



22 **Abstract**

23 Micromagnetic Tomography (MMT) is a new technique that allows to determine magnetic moments  
24 of individual grains in volcanic rocks. Current MMT studies either showed that it is possible to obtain  
25 magnetic moments of relatively small numbers of grains in ideal sample material, or provided  
26 important theoretical advances in MMT inversion theory and/or its statistical framework. Here we  
27 present a large-scale application of MMT on a sample from the 1907-flow from Hawaii's Kilauea  
28 volcano producing magnetic moments of 1,646 individual grains. To assess the robustness of the MMT  
29 results, we produced 261,305 individual magnetic moments in total: an increase of three orders of  
30 magnitude compared to earlier studies and a major step towards the number of grains that is  
31 necessary for paleomagnetic applications of MMT. Furthermore, we show that the recently proposed  
32 signal strength ratio is a powerful tool to scrutinize and select MMT results. Despite this progress, still  
33 only relatively large iron-oxide grains with diameters  $>1.5\text{-}2\ \mu\text{m}$  can be reliably resolved, impeding a  
34 reliable paleomagnetic interpretation. To determine the magnetic moments of smaller ( $< 1\ \mu\text{m}$ ) grains  
35 that may exhibit PSD behavior and are therefore better paleomagnetic recorders, the resolution of  
36 the MicroCT and magnetic scans necessary for MMT must be improved. Therefore, it is necessary to  
37 reduce the sample size in future MMT studies. Nevertheless, our study is an important step towards  
38 making MMT a useful paleomagnetic and rock-magnetic technique.

39

40 **Plain language summary**

41 The magnetic information of volcanic rocks is an invaluable archive of the behavior of the Earth's  
42 magnetic field through time. Recently a new technique was proposed, Micromagnetic Tomography,  
43 that promises to determine magnetic signals of individual iron-bearing grains in these rocks. This  
44 would greatly improve our ability to obtain and interpreted the magnetic information stored in them.  
45 Here we go beyond the proof-of-concept of this exciting new technique and show that it is indeed  
46 possible to obtain statistically robust results for set of rather large grains, with diameters  $>1.5\text{-}2\ \mu\text{m}$ ,  
47 in a sample from the 1907-flow from Hawaii's Kilauea volcano.

48

49 **1. Introduction**

50 Geological materials and archeological artefacts containing magnetic particles record the direction  
51 and intensity of the past geomagnetic field as they cool. These thermoremanent magnetizations  
52 (TRMs) are our primary source of information on the behavior of the Earth's magnetic field. Obtaining  
53 reliable paleointensities and paleodirections from samples with a large variation in grainsizes is a  
54 challenge, due to differences in magnetic behavior between grains that differ in size, shape, and  
55 chemistry (e.g. Dunlop and Özdemir, 1997; Tauxe and Yamazaki, 2015). In paleomagnetic  
56 measurement techniques that rely on bulk measurements, the contributions of individual grains are  
57 measured collectively, i.e. the signals of many millions of grains result in a single magnetic moment  
58 for the entire sample. This possibly obscures information from grains that record the paleofield well  
59 by the signal of non-perfect recorders in the sample. Especially the presence of large ( $\gg 1\mu\text{m}$ ),  
60 multidomain (MD), grains often prevent a reliable interpretation of a magnetic signal from a bulk  
61 sample. Samples consisting of predominantly single-domain (SD) grains or slightly larger (but  $< 1\mu\text{m}$ )  
62 pseudo-single domain (PSD) grains with complex domain structures such as vortices or 'flower states'  
63 generally produce more reliable paleomagnetic data (e.g. Nagy et al., 2017; 2019).

64

65 Over the past decades, a number of studies have focused on high-end magnetometry techniques to  
66 assess the magnetic state of magnetic recorders and micromagnetic processes in them on a (sub)  
67 micrometer scale (e.g. Almeida et al., 2016; Farchi et al., 2017; Lima & Weiss, 2009; Nichols et al.,  
68 2016; Weiss et al., 2007). These magnetometry techniques, such as scanning SQUID microscopy (Egli  
69 and Heller 2000; Weiss et al. 2007; Lima and Weiss 2016), Electron Holography (Harrison et al. 2002;  
70 Feinberg et al. 2006; Almeida et al. 2016), and the Quantum Diamond Microscope (Glenn et al. 2017;  
71 Levine et al. 2019; Fu et al. 2020) allow to zoom in on individual grains or magnetically well-behaved  
72 small regions in a sample, and thus can avoid magnetically adverse behaved regions or grains.

73

74 The recently developed Micromagnetic Tomography method (MMT, de Groot et al. 2018; 2021) builds  
75 on this collection of magnetometry techniques. MMT proposes to overcome the differences in  
76 recording properties between grains of various sizes in a bulk sample, by separating the contributions  
77 of individual grains to the bulk magnetic signal. It relies on supplementing the results of scanning  
78 magnetometry on the surface of a sample with spatial data of the magnetic recorders within a non-  
79 magnetic matrix. This enables a three-dimensional interpretation of the magnetometry through a  
80 least-squares inversion that allocates magnetizations to individual magnetic grains. The spatial data  
81 on the magnetic recorders is acquired from X-ray computed tomography (MicroCT) scans.

82

83 As recognized by de Groot et al. (2021), the development of MMT is promising, but currently there  
84 are major challenges left to solve before MMT can be routinely used for paleomagnetic and rock-  
85 magnetic studies. Most previous MMT studies used synthetic samples and produced results for a  
86 limited amount of grains (<150) in them. Here we present the first results from MMT applied to a  
87 natural volcanic sample. We show (1) that it is possible to determine magnetic moments for large,  
88 MD, grains in our sample, (2) that the current computational setup is capable of solving for a  
89 statistically relevant number of grains and (3) which challenges currently remain and will be the topic  
90 of future research. Despite our progress, our MMT results cannot yet be interpreted in  
91 paleomagnetic terms because we only solve for the moment of large, MD, grains. To obtain  
92 information from the paleomagnetically more relevant, PSD, grains, the resolution of the MicroCT scan  
93 used must be increased. Nevertheless, our study provides an insight into how the development of  
94 MMT for paleomagnetic uses can progress, and we illustrate how MMT data can be scrutinized and  
95 selected based on the recently proposed statistical framework by Out et al. (2022).

96

## 97 **2. Sample description**

98 For this study we selected material from the 1907 lava flow of Kilauea, Hawaii. This material was  
99 sampled as site HW03 (de Groot et al. 2013), and its rock-magnetic properties were described in depth  
100 (Ter Maat et al. 2018). Our sample is a cylindrical cutout from a thin section with a sample layer of 30  
101  $\mu\text{m}$  thick, with a diameter of 3 mm. The same sample was already used for one of the case studies in  
102 de Groot et al. (2021). HW03 consists of tholeiitic basalt with minor alteration, consisting of a low  
103 percentage (< 5%) iron-oxides.

104

### 105 *2.1 Petrography*

106 The iron-oxides are copious and present throughout the sample. They are <30  $\mu\text{m}$  in size, have a cubic  
107 shape and have experienced minor oxyexsolution (Fig. 1 and Ter Maat et al. 2018). A micrograph of  
108 the thin section shows that the sample contains areas with a lower number of relatively large iron-  
109 oxide grains (> 3 $\mu\text{m}$ , 'A' in Fig. 1a) and areas with a higher number of relatively small (< 3  $\mu\text{m}$ , 'B' in  
110 Fig. 1a) iron oxide grains. The circular cut-out that was used for this study contains an area with mainly  
111 relatively large grains, and thus not so many of them. The material surrounding the iron-oxide grains  
112 has a porphyritic structure with large (100-200  $\mu\text{m}$ ) clinopyroxene and plagioclase crystals and  
113 contains some equant olivine phenocrysts. The remaining material mainly consists of fine-grained  
114 clinopyroxene, plagioclase, and glass (Fig. 1b, c, and d).

