

Upper and lower plane bed revisited

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Key Points:

- Compiled experimental and observational data suggest a new definition of plane beds
- Previous definitions of upper and lower plane beds do not explain the dataset well
- The lower and upper plane beds are redefined as flat bedforms that appear when suspended load is inactive and active, respectively

Abstract

This study proposes a new definition of plane bed regimes. The upper plane bed is defined as a flat bedform that appears under conditions of active suspended load, whereas the lower plane bed forms when sediment particles are moved only by bed-load. Previous studies have recognized that two types of plane bed could be distinguished by shear stress, particle size, or flow regime. However, the compilation of a large amount of existing open-channel flow data indicates that the plane bed regimes cannot be differentiated by existing definitions. Newly produced phase diagrams indicate that plane bed data plot in two separate regions, and the gap between the two regions corresponds to the threshold condition of suspension. Further investigation is needed to understand the physical mechanisms generating lower plane beds and characterizations of parallel laminae from lower and upper plane beds.

Plain Language Summary

Bedforms are topographic features that appear at fluid–solid interfaces, such as river beds. Several types of bedform have been found, and “plane bed” is a type of bedform commonly observed in rivers. Previous studies have recognized that there are two types of plane beds and they could be distinguished by flow properties or bed particle size. However, this study finds that previous definitions of plane beds are inconsistent with existing river observation and experimental data. Therefore, we propose a new definition for plane beds. This study establishes that plane beds can be formed with a wider range of flow properties and bed particle sizes than previously considered.

1 Introduction

Bedforms and sediment transport have attracted great interest of sedimentologists and engineers, as well as geomorphologists. Although the formation conditions of bedforms have been examined by theoretical methods, such as linear stability analysis (e.g., Colombini & Stocchino, 2008), predicting the flow conditions of bedform phases remains challenging owing to the complexity of relationships between flows and bedforms. Therefore, numerous laboratory experiments (e.g., Gilbert, 1914; Guy et al., 1966; Taylor, 1971) and field observations (e.g., Colby & Hembree, 1955; Culbertson et al., 1972) have been conducted to obtain empirical relationships between flows and bedforms. The compilation of large data sets, including experiments and field observations, facilitates the understanding of the formation conditions of bedforms in response to flow behavior in alluvial rivers using bedform phase diagrams. Bedforms are vital for river management and reconstruction of paleo-flow conditions (Best, 2005). Thus, bedform phase diagrams have been proposed based on laboratory and field observations (e.g., van den Berg & van Gelder, 1993). Recently, Ohata et al. (2017) proposed bedform phase diagrams in three-dimensional parametric space, and their diagrams successfully reconstructed paleo-flow conditions from sedimentary structures.

The plane bed is a representative type of bedform that is commonly observed on alluvial river beds (e.g., Hauer et al., 2019), beach surfaces (Vaucher et al., 2018), and can be recorded as parallel lamination in sandstone beds formed in various environments such as rivers (e.g., Mazumder & Van Kranendonk, 2013; Umazano et al., 2012) and submarine fans (e.g.,

62 Eggenhuisen et al., 2011; Jobe et al., 2012). In submarine environment, it is known that parallel
63 lamination commonly occurs in turbidites as T_b or T_d divisions (Bouma, 1962).

64 The plane bed phase is characterized by nearly flat topography and was initially referred
65 to as “the smooth phase of bedforms” (Gilbert, 1914; Owens, 1908). Later, Simons et al. (1961)
66 used the term “plane” to avoid confusion with the hydraulically smooth phase of flows. Owens
67 (1908) and Strahan et al. (1908) observed that the bedform phase changes from ripples to
68 antidunes with increasing flow discharge, and the smooth phase of bedform (i.e., plane bed)
69 develops in the intermediate condition between ripples and antidunes. In addition, Simons et al.
70 (1961) recognized that the plane bed phase also appears under the condition when bed shear
71 stress is lower than that forming ripples or dunes. Thus, two conditions of plane bed formation
72 have been recognized.

