persistent structures such as blockings or Rossby Wave Packets (RWPs).

40 RWPs are packets of upper-level atmospheric waves that are able to travel coherently for several
41 days by downstream development (Chang and Yu 1999; Chang 2000), and transport large
42 quantities of energy. These packets are constantly being created and destroyed in mid-latitudes
43 and most of them survive only for a few days, but under certain conditions these packets can
44 achieve such stability that they live and propagate for several days or even weeks before
45 disappearing (Grazzini and Vitart 2015).

46
47 RWPs represent high-amplitude meanderings of the jet stream and thus are related to storm track
48 variability (Souders et al., 2014a). Moreover they have been referred as precursors of extreme
49 climatological events such as heat waves or droughts, (Chang 2005; Obrien and Reeder 2017;
50 Wirth et al., 2018), as well as extratropical cyclone development, (Chang et al 2005; Sagarra and
51 Barreiro 2019). Additionally, Souders et al., (2014b) remarked that if they are not well
52 represented in models, RWPs can also affect the weather forecasts in the areas they cross,
53 increasing the uncertainty on short and middle range forecasting. Increasing our understanding of
54 processes that control RWPs is thus crucial to understand the mechanisms that governs weather
55 and climate in mid-latitudes, and provides the possibility of extending the forecast beyond
56 synoptic time scales (Grazzini and Vitart 2015).

57
58 As mentioned before, under certain conditions RWPs can maintain their coherence longer than
59 usual and are able to survive for weeks in the atmosphere. The lifespan, extension and
60 propagation of RWPs are highly dependent on the potential vorticity gradient and the distribution
61 of adiabatic heating sources (Grazzini and Vitart 2015). Potential vorticity (PV) gradients
62 controls the strength and position of the waveguide where the RWP propagates (Hoskins and
63 Ambrizzi, 1993) in such a way that a very intense and narrow gradient of PV favors the
64 development of coherent RWPs that will last longer in the atmosphere before they disappear
65 (Chang and Yu 1999; Souders et al., 2014b) while weaker gradients tend to stop or damp wave
66 propagation (Grazzini and Vitart 2015). RWPs are more easily detected in the Southern
67 hemisphere due to the presence of baroclinically unfavorable continental areas (Grazzini and
68 Vitart 2015). Moreover, during summertime the storm track is almost zonally symmetrical at
69 50°S, favoring the propagation of packets (Chang 1999).

70
71 RWPs have been extensively studied in recent years in the northern hemisphere (NH). However
72 in the southern hemisphere (SH) there are fewer studies and mainly focused on climatological
73 statistics. Their interannual variability and the influence of global climate modes have not been
74 addressed in detail. Here we will study the impact of El Niño-Southern Oscillation (ENSO) and
75 the Southern Annular Mode (SAM) on the characteristics of RWPs during summertime. It is well
76 known that these climate modes alter the large scale mid-latitude flow and therefore the
77 waveguide where the RWPs propagate. Moreover, we will focus on long-lived RWPs, that is,
78 those packets that live longer than 8 days (Grazzini and Lucarini 2010). Due to their link with
79 extreme climatological events, the study of long-lived RWPs and their areas of

80 formation/detection may contribute to the improvement of extended range forecasting of extreme
81 events (Grazzini 2007; Wirth et al 2018). Next we describe the main results found in the
82 literature about for RWPs in the SH.

83 Chang (1999) first describes the properties of upper level wave packets by means of time-lagged
84 one point correlation maps of the meridional velocity. He found that wave packets tend to
85 propagate along the mid-latitude jets and that their zonal group velocity in the SH summer is
86 about 20-25 m/s. Chang (2000) showed that most wave packets propagate dominated by
87 downstream development, and that a surface cyclone developed in nearly all cases to the east of
88 the upper-level troughs associated with the wave packets.

89 Souders et al., (2014b) found that in the SH, RWPs have a mean lifetime of 7.9 days and
90 travelled a mean of 151 degrees. In contrast, during NH summer RWPs last for about 6 days and
91 travel only 119 degrees. Additionally, RWPs activity in the SH does not show a strong seasonal
92 cycle and the packets detected are faster and more coherent than in the NH. They also found a
93 positive trend in the annual mean activity of RWPs that they hypothesize may be related to an
94 improvement of the quality of the reanalysis (see e.g Barreiro et al 2014) or to the observed trend
95 in the SAM during summer.

96 Barreiro (2017) studied the interannual variability of wave activity during summertime in the
97 southeast Pacific-Atlantic sector and found that the leading pattern of variability in this region is
98 correlated with ENSO such that there is an increase in transient wave energy in the Pacific
99 during El Niño years. However, this study did not address the characteristics of RWPs.

100 Sagarra and Barreiro (2020) performed a climatological study of RWPs during the austral
101 summer in the SH. They found a mean of 32 packets per season, 90% of the trajectories have a
102 lifespan of 3-8 days and only 2% above 14 days. In addition 80% of the waves propagated
103 between 30-170° degrees of longitude and only 2% surpasses 360 degrees with a median of 107°.
104 No main area of dissipation/formation on seasonal or monthly time scales was detected. The
105 study did not find a relation between the interannual count of RWPs and the ENSO, but
106 suggested a possible relationship with the SAM (Southern Annular Mode).

107 In this study we use a similar methodology as in Sagarra and Barreiro (2020) but, as mentioned
108 above, focus on the impact of ENSO and SAM on the RWPs whose lifespan is longer than 8
109 days.

110 The manuscript is organized as follows. Chapter 2 describes the data and the algorithm used.
111 Chapter 3 describes the climatology and interannual variability of RWPs as well as the impact of
112 ENSO and SAM on long-lived packets. Lastly Chapter 4 provides a summary of the results.

113

114 **2. Data and methods**

115

116 **2.1 Data**

117

118 To track RWPs we use daily mean meridional winds (m/s) at 300 hPa as done previously by
119 several authors (e.g. Chang and Yu 1999; Sagarra and Barreiro 2019). This data comes from

120 NCEP/DOE Reanalysis 2 (NOAA/OAR/ESRL) available in <http://www.esrl.noaa.gov/psd/>
121 (Kanamitsu et al., 2002) with an horizontal spatial resolution of $2.5 \times 2.5^\circ$. The period of study
122 chosen is the austral summer here considered as December, January, February and March
123 (DJFM) from 1 December of 1979 to 31 March of 2020, retaining data corresponding to 41
124 summer seasons. After subtracting the daily climatology we calculated the amplitude of the
125 wave packet envelope (m/s) using the method of Zimin et al., (2003) retaining wave numbers
126 between 4 and 11 which are representative of the atmospheric transients in the SH (Trenberth
127 1981). For the methodology of Zimin et al., (2003) to work well the propagation of RWPs has to
128 be in the zonal direction (e.g. Zimin et al., 2006), which is the case for the SH summer as shown
129 by Chang (1999). This latter study also showed that the jet and maximum variance in meridional
130 wind anomalies are in a band centered at 50°S , and therefore in our study we averaged
131 latitudinally data in the band 40°S - 65°S .