115

116 2.2 Chemical and magnetic analysis

117 Ter Maat et al. (2018) chemically analyzed 22 iron-oxide grains from HW03 using a microprobe. Based  
118 on the microprobe analysis the grains can be roughly divided in two mineralogical families. The first  
119 mineralogical family is ilmenite with a relatively equal Ti and Fe content of ~19 atom% each (6 grains).  
120 The other family is titanomagnetite with a high Fe and low Ti content varying in ratio from 3:1 to 7:1  
121 (15 grains). Lastly, there is one grain of which two spot readings show that it is partly ilmenite and  
122 partly titanomagnetite. Further investigation by Scanning Electron Microscope (SEM) using  
123 Backscatter Electrons (BSE) to characterize the crystallographic lattice shows mineralogical domains  
124 within grains that are either titanomagnetite (cubic) or ilmenite (hexagonal). Within these  
125 mineralogical domains magnetic domains can exist, but it is important to keep in mind that ilmenite  
126 is paramagnetic and titanomagnetite is ferrimagnetic at room temperature (Readman and O'Reilly  
127 1972), i.e. only the titanomagnetite grains may hold a remanent magnetization. Investigation by  
128 Magnetic Force Microscopy (MFM) of four grains indeed showed multiple magnetic domains within

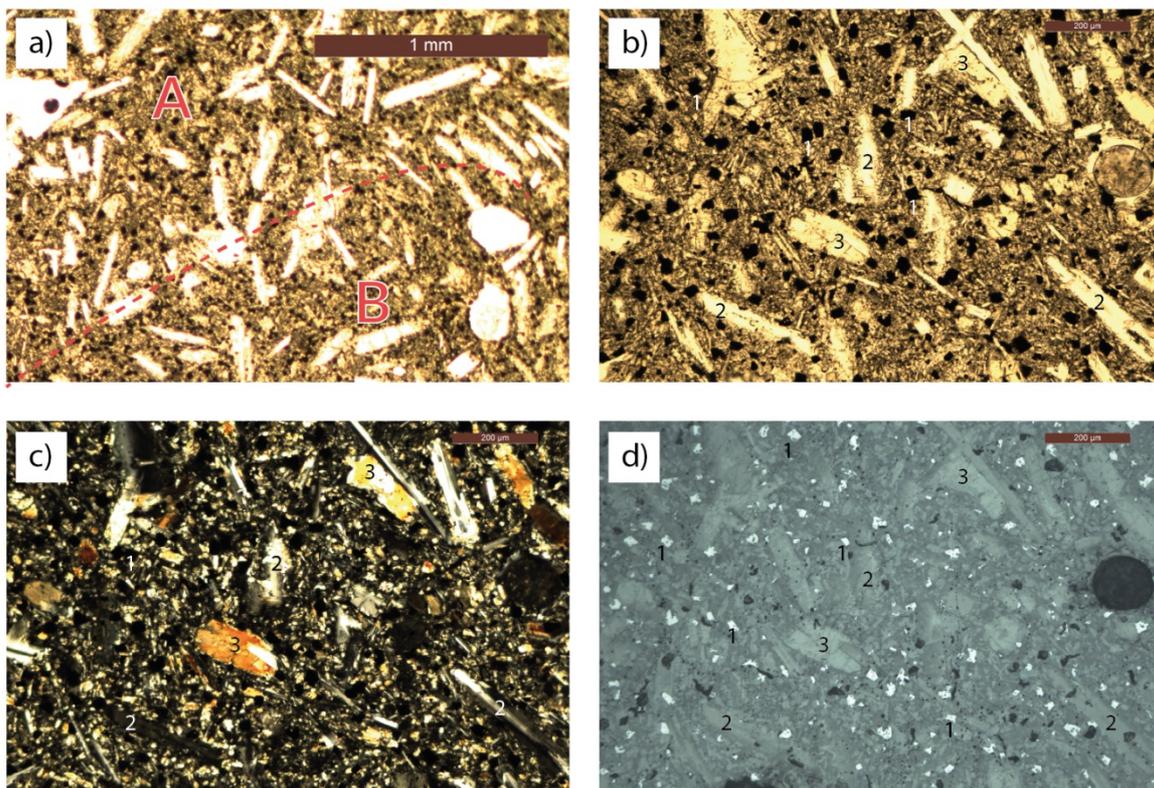


Figure 1 - Petrology and mineralogy of HW-03. Panel a shows a portion of the sample under plane polarized light (PPL) in which two areas A and B are loosely defined. Area A is an area with a sparse amount of large magnetite grains and area B an area with a higher density of smaller magnetite grains. The magnification in panels b, c, and d is larger and show the sample under PPL, crossed polarized light (XPL) and reflected light, respectively. Each shows the same field of view in which the different minerals are indicated: magnetite (1), clinopyroxene (2) and plagioclase (3). The cubic nature of the magnetite grains is especially visible under reflected light (d).

129 the titanomagnetite mineralogical domains (Ter Maat et al. 2018), but no magnetic signal in the  
130 ilmenite part of this particular grain. As it is currently impossible to discriminate between ilmenite and  
131 titanomagnetite in the MicroCT scans used to characterize the iron-oxides for our MMT study, some  
132 (parts) of the grains selected by the MicroCT analyses do not have a magnetic signal at room  
133 temperature.

134

### 135 **3. Methods**

136 The MMT method relies on combining spatial information of the location and shape of individual  
137 magnetic grains with the magnetic surface expressions of those same grains to solve their individual  
138 magnetic moments (de Groot et al. 2018; 2021). Supplementing the magnetic surface scan with spatial  
139 information is necessary to overcome the traditional non-uniqueness in potential field inversion  
140 problems (Fabian and de Groot 2019).

141

#### 142 *3.1 Spatial data acquisition*

143 The spatial information is derived from MicroCT scans that produce a three-dimensional image of the  
144 X-ray attenuation contrast in a sample that is often interpreted in terms of variations in density in the  
145 sample (e.g. Sakellariou et al. 2004). The interpretation of the three-dimensional image relies on the  
146 large attenuation contrast between the highly dense magnetic grains and the lower density  
147 surrounding matrix. Using a threshold, voxels with high density are retained, voxels with low density  
148 removed. Subsequently, adjacent, interconnected, high density voxels are grouped together to form  
149 a grain, for which the size, shape, and location are known. From this data physical properties such as  
150 volume and the distance to the surface can be calculated for each grain individually. The sample for  
151 this study was imaged by the Nanotom-S MicroCT at TU Delft. The scan had a resolution of 0.75  $\mu\text{m}$   
152 (voxelsize: 0.75x0.75x0.75  $\mu\text{m}$ ) and a field of view of 1901x2301 pixels. To suppress noise, grains that  
153 consist of 10 voxels or less were discarded, which inherently implies that the smallest grains included  
154 in the inversion routine are approximately 1.5-2  $\mu\text{m}$  in diameter. This is a major shortcoming of our  
155 MMT study, as smaller grains are undoubtedly present in our sample and are usually assumed to be  
156 the most reliable magnetic recorders. Also, it is important for the inversion routine that all magnetic  
157 sources in the system are known, otherwise there are magnetic contributions in the surface  
158 magnetometry that cannot be attributed to their source. Missing grains with diameters <1.5-2  $\mu\text{m}$  will  
159 therefore lead to less reliable MMT results. Nevertheless, the mineralogical investigation of this  
160 sample shows that many grains in our sample are >3  $\mu\text{m}$ . Furthermore, a major part of the discussion

161 of this paper consists of assessing how well the grains that are detected by our MicroCT scan are  
 162 resolved, in spite of missing smaller grains in our sample and other sources of uncertainty.