73 However, researchers have had differing opinions on the definition and formation
74 conditions of the two types of plane bed. The two types of plane bed have been classified
75 depending on (1) criticality of sediment motion (Simons et al., 1961; Simons & Richardson,
76 1962), (2) criticality of Froude number (Dey, 2014; Venditti, 2013), and (3) size of bed material
77 and flow intensity (Allen, 1968; Southard, 1971). Simons et al. (1961) recognized two conditions
78 of plane bed in their flume experiments—plane beds with and without sediment motion. Simons
79 et al. (1961) and Simons & Richardson (1962) also classified alluvial flows based on bedforms
80 into three regimes—a lower flow regime (plane beds without sediment movement, ripples, dunes
81 with ripples superposed, dunes), a transition zone (washed-out dunes), and an upper flow regime
82 (plane beds with sediment movement, standing waves, antidunes, chutes-and-pools). Although
83 Simons et al. (1961) noted that flows change from the lower to upper flow regimes at a Froude
84 number less than unity, Dey (2014) and Venditti (2013) stated that lower and upper flow regimes

are consistent with Froude-subcritical (Froude number < 1) and Froude-supercritical flow regimes (Froude number > 1), respectively. In other words, two types of plane bed are differentiated at Froude number = 1. On the other hand, Liu (1957) and Bogárdi (1961) reported a region of flow conditions in which sediment motion occurred, but ripples or other bed relief did not appear. Furthermore, the phase diagrams proposed by Allen (1968) and Southard (1971) suggested that there is a stable plane bed region where coarse sediment (median diameter > 0.6 - 0.7 mm) is transported at low flow velocities. This region is termed lower plane beds. Therefore, Allen (1968) and Southard (1971) recognized two types of plane bed in addition to the plane bed without sediment motion, which are lower plane beds on coarse-grained beds and plane beds in the upper flow regime.

In this study, we propose new definitions of plane beds that explain existing experimental and observational datasets well, compared to previous definitions. We compiled data pertaining to plane beds in unidirectional open-channel flows, and plotted those using dimensionless parameters as axes without any *a priori* assumption of types of plane bed. The compilation of a large amount of existing data from the literature indicates that two classes of plane bed regimes are recognized and the boundary between the two plane bed regimes matches the threshold condition of suspension.

2 Data Sources and Dimensionless Parameters for Morphodynamic Conditions

We compiled a total of 937 sets of data pertaining to plane beds, from the literature, to investigate the plane bed regimes. The dataset consisted of 892 sets of laboratory data and 45 sets of field data. Relevant references and the basic parameters are summarized in Table S1. A

wide range of hydraulic and sediment parameters are included in the datasets. The median diameter D_{50} ranges from 11 μm to 44.3 mm, the flow depth h ranges from 1.2 mm to 2.74 m, and the flow velocity U ranges 0.0448 to 2.38 m/s. Further, the data is classified into three types based on the suspended sediment concentration C_s , as follows: (a) the suspended sediment concentration was measured (hereafter referred to as data $C_s > 0$), (b) the suspended sediment concentration was recorded as zero or the bed material did not move (hereafter referred to as data $C_s = 0$), and (c) the suspended sediment concentration was not available (hereafter referred to as data no C_s).

First, the hydraulic conditions of plane bed regimes were examined with the sediment diameter and the sediment transport mechanism by which the sediment was moved, i.e., no sediment movement, bed-load, and suspended load. Bedform phases were expressed in the space of dimensionless parameters that reflected the properties of flows and sediment particles (Ohata et al., 2017; van den Berg & van Gelder, 1993). In this study, we employed the following dimensionless parameters to represent hydraulic conditions and sediment properties: the particle Reynolds number Re_p , Shields number τ_* , and the suspension index u_*/w_s .