132
133 To characterize the interannual occurrence of the global climate modes we considered the ONI
134 index (Oscillation Niño Index) in the case of ENSO and the AAO index (Antarctic Oscillation)
135 for SAM, both datasets available in the NOAA website <https://origin.cpc.ncep.noaa.gov/>. Of the
136 41 seasons considered 25 correspond to positive SAM and 16 to negative SAM, which is
137 consistent with the observed positive trend towards positive phases in SAM since the mid-90s. In
138 the case of ENSO, 14 years are classified as “El Niño”, 13 years as “La Niña” and 14 years as
139 Neutral events. We also used geopotential height at 300 hPa (Z300) from NCEP/DOE Reanalysis
140 2 and sea surface temperature (SST) from the ERSSTv5 dataset (Huang et al., 2017) to study the
141 upper level circulation anomalies and surface ocean conditions, respectively associated with the
142 interannual variability in the occurrence of RWPs.

143
144 Finally, to characterize changes in the position and intensity of the waveguide we calculated the
145 meridional gradient of the absolute vorticity (AV) that can be used as an approximation of the
146 isentropic gradients of potential vorticity (Wirth et al., 2018). This was done using 300hPa
147 meridional and zonal monthly mean winds from the NCEP/DOE 2 Reanalysis.

148

149 **2.2 Methodology**

150

151 **2.2.1 Tracking algorithm**

152

153 Previously to the application of the algorithm to track RWPs, it is necessary to set a minimum
154 threshold for the amplitude of the envelope to avoid tracking noise. However the value of this
155 threshold is not well-defined due to the lack of intrinsic physical properties that distinguish one
156 packet from another (Souders et al., 2014b). As a result, a low threshold tracks noise and a large
157 one will miss information. Grazzini and Vitart (2015) considered a threshold value of 16 m/s,
158 Souders et al., (2014b) 14 m/s and Sagarra and Barreiro (2020) 15 m/s. In this study we
159 considered a threshold of 15 m/s in the calibration of the algorithm and then we analysed the

160 robustness of our findings by considering higher and lower threshold values of 13, 15 and 17
161 m/s.

162
163 The algorithm tracks the amplitude of the envelope above the chosen threshold allowing to
164 follow the trajectory of a coherent wave packet in its trajectory (Grazzini and Vitart 2015;
165 Sagarra and Barreiro 2020). The algorithm consists of following steps:

166 1.-Detect of the highest value of the amplitude in the longitudinal axis (p_n) in the first day of
167 the data matrix.

168 2.- Search for the position of the maximum amplitude eastwards the next day (p_{n+1})

169 3.-If the points p_n and p_{n+1} are between [p_n+15° to p_n+45°], p_n and p_{n+1} are registered as part of
170 the same trajectory and we repeat steps 2 and 3 for the following days. The distance is chosen so
171 that the wave packet travels at a speed between 15-45° per day as found in the literature (Sagarra
172 and Barreiro 2020).

173 4.- When we find a maximum p_m that it is not within the range specified in step 3 or when we
174 reach the limit of the data matrix, the trajectory is saved and we resume the tracking from the
175 last day we detected the start of a trajectory.

176 5.-After analyzing all longitudes for the day, we proceed to the next one and repeat steps 1-4
177 until we analysed the whole data matrix.

178 6.-After finishing the analysis of the data matrix we check for trajectories that may have been
179 truncated and apply proximity criteria shown in Table 1 to link them. That is if two trajectories
180 are closer than 1000 km and their difference in slope are less than 20°/day they are joined.
181 Additionally, for slow trajectories that were separated by 1-2 days but the mean speed between
182 the end of a trajectory and the start of the next one is less than 15°/day, they were considered as a
183 single trajectory if the mean speed between points L-1 or L (being L the last point of the first
184 trajectory) to I or I+1 (being I the initial point of the second trajectory) is 15°/day or above and
185 fulfilled the rest of the conditions set in Table 1.

186 7.-Finally, we filtered out trajectories with a duration shorter than three days because they are
187 not relevant to this study.

188

Minimum threshold	15 m/s
Distance	1000 km
Slopes differences	20°/day
Mean speed per day	[15-45 °/day]
Days separated	1-2 days

189 Table 1: Proximity criteria used in the tracking algorithm

190

191 After running the algorithm, we performed a statistical analysis of the RWPs detected. The
192 following measures are considered: (i) mean propagation speed (m/s) taking into consideration
193 that 1°/day is approximately 0.82 m/s in the latitudinal region analysed (40-65°S), (ii) duration
194 (days during which a trajectory can be tracked until it disappears), (iii) length of the trajectory

195 extension (difference between the starting and ending point of the trajectory) in degrees and (iv)
196 areas of formation and dissipation, being formation the area where the RWPs surpasses the
197 minimum envelope threshold so the packet is sufficiently coherent to be tracked and dissipation
198 the last point recorded of the trajectory. Then, we studied the interannual variability of wave
199 packets and the impact of global climate modes.

200

201 **2.2.2 Classification of the identified trajectories**

202

203 RWPs are classified into three categories based on their lifespan: short-lived RWPs (SL) if their
204 lifespan ranges between 3-6 days, medium-lived RWPs (ML) when their lifetime ranges between
205 7-8 days and long-lived RWPs (LL) when they last more than 8 days. Although this study is
206 focused on long-lived packets we also registered RWPs of different lifespans to study whether
207 their frequency of occurrence is correlated. As pointed out above we considered several
208 thresholds and as expected we found a decreasing number of LL RWPs with increasing
209 threshold, so that for 17m/s the algorithm detects very few RWPs.

210

211 **2.2.3 Identification of the formation and dissipation areas and correlation with ENSO and 212 SAM events.**

213

214 To study the formation and dissipation areas we considered the following 6 regions in order to
215 condense the information: 0-60°E,61-120°E,121-180°E,181-240°E,241-300°E and 301-359°E.

216

217 A linear regression and correlation analysis between the time series of interannual occurrence of
218 RWPs with Z300 and SST anomalies was conducted. This allows to determine the large scale
219 circulation anomalies and oceanic conditions that accompany changes in the frequency of
220 occurrence of RWPs, which may be in turn related to ENSO and SAM. The statistical
221 significance of the result is assessed using a student t- test at 10 % significance level.

222

223 We also constructed composite maps considering the four years with largest and smallest
224 frequency of occurrence of long-lived RWPs using a threshold of 15 m/s. We plotted the upper-
225 level zonal winds and the meridional gradient of absolute vorticity in order to study the main
226 changes in the structure and intensity of the waveguide. To complement the analysis composite
227 maps were also constructed separately during El Niño and La Niña years as well as during
228 positive/negative SAM events using only three years per event in these cases due to the short
229 record available. Lastly we also performed one point lag linear regression maps between days -4
230 and +6 of the RWP amplitude to illustrate changes in the propagation under different conditions.