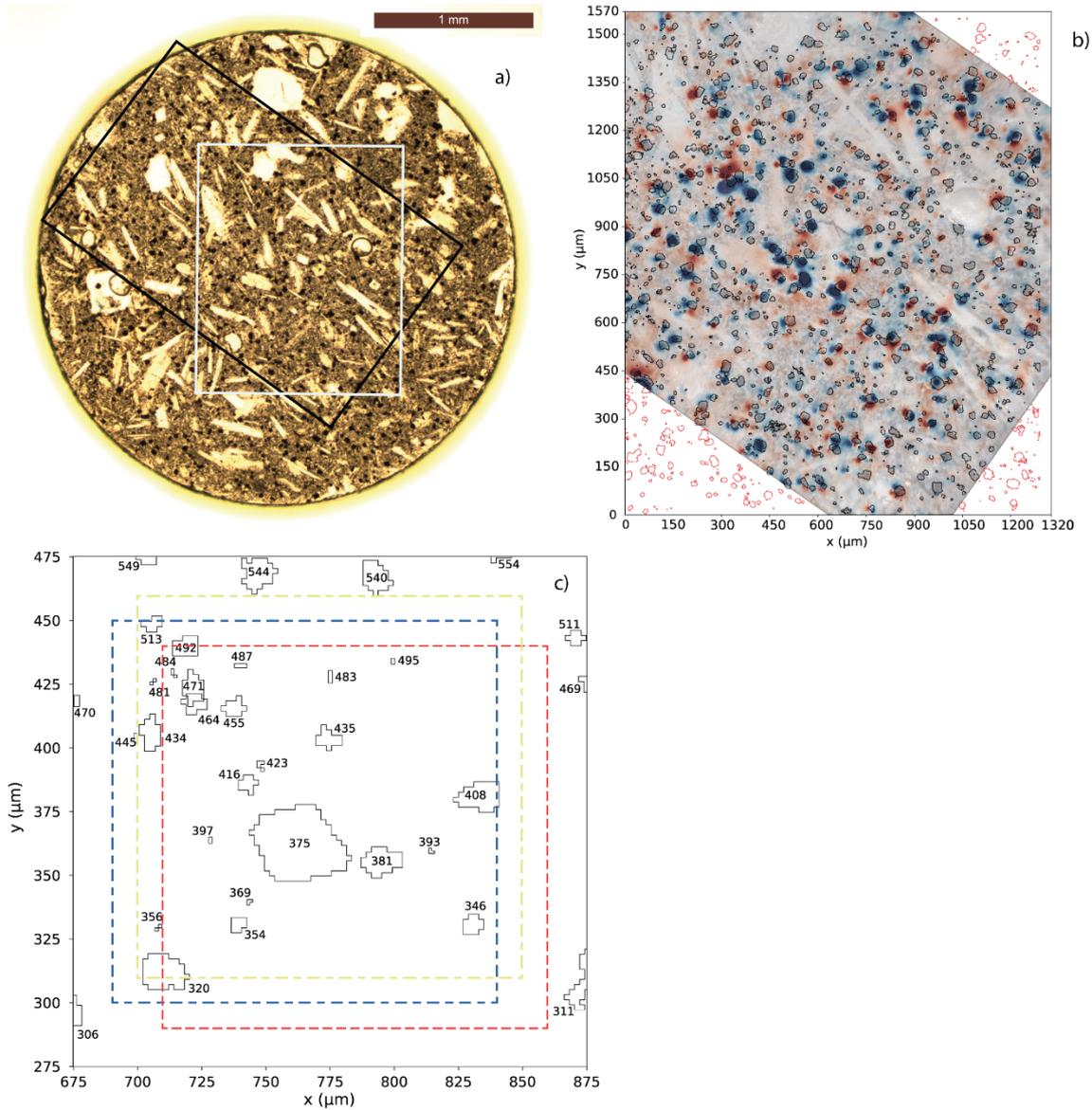


Figure 2 - The mapping of the QDM and MicroCT-scans. Panel a shows a microscope image of the entire sample, with the MicroCT field-of-view outlined in white and the QDM's field-of-view in black. In panel b the magnetic scan of the QDM is in the background (positive flux in blue, negative in red, an absolute scale lacks due to distortion from the other layer in the image), overlain by the QDM's LED image in 50% transparency. The grains as obtained from the MicroCT scan are outlined in black when they are in the QDM's field-of-view, in red when they are outside. Three unique subareas are shown in panel c, each subarea is  $150 \times 150 \mu\text{m}$  and they are spaced at  $10 \mu\text{m}$  intervals. The horizontal (x) and vertical (y) axes in panels b and c have the same origin and are in  $\mu\text{m}$ . The exclusion of grains due to intersecting the boundary of a subarea is illustrated by e.g. grain 513: it will be included in the blue and yellow subareas during the inversion but its results in the blue subarea will be discarded after the inversion was performed, because it intersects the boundary of the subarea.

163 *3.2 Magnetic surface flux*

164 The surface magnetizations for MMT can be imaged using a variety of instruments and techniques (as  
165 summarized in de Groot et al. (2021). The sample in our study was imaged by a Quantum Diamond  
166 Microscope (QDM) at Harvard (Glenn et al. 2017). The QDM determines the magnetic flux above the  
167 surface of a sample by measuring (dips in) fluorescence arising from nitrogen-vacancy (NV) centers in  
168 a diamond chip. The magnetic data is thus acquired from an optical image with a field of view of  
169 1920x1200 pixels and a spatial resolution of 1.2  $\mu\text{m}$ . During operation a 0.9 mT bias field is applied; its  
170 polarity is switched many times during measurement. Potentially induced (paramagnetic)  
171 magnetization in the sample can be removed from the remanent (ferrimagnetic) part of the signal by  
172 taking the average of the images acquired with switched polarity. The coercivities of naturally  
173 occurring iron-oxide grains are generally  $\gg 0.9$  mT (Readman and O'Reilly 1972; Dunlop and Özdemir  
174 1997); this bias field therefore would not prevent a paleomagnetic interpretation of QDM results.

175

176 *3.3 Co-registration*

177 The two datasets must be co-registered in the same coordinate system to apply the inversion routine.  
178 The QDM is an optical acquisition technique, and its camera can also be used to optically image the  
179 surface of the sample in the same coordinates as the magnetic scan. The MicroCT data can be sliced  
180 to only show the grains close to the surface of the sample. This enables co-registration based on the  
181 geometry of shallow grains (Fig. 2a). Since the two scans do not overlap entirely, not all grains imaged  
182 by the MicroCT have a co-registered magnetic flux signal (Fig. 2b). In total 1,646 grains were imaged  
183 by both the MicroCT and QDM analyses. Also, the scan height, i.e. the sample-sensor distance needs  
184 to be precisely known to properly co-register the MicroCT and QDM datasets. This distance was 6.0  
185  $\mu\text{m}$  and is derived from the actuators of the QDM set-up.