The particle Reynolds number Re_p is defined as:

$$\text{Re}_p = \frac{\sqrt{RgD_{50}}D_{50}}{\nu} \#(1)$$

where R represents the submerged specific gravity of the sediment, g denotes the gravitational acceleration, and ν denotes the kinematic viscosity of the fluid. The submerged density is defined as $R = (\rho_s - \rho_f)/\rho_f$, where ρ_s and ρ_f represents the sediment and fluid densities, respectively. The kinematic viscosity ν was assumed to be a function of temperature according to the relationship for clear, fresh water (van den Berg & van Gelder, 1993):

$$\nu = [1.14 - 0.031(T - 15) + 0.00068(T - 15)^2]10^{-6} \#(2)$$

128 where T represents the water temperature in degrees Celsius. A value of 20 °C was assumed for
129 data where T was not reported. Shields number τ_* is defined as:

$$\tau_* = \frac{u_*^2}{RgD_{50}} \#(3)$$

130 Here, the shear velocity u_* for the field data was computed as $u_* = \sqrt{ghS}$, where S represents
131 the slope. For laboratory data, the side-wall effect was removed using the method proposed by
132 Chiew & Parker (1994) (see Ohata et al. (2017) for details), and we calculated the bed
133 component of the shear velocity. The suspension index is expressed as the ratio of the shear
134 velocity u_* to the settling velocity of sediment w_s . The settling velocity w_s was estimated using
135 the relationship formulated by Ferguson & Church (2004):

$$w_s = \frac{RgD_{50}^2}{C_1\nu + \sqrt{0.75C_2RgD_{50}^3}} \#(4)$$

136 where constants C_1 and C_2 were set to 18 and 1, respectively.

137 The sediment transport regimes are classified based on the transport mechanism—no
138 sediment movement, bed-load dominated, mixed-load, and suspended-load dominated regimes
139 (Church, 2006). The boundary between no sediment movement and other regimes is defined by
140 Shields curve, which is the critical condition for the initiation of particle motion (e.g., Shields,
141 1936). Based on Shields' experimental data, Brownlie (1981a) proposed a function describing
142 the threshold condition of particle motion as:

$$\tau_c^* = 0.22\text{Re}_p^{-0.6} + 0.06 \exp(-17.77\text{Re}_p^{-0.6}) \#(5)$$

Next, the critical condition for the initiation of suspension is expressed using the suspension index u_*/w_s (Bagnold, 1966; van Rijn, 1984; Niño et al., 2003). The critical condition of suspension represents the boundary between the bed-load dominated regime and a regime where the bed materials are transported with suspension. The threshold condition for suspension, obtained by Niño et al. (2003), is:

$$\frac{u_*}{w_s} = \begin{cases} 21.2\text{Re}_p^{-1.2}, & 1 < \text{Re}_p < 27.3 \\ 0.4, & 27.3 \leq \text{Re}_p \end{cases} \quad \#(6)$$

Therefore, the region below equation (5) denotes the no sediment movement regime, the region between equations (5) and (6) denotes a bed-load dominated regime, and the region above equation (6) denotes mixed-load or suspended-load dominated regimes where the bed materials are moved with suspension.

Second, the formation conditions of the plane bed regimes were investigated using Froude number. The phase diagrams were obtained using Re_p , u_*/w_s , and Froude number as the axes. Froude number is the ratio of the inertial force to the gravitational force, defined as:

$$\text{Fr} = \frac{U}{\sqrt{gh}} \quad \#(7)$$

Here, the inertial force is represented by the flow velocity U and the gravitational force \sqrt{gh} denotes the velocities of the free-surface waves. Flows with $\text{Fr} < 1$ and $\text{Fr} > 1$ are called Froude-subcritical and Froude-supercritical flows, respectively.

3 Results

3.1 Sediment transport mode and plane bed formation

First, the plane bed data were plotted in $Re_p - \tau_*$ space (Figure 1a) to investigate the relationships among the plane bed regimes, the particle diameter, and the criteria for sediment movement. The threshold condition of particle motion (equation (5)) and suspension (equation (6)) are described in Figure 1a. Threshold condition of suspension (equation (6)) was rearranged using equations (1), (3), and (4) where $R = 1.65$ and $\nu = 1.0 \times 10^{-6}$. The median diameter D_{50} is also recasted from Re_p with $R = 1.65$ and $\nu = 1.0 \times 10^{-6}$, which is represented on the top axis of Figure 1a.