231

232 **3. Results**

233

234 **3.1 Mean and interannual variability of RWPs**

235 Before applying the algorithm, we calculated the climatological amplitude of the envelope in the
 236 SH. The highest values of the amplitude are located in the latitudinal band between 40-60°S and
 237 from 50-250 ° E (Figure 1). There is a minimum located around the southern tip of South
 238 America consistent with the results of Souders et al., (2014a) (in their case on annual time
 239 scale). The spatial structure is similar on individual months (not shown). As a result RWPs
 240 tracked using restrictive thresholds could be abruptly interrupted in this area and may cause the
 241 detection of high dissipation rates between 280-330°E.

242

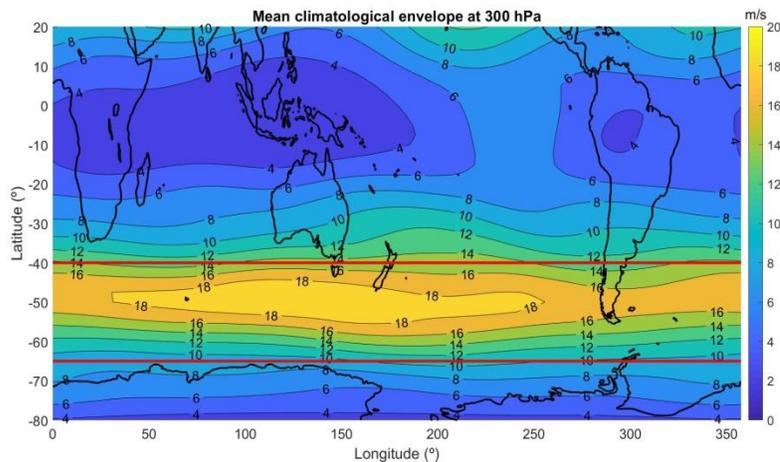


Figure 1. Mean climatological amplitude (m/s) during DJFM solid red lines shows the limitations of the area of study.

253

254 Figure 2 presents a summary of statistics of RWPs properties. It shows an exponential decrease
 255 in the duration of the wave packets such that about 70% do not reach 7 days and barely 2% of
 256 the trajectories have a lifespan of 14 days or above during the season. In addition, Figure 2
 257 shows that around 80% of the trajectories travel between 30°-170° in the longitudinal direction
 258 and only about 2.5% are able to go around the globe. The mean distance travelled by a packet is
 259 128 °, the median is 84° and the inter-quartile is 85°. In the case of the waves' lifespan they show
 260 a mean duration of 5 days, a median of 4 days and an inter-quartile range of 3 days. These
 261 results are close to those obtained by Sagarra and Barreiro (2020). Also Sagarra and Barreiro
 262 (2020) found that the mean speed of the packets is 20 m/s with a standard deviation of 6,6 m/s,
 263 while in our study we have obtained a mean speed of 21 m/s and a standar deviation of 4.3 m/s.
 264 The mean duration found here is lower than in Souders et al., (2014b), which may be, at least
 265 partially due to the use of a less restrictive threshold in that study (14 m/s).

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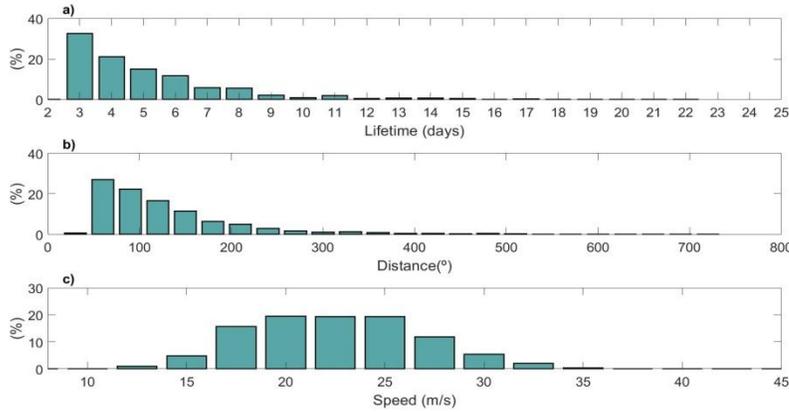
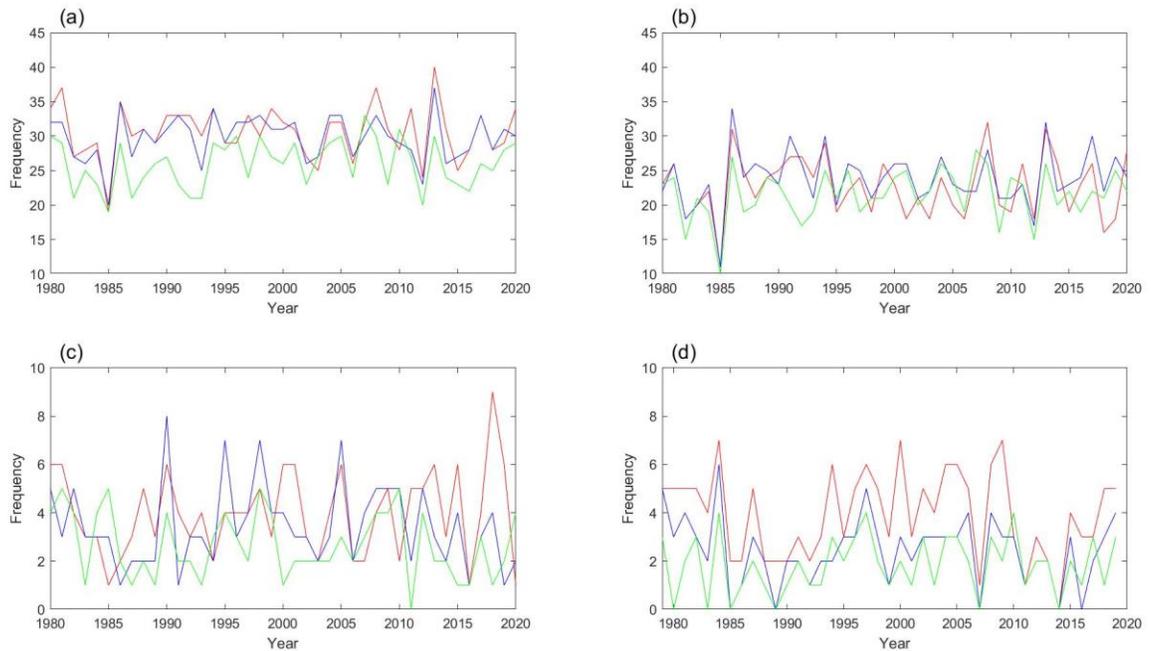


Figure 2. Climatological properties of RWP tracked between 1979-2020 during DJFM using a minimum threshold of 15m/s; (a) lifespan, of RWP (b) distance travelled, and (c) the mean speed per packet.