186

187 *3.4 Inversion set-up*

188 The inversion routine is identical to the inversion scheme used in (de Groot et al. 2018; 2021) and  
189 based on the theory in (Fabian and de Groot 2019). The data in those publications was processed on  
190 a desktop computer and performing the inversion routine for one area of 150x150  $\mu\text{m}$  with the grain  
191 density of our sample took approximately 2 weeks computational time. For this study, the inversion  
192 routine was highly optimized and moved to Python to use state-of-the-art libraries. Furthermore, we  
193 now use a computational server with 52 double-threaded cores and 192 GB RAM. These  
194 improvements allow to perform a single inversion of an area of 150x150  $\mu\text{m}$  in approximately 1.5  
195 minutes. Because the inversion has a high RAM demand, the size of the area that can be inverted in  
196 one calculation is still limited to 150x150  $\mu\text{m}$ . Therefore, to analyze the entire sample, it is inverted in

197 many overlapping subareas: so called ‘tiling’ (Fig. 2c). The subareas have an interspacing of 10  $\mu\text{m}$ ,  
198 leading to 118 subareas to be inverted along the x-axis and 143 subareas along the y-axis, so there are  
199 16,874 subareas in total. Tiling, however, has drawbacks. First, the inversion in subareas implies that  
200 grains can be intersected by the subarea boundary. A boundary condition of the inversion is that all  
201 magnetic sources in the system and their magnetic signals are imaged, and no sources or signals are  
202 missed. The inversion in subareas violates this boundary condition, and more so towards the edges of  
203 each subarea. Therefore, results from grains that intersect the subarea boundary are discarded after  
204 the inversion routine is performed. Grains that intersect the boundary and hence are partially in the  
205 subarea are, however, included in the inversion of that subarea to make sure their signal is not falsely  
206 attributed to other grains. Tiling also does have a major advantage. The magnetic moment of each  
207 individual grain is solved in multiple unique subareas, this enables a statistical analysis of the results.  
208 If the magnetic solution of a grain is robust, we expect only minor variations between the magnetic  
209 moments produced by the inversions of different subareas; if the results for a grain are dispersed, its  
210 solution is less accurate.

211

212 The inversion routine uses a least squares minimization to calculate the magnetic moment in three  
213 orthogonal directions per grain. The calculated magnetic moments and the spatial data from the grains  
214 are then used to create a calculated magnetic flux map at the surface in a forward model. This  
215 calculated magnetic flux map is subtracted from the original magnetic flux map to produce a map of  
216 the residuals, i.e. the measured magnetic signal that is not accounted for by the inversion.

217

### 218 *3.5 Data processing*

219 Tiling leads to multiple solutions for individual grains: on average there are 159 solutions per grain,  
220 with a minimum of 10 and a maximum of 225 solutions. From each individual solution (i.e. magnetic  
221 moment in three orthogonal components), the total magnetic moment, its declination with respect  
222 to the positive x-axis, and inclination with respect to the horizontal plane (downwards into the sample)  
223 are calculated. Also, for each individual solution, the distance to the nearest boundary for that grain  
224 in that subarea is determined. The entire dataset is subsequently grouped per grain; for each grain we  
225 determine its mean and median magnetic moments. Furthermore, the Fisher mean (Fisher 1953) is  
226 calculated, together with its precision parameter ( $k$ ) and confidence interval ( $\alpha 95$ ) for each grain. To  
227 assess the accuracy of the magnetic direction of each individual solution the angle between the Fisher  
228 mean for that grain and each of its individual solutions is determined ( $\Delta\text{Angle}$ ). It is important to note,  
229 however, that Fisher statistics only considers the directions of the magnetic moments, and not their  
230 magnitudes. To quantify the accuracy of the magnetic moment of each individual solution we calculate

231 the percentual difference between the magnetic moment of each individual solution and the median  
232 magnetic moment for that grain ( $\Delta m$ ). The  $\Delta$ Angle and  $\Delta m$  proxies can be used to investigate the  
233 stability of the results per grain between solutions stemming from different subareas.

234

### 235 *3.6 Mr/Ms ratio*

236 To determine whether an allocated magnetic moment is reasonable, the allocated magnetic moment  
237 is divided by the theoretical maximum magnetic moment for that grain. The theoretical maximum  
238 magnetic moment is calculated based on the saturation magnetization of pure magnetite (480 kA/m,  
239 e.g. Dunlop and Özdemir 1997) and the volume of the grains as produced by the MicroCT analysis.  
240 This ratio is equal to the Mr/Ms ratio, as the volume factor that converts magnetic moment to  
241 magnetization is present in both the nominator and denominator and is a constant for an individual  
242 grain. The allocated magnetic moment can never be higher than this theoretical maximum, therefore  
243 results for which Mr/Ms >1 are inherently inaccurate. Since our grains have diameters >1.5-2  $\mu\text{m}$ , they  
244 are firmly outside the single domain (SD) grainsize range for which a magnetic moment close to Mr/Ms  
245 = 1 can be expected. Realistic Mr/Ms values for individual grains are therefore well below 1, although  
246 there currently is no theoretical framework for their expected values. Furthermore, the mineralogical  
247 analyses showed that the iron-oxide grains are not pure magnetite, but titanomagnetite and ilmenite.  
248 Grains with a larger percentage of Ti may have three to four times lower Mr-values compared to  
249 magnetite (Dunlop and Özdemir 1997), and ilmenite does not hold a remanent magnetization at room  
250 temperature (Readman and O'Reilly 1972). To make it even more complex, parts of iron-oxide grains  
251 in our sample can be ilmenite while other parts are titanomagnetite (Ter Maat et al. 2018). This makes  
252 their 'magnetic grain size' smaller than the physical grain size as determined by MicroCT, lowering  
253 their expected Mr even further. All this implies that the calculated theoretical maximum magnetic  
254 moment is an absolute upper limit to the allocated magnetic moment by the inversion. For our grains  
255 it is safe to expect Mr/Ms values to be in the order of 0.1, and possibly even (much) lower.

256

### 257 *3.7 Signal Strength Ratio*

258 MMT's potential to accurately resolve a grain's magnetization depends on (1) the magnetic signal from  
259 that grain and (2) its distance to the surface of the magnetic scan (de Groot et al. 2018; 2021).  
260 Recently, Out et al. (2022) defined the signal strength ratio (SSR), a parameter that indicates the  
261 expected magnetic signal at the surface for each grain in the sample. It is based on the grain's depth  
262 in the sample (R), volume (V) and diameter (d) and calculated by:  $SSR = \frac{V}{R^3 d}$ . Furthermore, they  
263 numerically modelled the performance of MMT inversions as function of a grain's SSR and showed  
264 that the SSR can be used to select a subset of grains of which a pre-defined percentage of grains

265 produces an accurately resolved magnetic moment. The proper SSR cut-off depends primarily on the  
266 noise in the magnetic scan and the concentration of grains in the sample (Out et al. 2022).

267

#### 268 **4. Results**

269 We performed 16,874 MMT inversions to characterize the magnetic moments of 1,646 grains in our  
270 sample. Due to tiling of the subareas some grains were solved for more often than others, but 261,305  
271 unique magnetic moments were obtained in total. Previous MMT studies were limited to samples or  
272 (sub)areas with <150 grains that were inverted only once (de Groot et al. 2018; 2021). Here we gain  
273 one order of magnitude in the number of grains that were analyzed and more than three orders of  
274 magnitude in the amount of unique magnetic moments that were obtained in a single MMT study.  
275 Below we first present results for individual grains, then for single subareas, and lastly consider the  
276 entire sample - including selecting the grains that are best resolved in our study.