Shields number of the plane bed ranges from 0.05 to 10, and the plane bed data are divided into two separate regions at $\tau_* = 0.1$ – 0.2 (Figure 1a). The plane bed data with low τ_* plot around the threshold condition of particle motion, and the region with high τ_* plots above the threshold condition of suspension. The threshold condition of particle motion does not divide the plane bed regime; rather, the threshold condition of suspension runs between the two regions of data. The particle Reynolds number of the data with low and high τ_* ranges as follows: $9 < Re_p < 4.0 \times 10^4$ ($0.2 \text{ mm} < D_{50} < 40 \text{ mm}$) and $0.1 < Re_p < 300$ ($0.01 \text{ mm} < D_{50} < 2 \text{ mm}$), respectively. The data $C_s > 0$ and field observation data are included in the region of high τ_* . Most data $C_s = 0$ plot around the criteria for particle motion, and three points of data $C_s = 0$ plot around $\tau_* = 1$.

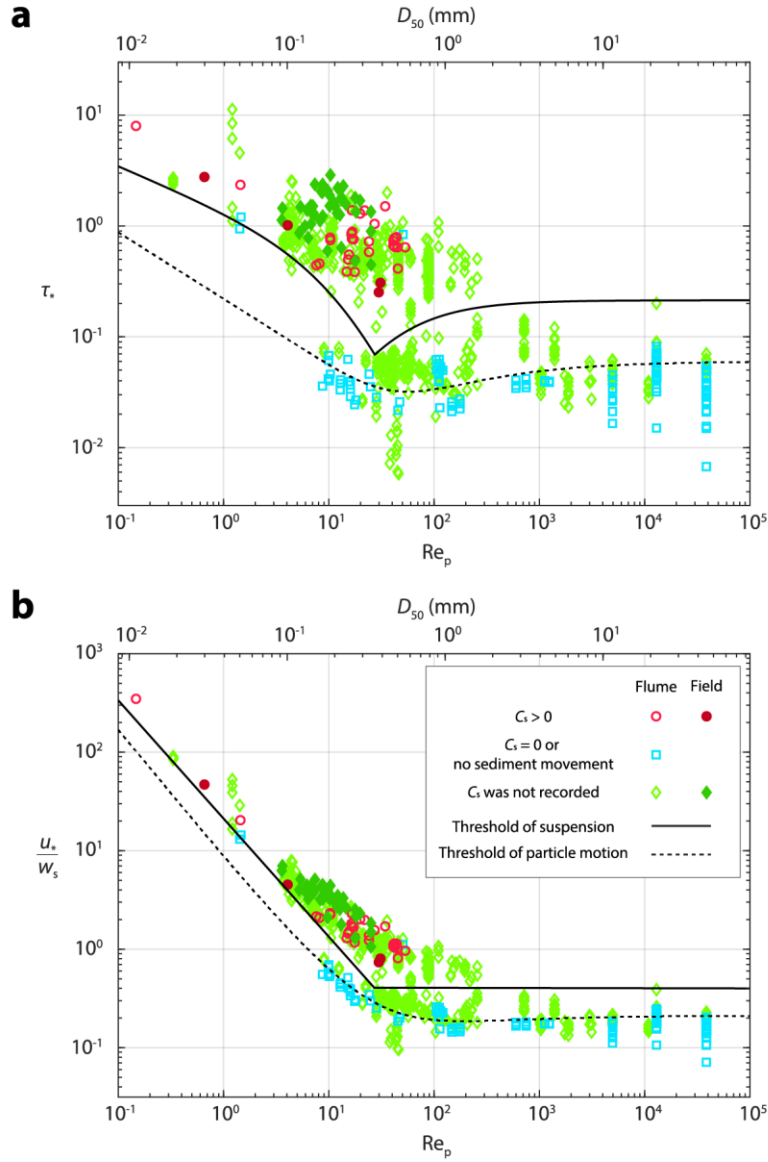


Figure 1. Plane bed regime in (a) Re_p – τ_* space and (b) Re_p – u_*/w_s space. The dotted line denotes the threshold condition for particle motion (equation (5)). The solid line denotes the threshold condition for suspension (equation (6)). Both lines were extended to $Re_p = 10^{-1}$ and 10^5 .