283 RWP shows large interannual variability, such that there are years with a minimum of 20 and
 284 years with a maximum of 40 events (Figure 3a). In addition the time series of RWP detected are
 285 highly correlated for different thresholds, but the number diminishes significantly for 17m/s.
 286 Looking at subsets of RWP, it is clear that the overall behaviour of the total number of events is
 287 dominated by the variability in occurrence of SL waves. For ML and LL waves the number of
 288 events per year is similar, relatively small and sensitive to the threshold used so that their number
 289 increases with a lower threshold. In contrast, they show large interannual variability percentage
 290 wise, suggesting that large scale circulation conditions set up by low-frequency climate modes
 291 can modulate their occurrence.

The frequency of occurrence of the total number of



292
 293
 294 Figure 3. Interannual variability of the total RWP tracked (a), short-lived waves (b), medium-lived waves (c), and long-lived
 295 waves (d). Red blue and green lines indicate the results obtained using thresholds of 13,15 and 17 m/s, respectively.

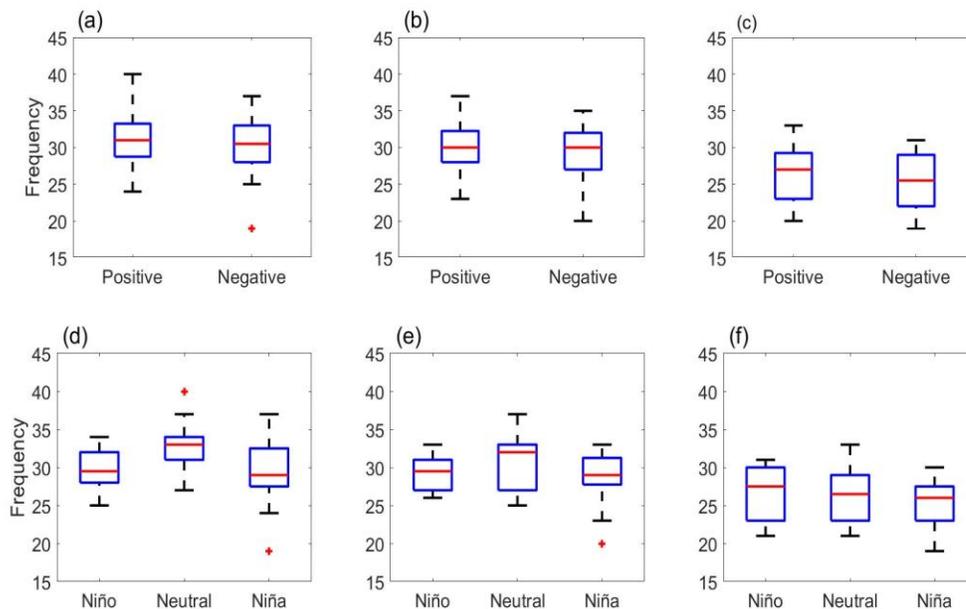
296 In the next section we study this focusing on the LL RWPs, given that a better understanding can
 297 provide maximum improvements on prediction on long-range to subseasonal scales. It is also
 298 worth mentioning that there is a negative correlation between the occurrence of SL and LL
 299 waves with $r=-0.47$ for a threshold of 15 m/s, so that the detection of several LL waves also
 300 mean less SL waves during that particular year. This correlation value is statistically significant
 301 according the t student test with a 90% of significance.

302

303 3.2 Impact of SAM and ENSO on long-lived packets

304

305 Figure 4 shows the frequency of occurrence of the total number of RWPs stratified according to
 306 the different phases of ENSO and SAM. In the case of SAM, both phases show a similar median
 307 for thresholds of 13 and 15 m/s, while a positive phase favours the occurrence of RWPs against
 308 negative SAM for a threshold of 17 m/s. This latter result is in agreement with Sagarra and
 309 Barreiro (2020). In the case of ENSO, El Niño and La Niña events present less RWPs than
 310 Neutral years for thresholds of 13 and 15 m/s while for a threshold of 17 m/s El Niño events
 311 tends to favour more RWPs than during La Niña. Similar results are found considering only SL
 312 waves, consistent with the fact that they represent the largest percentage of waves.



313

314 Figure 4. Boxplots of interannual variability of all RWPs detected during different SAM (a-c) and ENSO (d-f) phases for
 315 different minimum thresholds.

316

317 When we focus on long-lived RWPs, the algorithm detected a larger number of waves during El
 318 Niño compared to Neutral and La Niña years, particularly for a threshold of 13 m/s, although the
 319 result holds for other thresholds as well (Figure 5a). Additionally the dispersion of LL RWPs
 320 detected during El Niño is lower compared to Neutral and La Niña years. During La Niña years

321 the frequency of occurrence of RWPs shows much higher interannual variability than during El
 322 Niño events for the lowest threshold.