277

##### 278 4.1 Results for individual grains

279 As each grain is included in multiple unique subareas, the magnetic moment for each grain is  
280 determined between 10 and 225 times, with 159 times as average. In Fig. 3 we present the results for  
281 three typical grains in our sample that illustrate the differences in stability between the different MMT  
282 inversions. Grain 200 is poorly resolved: there are 225 solutions, but their directions are dispersed as  
283 indicated by a precision parameter  $k$  of 1.18 (Fig. 3a). Grain 319 performs better, with 180 solutions  
284 and a precision parameter of 5.63 (Fig. 3b). Grain 693 exhibits truly stable behavior with a  $k$ -value of  
285 363.46 with 169 solutions (Fig. 3c). We see a similar trend in the inaccuracies of the magnitude of the  
286 magnetic moments. The  $\Delta m$  parameter that gives the percentual difference between an individual  
287 solution and the length of the median solution, shows a tight distribution around 0% for grain 693,  
288 and increasingly flatter distributions for grains 319 and 200, respectively (Fig. 3d). It is important to  
289 note, however, that the deviations in direction and magnitude are not linked one-to-one, i.e. not all  
290 solutions for which the direction is poorly resolved produce an inaccurate magnitude as well, and vice  
291 versa (Fig. 3e).

292

293 To generalize these three examples, we consider the distribution of  $k$ -values for the individual grains  
294 in our sample. Grain 319 is representative for the most common precision of the results for individual  
295 grains: the median  $k$ -value in the entire population is 4.17, with the first and third quartile being 2.43  
296 and 9.57, respectively. The extremes are a grain with a  $k$ -value of 1.08 on the low end, and a grain  
297 with a  $k$ -value of 9,018.25 on the upper end.

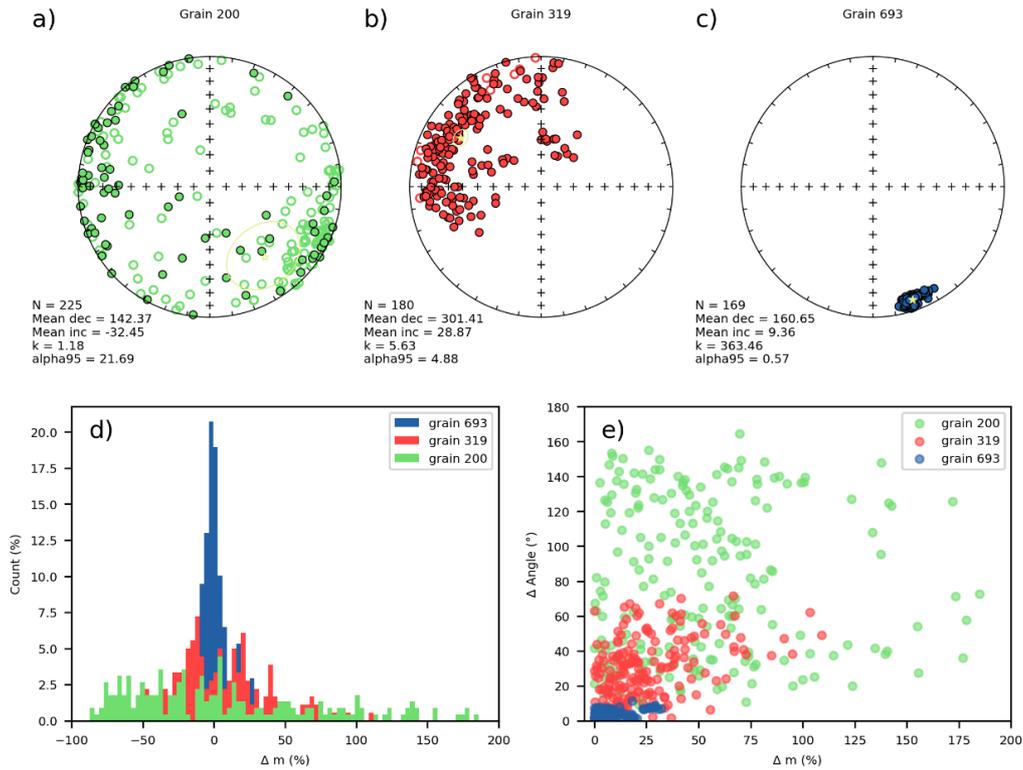


Figure 3 - Results for three typical grains. Panels a, b and c show the directional results in equal area projections for grains 200, 319 and 693 respectively; closed symbols are in the lower hemisphere corresponding to positive inclinations, open symbols are in the upper hemisphere with negative inclinations. The percentual deviation of individual magnetic moments with respect to the median magnetic moment for each grain ( $\Delta m$ ) is in d. The relationship between the angular deviation ( $\Delta \text{Angle}$ ) and  $\Delta m$  for each individual solution is in e. Note that the x-axis in panels d and e is cut-off at 200% deviation; this excludes 19 individual solutions for grain 200 from the plots.

298

#### 299 4.2. Results for an individual subarea

300 For each subarea the results consist of a list of calculated magnetic moments for the grains that are  
 301 present in that subarea, a calculated map based on the forward model of the results produced by the  
 302 inversion, and a map of the calculated residual field. The results of a typical subarea are presented in  
 303 Fig. 4. It contains 22 grains that are fully within this subarea; another 3 are intersected by its boundary  
 304 and are therefore included in the inversion, but their results are rejected. The list of solutions in this  
 305 subarea are in Supplementary Table S1. The measured magnetic flux map shows one major expression  
 306 that is not characteristically dipolar (Fig.6a, around [760,360]). Since our inversion only allocates  
 307 dipoles to the individual grains, this anomaly is only partially resolved in the forward field (Fig. 6b) and  
 308 leaves a distinct multipolar residual (Fig. 6c). The grain that is mainly beneath this multipole signal is  
 309 grain 375, which is a large and shallow grain, from which multidomain behavior can be expected.

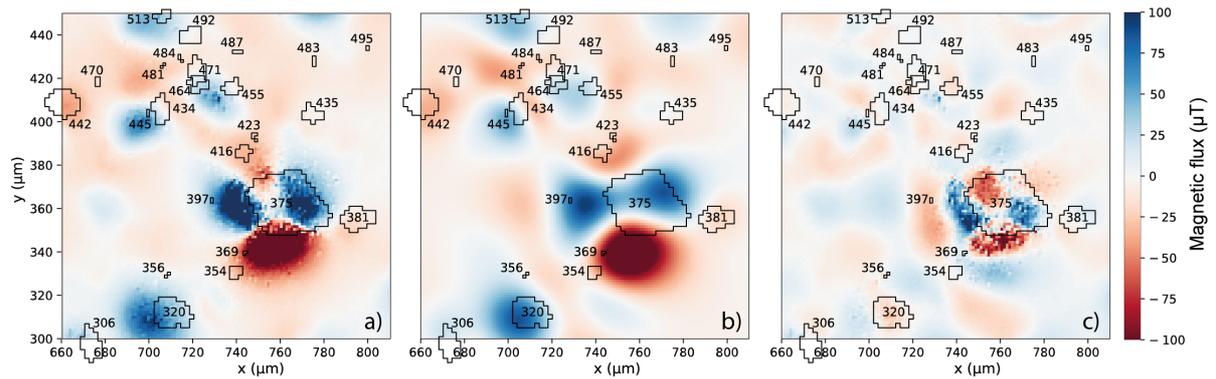


Figure 4 - Example of the results of a typical subarea; this subarea has its lower-left corner at [660,300] and its top-right corner at [810,450] in Fig. 3b. The measured magnetic flux is in panel a; the forward field based on the allocated magnetic moments in panel b; and the residual are in panel c. The grains in the subarea are outlined in black and indicated by their label (number).