Second, the relationships among the plane bed regimes, the sediment diameter, and the criteria for suspension were examined by plotting u_*/w_s versus Re_p (Figure 1b). The threshold condition of suspension (equation (6)), and for comparison, the threshold condition of particle motion (equation (5)) are described in Figure 1b. To plot equation (5) in Re_p-u_*/w_s space, R and ν were assigned the values 1.65 and 1.0×10^{-6} , respectively. The median diameter D_{50} is depicted on the top side of Figure 1b.

Figure 1b shows that the plane bed data plot in two regions separated by the threshold condition of suspension as well as in $Re_p-\tau_*$ space (Figure 1a). The threshold condition of particle motion runs through the middle of the data with a low suspension index. With respect to the particle Reynolds number, the data with high suspension index plot in the region $0.1 < Re_p < 3200$ ($0.01 \text{ mm} < D_{50} < 2 \text{ mm}$) and the data with low suspension index plot in the region $9 < Re_p < 4.0 \times 10^4$ ($0.2 \text{ mm} < D_{50} < 40 \text{ mm}$). The data $C_s > 0$ and field observation data plot above the threshold condition of suspension. Most data $C_s = 0$ plot along the criteria for particle motion, although a few points of data $C_s = 0$ plot above the suspension criteria.

3.2 Impact of Froude number

The consistency of the flow regime concept based on Froude number (Dey, 2014; Venditti 2013) to plane bed regimes was investigated using the diagrams described in Re_p-Fr (Figure 2a) and u_*/w_s-Fr spaces (Figure 2b).

Figure 2 shows that all the data covers a wide range of Froude number ranging from 0.1 to 3.5, and the separated regions cannot be seen on the Froude number axis. The field observations data have a lower Froude number ($Fr < 0.5$), whereas the field data exceed 0.7 in

the suspension index and plot above the threshold condition of suspension (Figure 1b). For $C_s > 0$, the data fall in the domain $0.1 < Fr < 1.3$, and the data $C_s = 0$ distribute in the domain $0.2 < Fr < 1.0$.