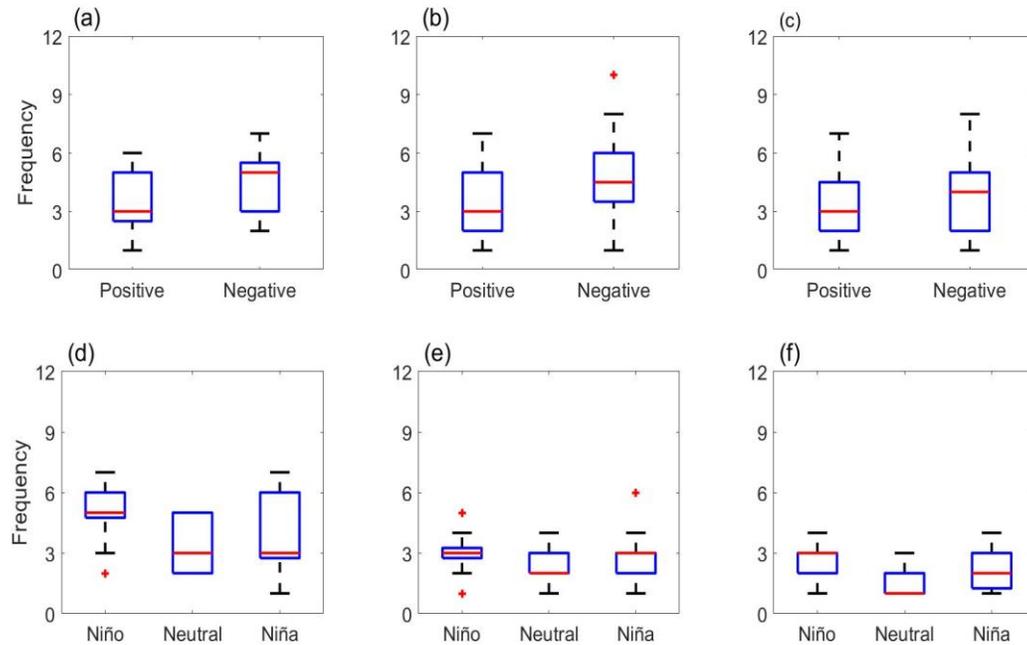
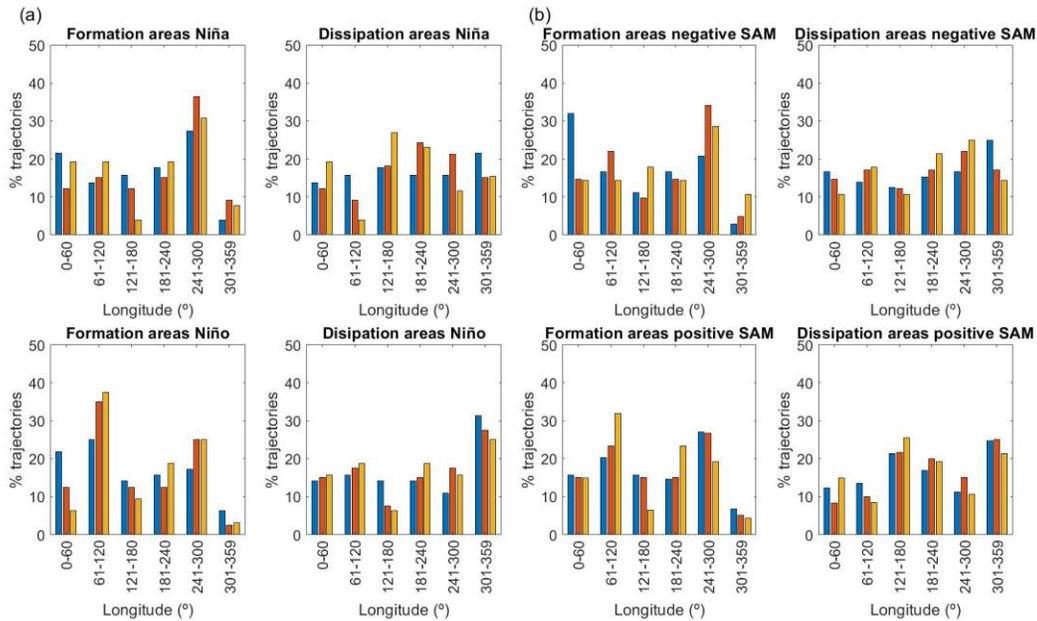


Figure 5. Analogous to figure 4, but for long-lived RWPs.

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 325

326 ENSO also affects the formation and dissipation areas of LL RWPs (Figure 6). During El Niño
 327 events, there is a main formation area between 61-120°E (South-Southwest of Australia) and a
 328 secondary area at 241-300°E (East of South America), whereas in the case of La Niña only the
 329 latter is a well defined formation region. On the other hand, during El Niño events there is a main
 330 area of dissipation in the South Atlantic (301-359°E) while in the case of La Niña no region
 331 stands out. These results are independent of the threshold considered. Note that in the case of El
 332 Niño there seems to be a compensation in detection of LL RWPs in the regions (0-60°E) and
 333 (61-120°E), depending on the threshold. For example, similar number of LL RWPs are detected
 334 in these two areas for 13 m/s, but for 17 m/s the number detected is much larger in (61-120°E)
 335 than (0-60°E). A plausible explanation is that most of the waves detected in the 0-60°E section
 336 using the lowest threshold are spotted in the following area when we applied more restrictive
 337 thresholds because the packet amplitude had more time to grow.

338



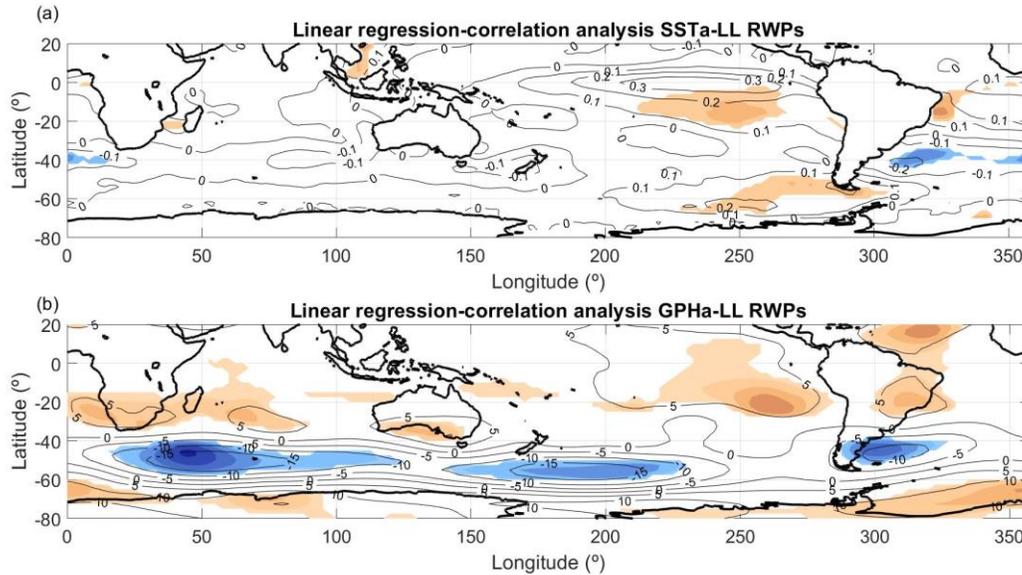
339
 340 Figure 6. Formation and dissipation areas found during el Niño and la Niña events (a), and during positive and negative SAM
 341 (b). Colorbars represent different minimum thresholds, being the blue 13 m/s, orange 15 m/s and yellow 17 m/s.

342

343 Regarding the impact of SAM (Figure 5b), we observe that independently of the threshold,
 344 negative SAM conditions favour the detection of LL RWPs. Both phases have large dispersion
 345 but the median is always higher in the negative phase, especially for the lowest threshold (13
 346 m/s) which statistics is more robust due to the higher number of cases. As for main areas of
 347 formation, during positive SAM, the main formation regions are located in the eastern Pacific
 348 (241-300°E) and over South-Southeastern Africa (61-120°E). During negative SAM phases only
 349 the region located in the eastern Pacific stands out for all thresholds (Figure 6). In the case of
 350 dissipation there are two main areas in positive SAM: over the South Atlantic basin (301-359°E)
 351 and to the Southwest of Australia (121-180°E), where about 50% of the LL RWPs are no longer
 352 detected. On the other hand, in the case of negative SAM there is no main area of dissipation
 353 consistent for all thresholds.