310 Remarkably, the results for this grain are relatively robust between the inversions of different  
 311 subareas in which this grain is present, as its results have a k-value of 28.43. For many other smaller  
 312 grains, the residual is low indicating that the allocated magnetic moment is in line with the measured  
 313 magnetic flux by the QDM.

314

#### 315 4.3 Results for entire scanned surface

316 The overlapping field-of-view of the MicroCT and QDM scans encompasses 1,646 grains for which  
 317 261,305 unique magnetic moments were calculated in 16,874 subareas. We already showed that the  
 318 results of some grains are more stable than others, so it is paramount to select only the most reliable  
 319 magnetic moments. We therefore assess the results as function of (1) the theoretical remanence  
 320 ratios ( $M_r/M_s$ ) of the grains; (2) the distance to the closest boundary in the subarea; and (3) the SSR  
 321 of the grains (Fig. 5). The  $M_r/M_s$  ratio has a theoretical maximum of 1, but accurate solutions of the  
 322 remanent magnetization of grains in our sample are most likely in the order of 0.1 or less (section 3.6).  
 323 Many solutions from our inversions have  $M_r/M_s$  ratios  $>1$ , but the SSR of these grains is generally low,  
 324 and there is no trend between the  $M_r/M_s$  ratio of a solution and its distance to the closest boundary  
 325 of the subarea (Fig. 5a). The accuracy of the direction of the magnetic moments does show a trend  
 326 with distance to the closest boundary of the subarea (Fig. 5b). Apparently, missing parts of the  
 327 magnetic expression of a grain does hamper a reliable determination of the direction more than the  
 328 intensity. Again, the SSR is able to discriminate between accurate and less accurate results well, but  
 329 the angular deviations are still in the order of tens of degrees, even for SSRs  $>0.1$ .

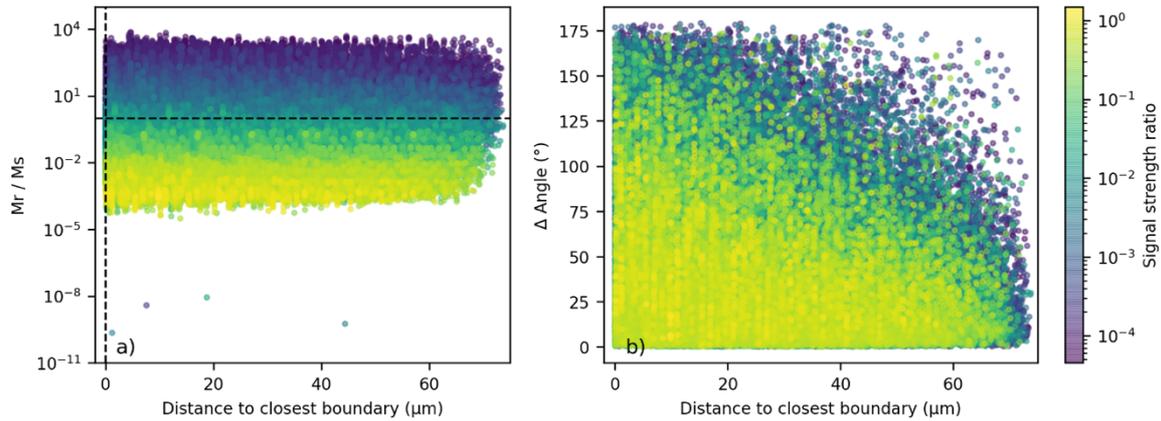


Figure 5 - Accuracy of the 261,305 uniquely determined magnetic moments from 1,646 grains as function of distance to closest boundary of the subarea and SSR. The accuracy is expressed as (theoretical)  $M_r/M_s$  ratio (a) which must be  $<1$  and is most likely in the order of 0.1 or less; and as the angle between individual results and the mean direction of a grain ( $\Delta Angle$ ) in b. In both panels the results are color coded based on the SSR of the grain. Results with a theoretical  $M_r/M_s$  ratio  $>1$  (indicated by the horizontal dashed line in a) are inaccurate as the allocated magnetization cannot be larger than the saturation magnetization.

330

## 331 5. Discussion

### 332 5.1 Mineralogy and grain sizes

333 The mineralogy of our sample was described in-depth by Ter Maat et al. (2018) and studied here  
 334 (section 2). For our purposes only the iron-oxide grains in the sample are of interest, and they are  
 335 abundant. To isolate the iron-oxides grains in the sample, we selected grains with high attenuation  
 336 contrasts from the MicroCT scan. A high attenuation contrast is often interpreted as a material with a  
 337 high density. Given the mineralogy of our sample it is likely that the vast majority of grains with a high  
 338 density are iron-oxide grains. The spatial resolution of the MicroCT scan, however, limits the size of  
 339 grains that are detected to  $>1.5\text{-}2\ \mu\text{m}$  diameter. This means that iron-oxide grains with diameters  $<1.5\text{-}$   
 340  $2\ \mu\text{m}$  that are undoubtedly present in our sample are missed in our analyses. Moreover, the  
 341 mineralogical study of our sample shows that the chemical composition of our iron-oxide grains varies  
 342 from ilmenite to (titano-)magnetite. Since ilmenite does not carry a remanent magnetization at room  
 343 temperature, not all iron-oxides in our sample will have a magnetic expression. Some grains even have  
 344 mineralogical domains of which some are magnetic at room temperature while others are not. This  
 345 implies that there are two major shortcomings in our MMT study: first, not all magnetic sources in the  
 346 system are known, and second, some sources that are identified do not carry a remanent  
 347 magnetization. It is difficult to assess the impact of these two shortcomings, and they are subject of  
 348 ongoing research. We expect, however, that the impact of missing magnetic sources on the accuracy  
 349 of MMT results is much larger than the impact of non-magnetic sources in the system. Missing a source

350 directly violates the important boundary condition of the uniqueness theorem of the inversion that  
351 states that all, but only, the magnetic sources and magnetic signals in the system must be known  
352 (Fabian and de Groot 2019). Since the inversion can allocate a near-zero magnetic moment to a source  
353 in the system, including a non-magnetic source in the inversions seems to be a lesser issue. In future  
354 studies increasing the resolution of the MicroCT scan should be considered to identify as much of the  
355 iron-oxide grains in the sample as possible. The lower limit of 50 nm below which iron-oxide grains  
356 behave superparamagnetically (Dunlop and Özdemir 1997) is challenging but not impossible for  
357 current NanoCT scanners. The field-of-view of scanners with such resolutions, however, is often  
358 narrow, limiting the volume of the sample and hence the number of grains present in a scan.