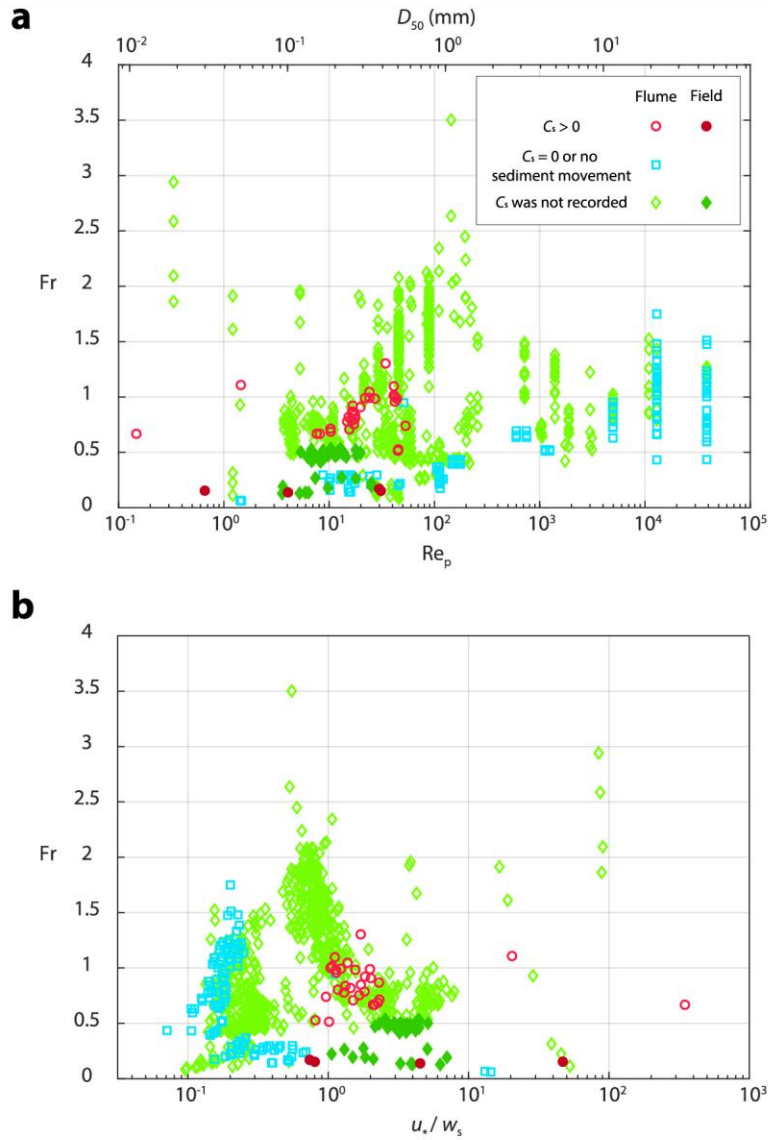


Figure 2. Plane bed regime in (a) Re_p – Fr space and (b) u_*/w_s – Fr space.

4 Discussion

Data from the existing field observations and flume experiments for plane beds were collected to revisit the formation conditions of plane beds. The phase diagrams indicate that the phase regions of the plane beds are not separated by particle diameter or the criticalities of sediment motion and Froude number, which were used to classify the lower and upper plane beds. In contrast, observational data of the plane beds plot in two separate regions in the $Re_p-\tau_*$ and Re_p-u_*/w_s diagrams, and the boundary between the two regions matches the threshold condition for suspension. Therefore, our phase diagrams suggest that the upper and lower plane bed phases are redefined—the upper plane beds are plane beds with suspension, and the lower plane beds are plane beds without suspension.

The new definition of upper plane beds is consistent with previous studies that describe the phase diagrams of bedforms based on theoretical analysis (Nakasato & Izumi, 2008). Nakasato & Izumi (2008) performed linear stability analysis considering active suspended loads and demonstrated that dune formation is suppressed by suspended loads. Further, the contribution of suspension to the formation of plane beds from dunes has been observed in flume and numerical experiments (Naqshband et al., 2015; Naqshband et al., 2017; Bradley & Venditti, 2019). The bed-load transport plays a vital role in the formation of stable dunes (Naqshband et al., 2014; Bradley & Venditti, 2019), whereas dune height decreases as more bed particles are transported as suspended loads (Naqshband et al., 2017; Bradley & Venditti, 2019).

According to Figure 1a, our compiled datasets indicate that the previous definitions of lower plane beds by Simons et al. (1961) and Southard (1971) do not demonstrate the two separated regimes of plane beds. In previous studies, lower plane beds were defined as plane beds without sediment motion (Simons et al., 1961) or plane beds just above the threshold of

sediment motion under subcritical flow with coarse grains (Southard, 1971). In Figure 1a, two separated regimes that have low Shields number ($\tau_* < 0.1$ – 0.2) and high Shields number ($\tau_* > 0.1$ – 0.2) are observed. For the data with low Shields number ($\tau_* < 0.1$ – 0.2), plane beds can appear with finer sediments ($D_{50} < 0.6$ – 0.7 mm) and the threshold condition of particle motion plot through the middle of the data scatter. This suggests that the traditional definitions of plane beds by Simons et al. (1961) and Southard (1971) are inconsistent with the dataset.