354

355 According to above results, LLRWPs are influenced by both ENSO and SAM, in such a way that
 356 there is an increase in the frequency of the occurrence during El Niño and negative SAM phases.
 357 This influence can be seen in the regression maps of sea surface temperature and 300hPa
 358 geopotential height anomalies against the (standardized) time series of occurrence of LL RWPs
 359 (Figure 7). Results are shown for a threshold of 13 m/s because the number of LL RWPs is larger
 360 and the statistics are more robust, (nonetheless the patterns for a threshold of 15 m/s are very
 361 similar but with weaker anomalies). As expected from the boxplots an increase in the number of
 362 detected LL RWPs is correlated with positive tropical Pacific sea surface temperature anomalies
 363 (Figure 7a). Moreover it is correlated positively with the southeaster Pacific and negatively with
 364 the southwestern Atlantic.



365
 366 Figure 7. Linear regression maps of SST anomaly (above) and 300hPa GPH anomaly (down) with respect to the time series of LL
 367 RWPs for a threshold of 13 m/s. The colored areas indicate where the correlation is significant at the 10% level and black lines
 368 the value of linear regression.

369
 370 Consistent with these results, the tropical upper level geopotential height shows positive
 371 anomalies during years of increased LL RWPs (Figure 7b). In addition, there is a decrease in
 372 Z300 in midlatitudes and an increase in high latitudes, signalling the occurrence of the negative
 373 phase of SAM. Thus this map is in agreement with the results of Figure 5 and suggests that the
 374 jet and by extension, the waveguide, are enhanced and shifted northward during years of
 375 increased LL RWPs. In the next section we describe in more detail the changes in the jet.

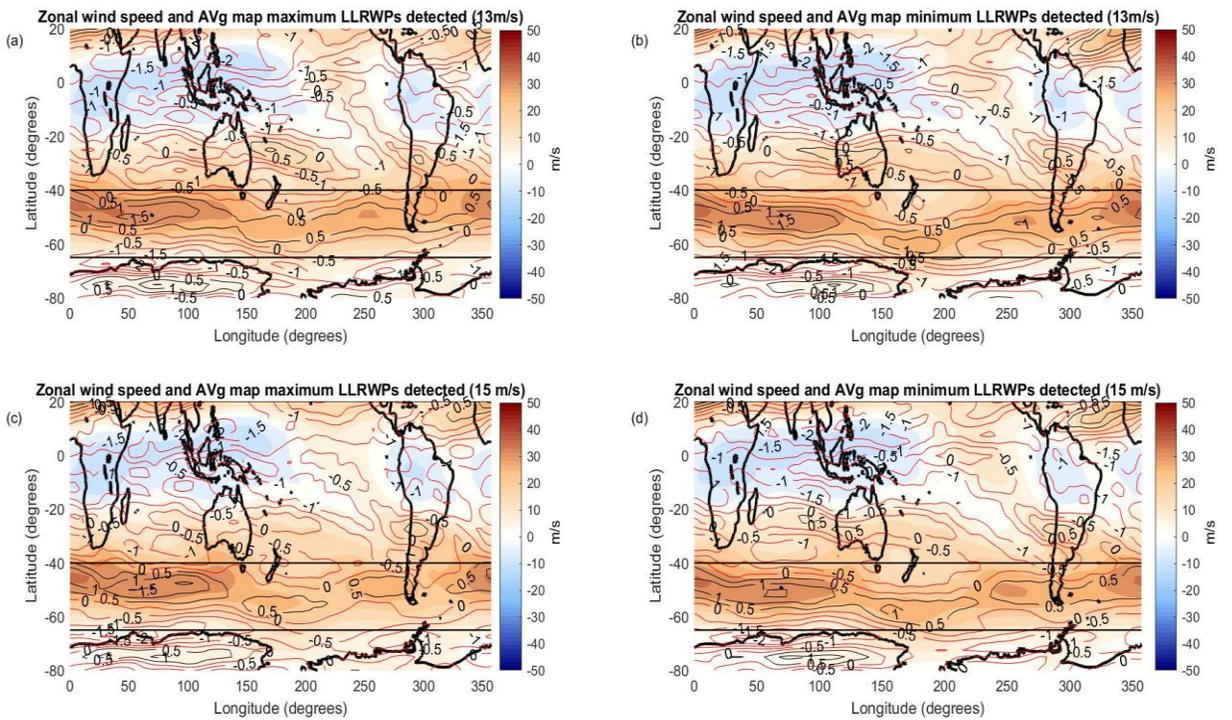
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377 3.3 Modification of the waveguide

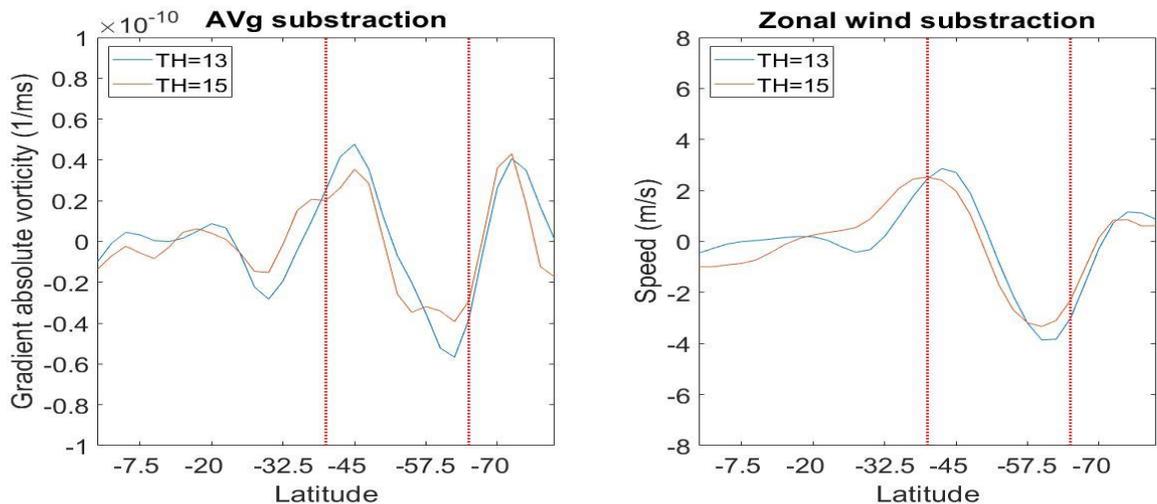
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379 Figure 8 shows a composite map of the upper level zonal winds and meridional gradient of
 380 absolute vorticity (GAV) constructed for the four years with the most and least number of LL
 381 RWPs detected for low (13 m/s) and medium (15 m/s) thresholds. We observe that the jet and the
 382 waveguide are more zonally symmetric during years of maximum number of LL RWPs.
 383 Moreover during the years with the smallest number of LL RWPs detected using a low threshold
 384 (13 m/s, Figure 8b) the jet and the waveguide are displaced southward, particularly between 130-
 385 220°E, compared with years with the highest occurrence of LL RWPs, (Figure 8a). Similar
 386 results are found using a medium threshold (15 m/s, Figures 8c, 8d). This is further illustrated in
 387 Figure 9 which show the difference of the zonal mean winds and GAV between years of
 388 maximum and minimum detection of LL RWPs for both thresholds.