359

360 Furthermore, the variation in mineralogy of the iron-oxides in our sample hampers our interpretation  
361 based on theoretical remanence ratios ( $M_r/M_s$ ) that are based on the saturation magnetization ( $M_s$ )  
362 of pure magnetite (480 kA/m). First, the iron-oxide grains in our sample are large, multidomain, grains  
363 for which  $M_r/M_s$  ratios  $\ll 1$  are expected. Second, the iron-oxide grains in our sample are  
364 titanomagnetites or ilmenites, instead of magnetites. Since titanomagnetite has a saturation  
365 magnetization three to four times lower than pure magnetite and ilmenite is paramagnetic at room  
366 temperature (e.g. Readman and O'Reilly 1972; Dunlop and Özdemir 1997), the expected  $M_r/M_s$  values  
367 are lowered even further. Third, some grains have mineralogical domains of both ilmenite and  
368 titanomagnetite. Their volume as identified by the MicroCT scan does not represent their 'magnetic  
369 volume' at room temperature. All this implies that the calculated  $M_r/M_s$  ratios are overestimated,  
370 and that the theoretical upper limit of 1 for the  $M_r/M_s$  ratios is too high. Therefore, we expect  $M_r/M_s$   
371 values in the order of 0.1 or (much) less, although a theoretical framework for this value lacks.

372

### 373 *5.2 Geometry of MMT experiments*

374 A precise co-registration of the MicroCT and QDM datasets is paramount for the accuracy of MMT  
375 results. Small perturbations in the mapping between the two datasets may lead to increased  
376 uncertainties and/or poorly resolved grains in the MMT inversion. Also, the scan height –the sensor-  
377 sample distance that dictates the distance between the grains and the surface of the magnetic scan–  
378 needs to be known with precision, as errors in this distance propagate with the power of three in MMT  
379 results (de Groot et al. 2018).

380

381 In the current workflow the co-registration of the spatial data onto the magnetic surface scan is done  
382 by hand. This introduces an uncertainty in the mapping that may influence the results, although the  
383 grains do line up over the entire field-of-view of 1320x1570  $\mu\text{m}$ , and there are no distortions evident

384 over the surface (Fig. 2b). The optical image of the QDM greatly eases the mapping compared to the  
385 alignment of the MicroCT scan with magnetic anomalies in the Scanning SQUID Microscopy scan as  
386 was necessary in de Groot et al. (2018). Nevertheless, the mapping is time-consuming and sometimes  
387 difficult because it is done based on the geometry of shallow grains, which have a unique pattern, but  
388 which takes time to recognize in both the LED image from the QDM set-up and the surface grains from  
389 the MicroCT scan. We briefly explored the possibility of using pattern recognition software to map the  
390 surface grains to the LED image from the QDM-setup. The amount of data that the pattern recognition  
391 can use, however, is relatively scarce, because the MicroCT data consists of only the high-density  
392 grains. The results of the pattern recognition routine were not very promising, so currently, the  
393 tedious process of mapping the LED image of the QDM onto the shallow grains in the MicroCT scan by  
394 hand seems the most accurate way of co-registration.

395

### 396 *5.3 Multidomain signals of large grains*

397 The magnetic expression of large multidomain grains on the surface of the sample often is rather  
398 complex (e.g. grain 375 in Fig. 4a). Since the MMT inversion that we used here only solves for dipole  
399 magnetizations in a grain, it is likely that parts of the complex magnetization of such a multidomain  
400 grain are attributed to surrounding grains which are then assigned incorrect magnetizations. Even  
401 though the signal of these large multidomain grains may not be of interest for a paleomagnetic  
402 interpretation, solving them correctly would prevent their signal to be erroneously attributed to grains  
403 in their close surrounding. In a future study it is therefore worth exploring whether solving for non-  
404 dipolar behavior for large grains and/or grains that are close to the surface using the routines  
405 presented by Cortés-Ortuño et al. (2021) would improve the accuracy of MMT results. Special  
406 attention should be paid to select the best degree of the spherical harmonics for each grain; expanding  
407 the inversion to also solve non-dipole signals most likely increases the computational time  
408 considerably. Solving for higher order degrees of spherical harmonics introduces more variables per  
409 grain, reducing the amount of datapoints in the magnetic scan per variable to solve. Lastly, it would  
410 require a higher signal-to-noise ratio in the magnetic scan, since small errors can be accommodated  
411 in more detailed descriptions of a grain's magnetization, while they would be averaged out in a dipole  
412 approximation.

413

### 414 *5.4 Effect of the subarea size*

415 The subarea size is limited to 150x150  $\mu\text{m}$  due to RAM memory requirements. The distance to the  
416 nearest subarea boundary influences the stability of the directional solution more than the stability of  
417 the length of the magnetic vector (Fig. 5). To explain the dependency of the stability of the solution to

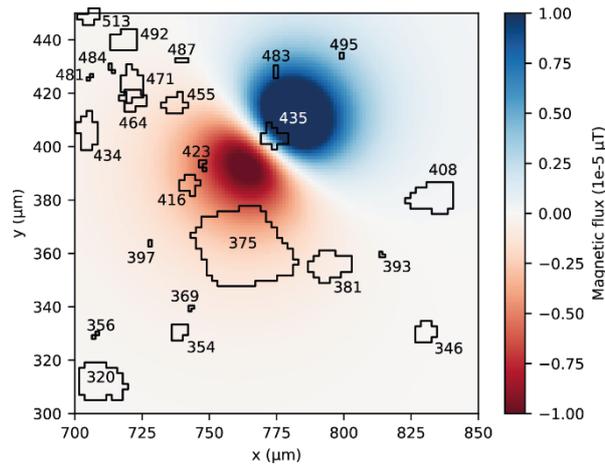


Figure 6 - The forward signal of grain 435 is present in a large part of the subarea. The contributions of the other grains are removed. Note that the color scale is exaggerated with respect to the other figures in this manuscript.

418 the distance to the closest boundary we determined the size of the surface area across which a single  
 419 magnetic grain expresses its signal. We therefore calculated the forward signal of a small grain with a  
 420 diameter of 2-3  $\mu\text{m}$  (grain 435) while ignoring the magnetic signals of the surrounding grains (Fig. 6).  
 421 Its magnetic signal on the surface is roughly circular with a diameter of 100  $\mu\text{m}$ , it is therefore present  
 422 in a major part of the subarea. Larger grains will exhibit even larger magnetic stray fields that  
 423 encompass larger surface areas. This implies that when using 150x150  $\mu\text{m}$  subareas (major) parts of  
 424 the magnetic expressions of grains will be outside the subarea and are not considered in the inversion  
 425 of that particular subarea. This violates a boundary condition of the uniqueness theorem that  
 426 stipulates that all but only the magnetic sources and their magnetic signals in the system must be  
 427 known (Fabian and de Groot 2019). It would therefore be desirable to work with larger subareas, or  
 428 even invert for the entire sample at once. This will be pursued in a future study, when the  
 429 computational limitations currently impeding larger (sub)areas are resolved.

430

### 431 5.5 Signal strength ratio cut-off

432 The signal strength ratio (SSR, Out et al. (2022)) provides an indication of the expected signal at the  
 433 surface for each grain based on its size and location in the sample. The SSR is thus a first order estimate  
 434 of the signal-to-noise ratio that can be expected for a grain, and as such an indication of how well the  
 435 MMT inversion will be capable of solving its magnetic moment. Out et al. (2022) proposed to use SSR  
 436 cut-offs above which MMT results can be trusted and illustrated that the SSR can be chosen such that  
 437 a predefined percentage of the grains that are selected is accurately solved for by MMT.