It is noteworthy that Froude number has a small impact on the formation of the plane beds, and there are no gaps in Froude number as shown in Figure 2. Froude number for the field data can be lower than unity, due to the great depth of natural rivers, and when the shear stress for the field data is high. Indeed, Ma et al. (2017) conducted field surveys at the Yellow River, China, measuring sediment loads and bed topography, and observed a nearly flat bed with low Froude number ($Fr = 0.23$). The ratio of length to height of the bed waves ranged from 500 to 2000, which is a much higher ratio than that of dunes in other natural rivers (Ma et al., 2017). They estimated the hydraulic resistance in the Yellow River based on Engelund–Hansen sediment transport theory (Engelund & Hansen, 1967) and found that the bed resistance was similar to that of plane bed. Shields number in the Yellow River was 1.05 with $D_{50} = 0.09$ mm, which would plot above the threshold condition of suspension (see Figure 1a), implying that the bedform observed in the Yellow River corresponds to the upper plane bed defined in this study. To summarize, the lower and upper plane beds cannot be defined by Froude number, as stated by Dey (2014) and Venditti (2013).

Similarly, our phase diagrams imply that the origin of parallel laminae in sandstones cannot be presumed as the upper plane bed. Previous studies suggest that the lower and upper plane beds can be distinguished by grain size, therefore, fine-grained parallel lamination has

256 been considered as the upper plane bed deposits. Best & Bridge (1992), whose experimental
257 results were compared with parallel laminated sandstones, assumed that the laminae were
258 produced by the upper plane bed. However, lower plane beds can also form with fine-grained
259 materials, as depicted in Figure 1. Subsequently, it is difficult to distinguish from the outcrops
260 whether the parallel laminae were formed with or without suspension. This study suggests that
261 much remains to be done theoretically and experimentally to interpret the flow conditions that
262 form parallel lamination. To this end, detailed characterizations of parallel laminae from lower
263 and upper plane beds are required under controlled conditions in flume experiments.

264 The origin of the lower plane bed phase in the new definition is puzzling. In previous
265 studies, the controlling factor in the lower and upper plane bed regimes was assumed to be the
266 grain size. Leeder (1980) interpreted that coarse-grained plane beds are formed because flow
267 separation is prevented by the bed roughness. Recently, Blois et al. (2014) proposed that bed
268 permeability may be another explanation for the formation of coarse-grained plane beds.
269 However, this study established that the region of coarse-grained plane beds extends
270 continuously to the fine-grained region, and the conditions for fine-grained plane beds were
271 distributed separately in the two regions: the high- and low-bed shear stress. Thus, the formation
272 of the lower plane bed cannot be attributed to bed roughness or permeability owing to the grain
273 size distribution.

274 It should be noted here that bed particles are also transported by sheet flows (Sumer et al.,
275 1996), which may potentially be related to the plane bed. The sheet flow consists of a shear layer
276 of bed load and occurs at high shear stress where ripples and dunes are washed out (Pugh &
277 Wilson, 1999; Sumer et al., 1996). The laboratory experiments for sheet flows have been
278 conducted using a closed conduit in order to achieve high shear stresses and avoid the free-

surface effect (Nnadi & Wilson, 1992; Sumer et al., 1996) or under oscillatory flows to simulate storm conditions that change the beach surface (Dick & Sleath, 1992; O'Donoghue & Wright, 2004). The sheet flow regime is not described in our phase diagram because this study focuses on the unidirectional open-channel flows. Experiments for sheet flows in open-channel flows are scarce. Therefore, currently, it is difficult to determine whether the sheet flow is another possible mechanism that generates plane beds. Thus, further laboratory experiments, numerical experiments, and theoretical consideration are required to provide explanations for the lower plane bed regime.

5 Conclusions

In the present study, we compiled 937 flume and field data sets pertaining to plane beds and analyzed the dataset in the nondimensional parametric space. The results of our analysis indicate that the formation conditions of plane beds are distributed in two separate regions, and the boundary between the regions matches the threshold condition for suspension. Therefore, the upper and lower plane beds can be redefined as plane beds with and without suspension, respectively. Conversely, the formation conditions of plane beds do not show a clear boundary at the threshold condition of sediment motion, a unique value of particle diameter, or a Froude number of unity, which have been traditionally used to define lower and upper plane beds. This study demonstrates that the existence of suspension significantly contributes to the formation of plane beds, whereas other mechanisms can also produce fine-grained plane beds in flows with low-bed shear stress.

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