389

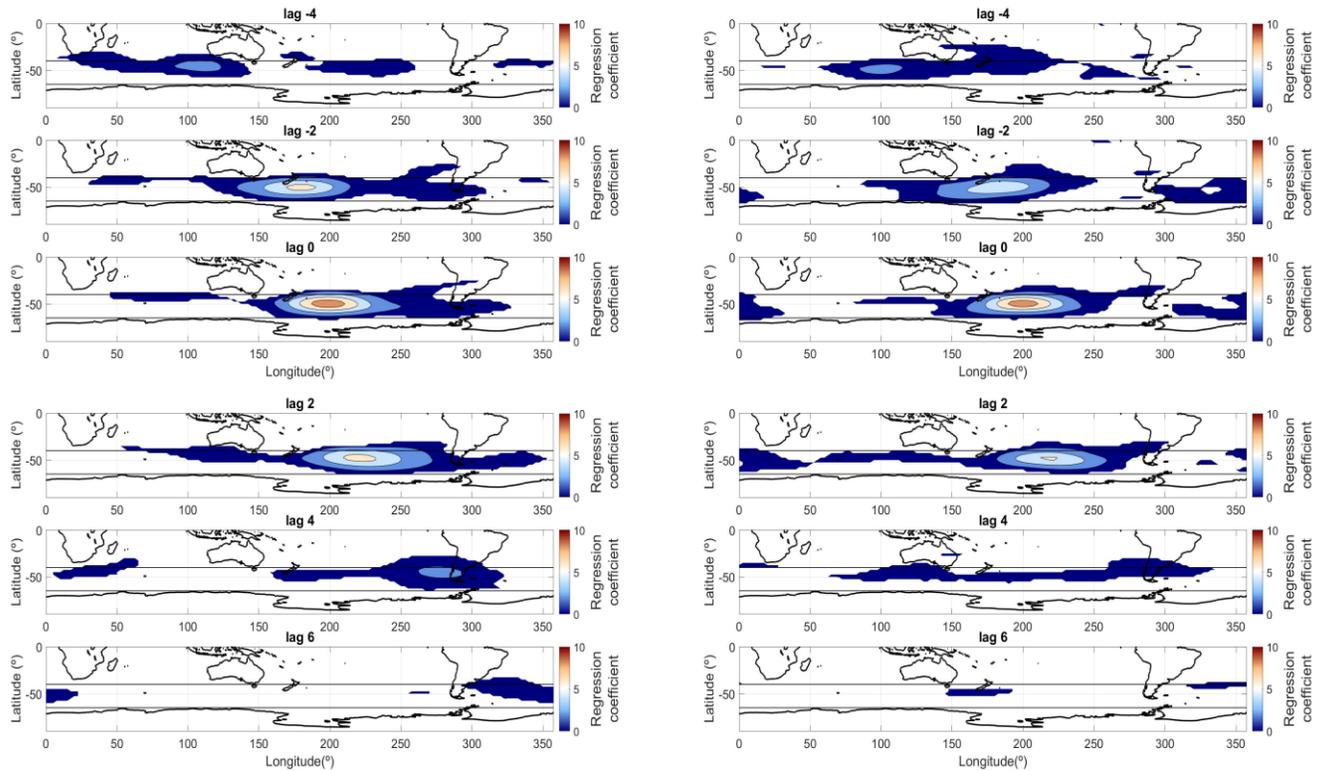


390
 391 Figure 8. Mean zonal wind and meridional absolute vorticity gradient (AVG) during the 4 years with maximum (a, c) and
 392 minimum (b, d) number of LL RWPs detected for thresholds of 13m/s (upper row) and 15 m/s (down row). Zonal wind speed
 393 (m/s) is indicated in the colorbar while GAV is represented with black (positive values) and red curve lines (negative values).
 394 The units of the AV gradient (1/ms) are multiplied by a factor of 10^{10} to enable an easier visual representation. Red straight lines
 395 limit the area of study.
 396
 397



398
 399 Figure 9. Difference of the zonal mean AV gradient (left) and zonal wind speed (right) between years with maximum and
 400 minimum number of LL RWPs for thresholds 13m/s (blue) and 15 m/s (orange). Red lines indicate the limits of the area of study
 401 (40-65°S).
 402
 403

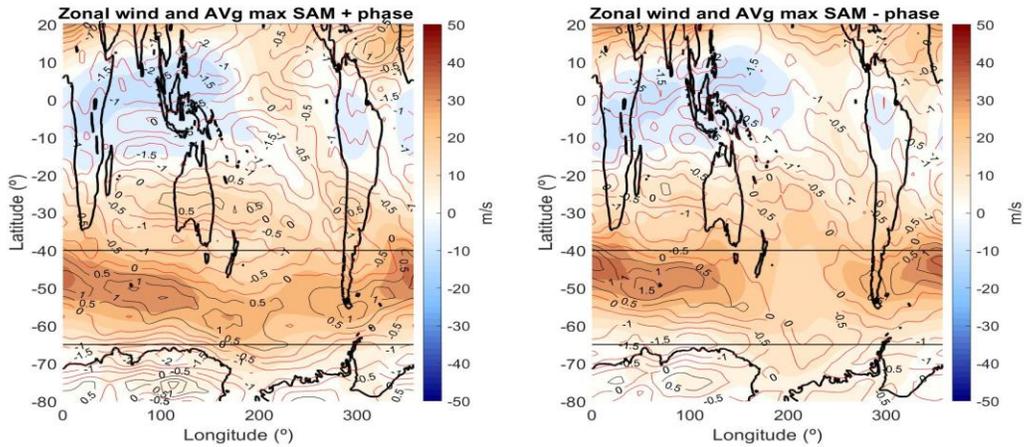
404 To illustrate the differences in the trajectory of wave packets we perform a linear regression of
 405 the wave packet amplitude with respect to that of a point located in 200°E,50°S from lag -4 to lag
 406 +6 (Figure 10). This location is chosen because it is in the middle of the Pacific sector where the
 407 largest differences in the zonal winds between extreme cases of RWPs was found. The maps
 408 show that the RWPs show largest amplitudes and are more coherent during the 11 days for the
 409 years with the largest number of LL RWPs detected, particularly on lags -2, +2 and +4. At lag +6
 410 it is still possible to capture the propagation of the packet in the case of maximum number of LL
 411 packets, whereas for years with the least packets the regression does not show a consistent area.
 412 This result further supports the hypothesis that the northward deviation of the mid-latitude jet
 413 favours the propagation of LL RWPs, as shown above.
 414



415
 416 Figure 10. Lag linear regression maps of the envelope amplitude with respect to that of a point located in the middle of the
 417 Pacific basin in 50°S 200°E during the years of maximum (left) and minimum (right) frequency of LL RWPs.
 418

419 We now assess changes in the jet associated with the phases of SAM in more detail. To do so we
 420 compare the most extreme three years of positive and negative SAM phases (Figure 11). As
 421 expected, during the positive SAM phase the jet stream is located further south compared to the
 422 negative SAM phase. Moreover, during the negative SAM the jet is zonally symmetric between
 423 the south Atlantic and the south of Australia (250°E and 150°E) and weak in the Pacific sector,
 424 interrupting the waveguide. On the contrary, during SAM positive phase the jet is overall
 425 stronger, but it has a more diagonal structure particularly in the Indian-Pacific sector that does
 426 not favour the zonal propagation of wave packets. Consistently this shape is similar to the one
 427 observed during the years of minimum number of LL RWPs (see e.g. Figure 8b).

428



429

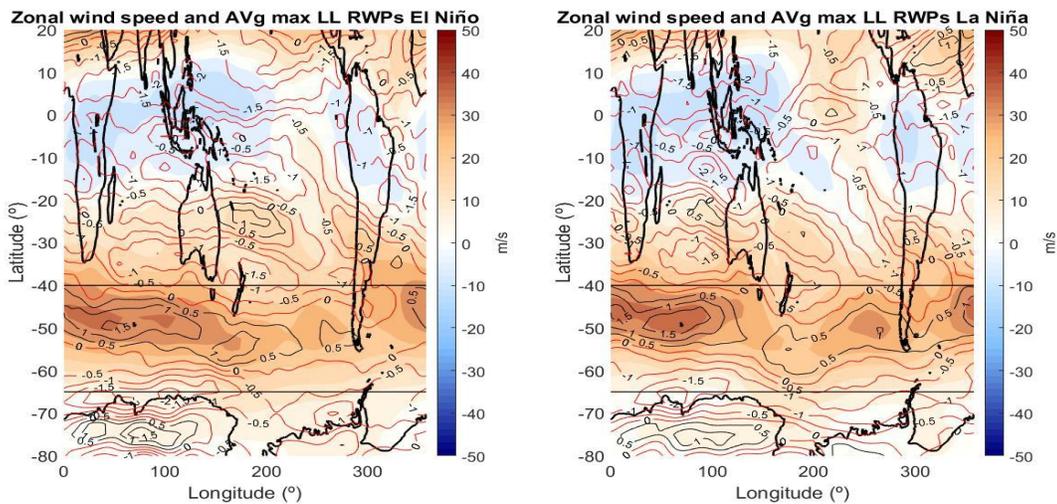
430

431 Figure 11. Mean zonal wind and meridional absolute vorticity gradient (GAV) during the 3 years with maximum positive (left)
 432 and negative (right) SAM phases. Zonal wind speed (m/s) is indicated in the colorbar, while the GAV is represented with black
 433 (positive value) and red lines (negative values). The units of the AV gradient are multiplied by a factor of 10^{10} /ms to enable an
 434 easier visual representation.

435

436 The importance of a zonally symmetric shape of the waveguide in favouring the detection of
 437 RWPs can be further shown comparing the flow during years with maximum number of LL
 438 RWPs detected during El Niño and La Niña events (for a threshold of 13 m/s) in Figure 12. In
 439 both cases we detected the same number of RWPs (6-7 per year) and the jet is both zonal and
 440 narrow. Nonetheless, it can be seen that during La Niña the jet begins to the Southwest of Chile
 441 and ends to the south of Australia, whereas during El Niño years it begins in the South Atlantic
 442 basin ending to the south of New Zealand. Thus, a high number of LL RWPs develop in a zonal
 443 and narrow flow independently on the longitudinal sector it occupies.

444



445

446

447 Figure 12. Analogous to Figure 11, but for the years with maximum frequency of LL RWPs during El Niño (left) and La Niña
 448 (right) events using the lowest threshold (13 m/s).

449

450 **4.Summary**

451
452 A detection algorithm was implemented to study the interannual variability of RWPs, and in
453 particular the impact of SAM and ENSO on those that have a lifespan longer than 8 days.

454 We have shown that ENSO and SAM control the frequency of occurrence and
455 formation/dissipation areas of LL RWPs, such that El Niño and a negative phase of SAM favour
456 higher numbers. Both conditions favour a more northward position of the jet stream which is also
457 narrow and more zonally symmetric. Moreover results show that La Niña can also provide
458 background conditions for the occurrence of large number of LL RWPs, but the dispersion is
459 large, suggesting that only a few particular events favour the generation of these long-lived
460 packets.

461
462 The formation and dissipation areas of LL RWPs differ depending on the phase of ENSO and in
463 some cases they depend on the threshold considered. In the case of formation areas, during La
464 Niña events it is shown that most of the LL RWPs (near 30%) are detected at the eastern Pacific
465 basin (241-300°E), which is consistent with the beginning of the jet to the southwest of Chile
466 seen in Figure 12. On the other hand, no main dissipation area was detected. In the case of El
467 Niño, the main formation area is located at 61-120°E, and the main dissipation area is located at
468 300-359°E.

469
470 During SAM both phases have the same main formation area as in La Niña events (eastern
471 Pacific 241-300°E) and positive SAM phases showed another maximum at 61-120°E (as in El
472 Niño). Thus there are two main formation areas of LL RWPs in the SH, whose importance vary
473 interannually depending of the dominant mode of variability. As in Niño events, the main
474 dissipation area for positive SAM occurs in the south Atlantic possibly associated with the
475 minimum in envelope amplitude seen in Figure 1.

476
477 Results found in this study thus suggest that LL RWPs are more common during years with El
478 Niño and negative SAM phases. Given their link with extreme weather events, our results are
479 indicative that extended range forecasting of extreme events may be more feasible during years
480 when the global modes of climate variability are present in these phases, and will be less skilful
481 when La Niña and positive SAM phases are present.

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483
484
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489 de Catalunya.

490 **Data Availability Statement**

491 NCEP/DOE Reanalysis 2 and ERRST v5 can be downloaded at the websites:
 492 <https://psl.noaa.gov/data/gridded/data.ncep.reanalysis2.html> and
 493 <https://psl.noaa.gov/data/gridded/data.noaa.ersst.v5.html>. ENSO and SAM indexes are available
 494 at <https://origin.cpc.ncep.noaa.gov/>

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