438

439 Here we can test the performance of SSR cut-offs empirically on a natural sample for the first time. To  
 440 determine the proper SSR cut-off for our dataset we consider the theoretical  $M_r/M_s$  ratios of the

441 grains as function of their SSR (Fig. 7a). The SSR is governed by the volume of a grain and its depth in  
442 the sample, grains that are deeper in the sample and/or have a smaller volume have a lower SSR and  
443 are more difficult to resolve. SSR cut-offs of 0.01, 0.02, and 0.03 all eliminate almost all solutions with  
444  $Mr/Ms > 1$ . A SSR of  $\geq 0.01$  accepts some smaller but shallow grains that are poorly resolved – possibly  
445 because they are near larger grains for which the MMT assumption to solve for dipoles breaks down.  
446 A cut-off at 0.03 on the other hand rejects some larger, deeper grains that are properly resolved. A  
447 SSR of 0.02 therefore seems to be an optimal cut-off for our dataset (Fig. 7a). Remarkably, Out et al.  
448 (2022) predicted this SSR value  $\geq 0.02$  in one of their computational models that has corresponding  
449 characteristics to our natural sample in terms of grain density, noise level and scanning geometry. In  
450 the numerical model, this SSR selects a subset of grains of which 99% are solved within 1% of their  
451 known values. On the other hand, this SSR only selects 42.9% of the grains that did reproduce their  
452 magnetic moment in the numerical model; i.e. 57.1% of the grains that are properly resolved by the  
453 inversion were unfortunately rejected.

454

455 In our sample, a SSR of  $\geq 0.02$  contains 419 grains (Fig 7b). Only one of these grains has results for  
456 which  $Mr/Ms \geq 1$ , namely grain 1870. Inspection of this grain shows that this is most likely a portion of  
457 a larger grain that was not completely imaged by the MicroCT scan. It has a very small volume ( $3.4$   
458  $\mu\text{m}^3$  or 11 voxels), and there are two other grains very close to it. For the directional results the SSR  
459  $\geq 0.02$  eliminates results over the whole range of the dataset (Fig 7c). However, when comparing Fig.  
460 5b (directional scatter for all grains) and Fig. 7c (directional scatter for grains with an SSR  $\geq 0.02$ ) it is  
461 evident that the cut-off value does reject many grains with a high angular deviation. To assess the  
462 performance of using the SSR as cut-off further, we determined the median  $\Delta\text{Angle}$  and  $\Delta m$   
463 parameters, and the precision parameter  $k$ , for the entire set of solutions (1,646 grains) and the  
464 solutions that are associated with an SSR  $\geq 0.02$  (419 grains). The median  $\Delta\text{Angle}$  changes from  $25.0^\circ$   
465 to  $17.8^\circ$ , the median  $\Delta m$  goes from 22.7% to 15.2%, and the median  $k$  changes from 4.17 to 8.07. This  
466 illustrates that the SSR cut-off indeed selects grains that exhibit more stable behavior in the inversions  
467 of the different subareas in which it is included, and that the SSR is a powerful tool to scrutinize MMT  
468 results from natural samples.

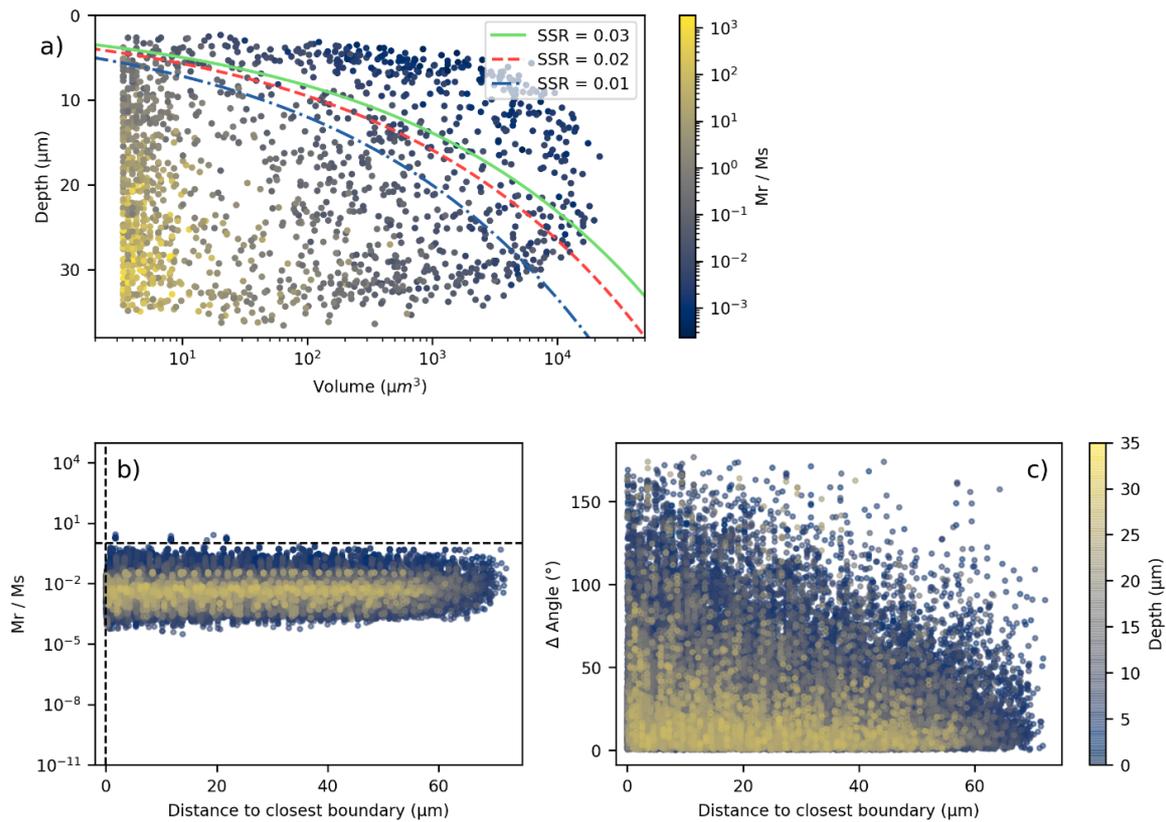


Figure 7 - The SSR is a powerful tool to select MMT. In the  $Mr/Ms$  ratio for each grain present in the sample is indicated as function of its volume and depth in the sample; three SSR cut-offs are indicated and show which grains are selected. The grains with a  $SSR \geq 0.02$  are in b and c; the  $Mr/Ms$  ratio (b) and  $\Delta Angle$  (c) are given as function of distance to closest boundary and color coded by their depth in the sample. Panels b and c are similar to panels a and b in Fig. 5, but only show data points with a  $SSR \geq 0.02$ .

469

## 470 5.6 Paleomagnetic interpretation

471 The paleomagnetic potential of MMT is based on the assumption that a certain subset of grains in a  
 472 sample contains the reliable remanent signal that is indicative of the past state of Earth's magnetic  
 473 field. Beyond SD grains, for which (Berndt et al. 2016) provided important boundary conditions, it is  
 474 enigmatic which grains are reliable recorders of the paleofield and if so, how many of such grains are  
 475 necessary to provide a meaningful statistical ensemble for a paleomagnetic interpretation. Grains in  
 476 the PSD realm ( $<1 \mu\text{m}$ ) get increasingly more attention as possible stable recorders because of their  
 477 potential vortex states (e.g. Nagy et al. 2017; 2019). Multidomain grains are often regarded as  
 478 unreliable recorders over (geologic) time scales (e.g. de Groot et al. 2014). The detection limit of the  
 479 MicroCT scan used here ( $>1.5\text{-}2 \mu\text{m}$ ) only allows to characterize large, MD, grains in our sample. This  
 480 implies that we miss gains that are of SD or PSD nature –the grains that are often believed to be more  
 481 reliably recorders of the Earth's paleofield– in our study. In future studies the resolution of the  
 482 MicroCT scan must be lowered by approximately one order of magnitude, before a meaningful

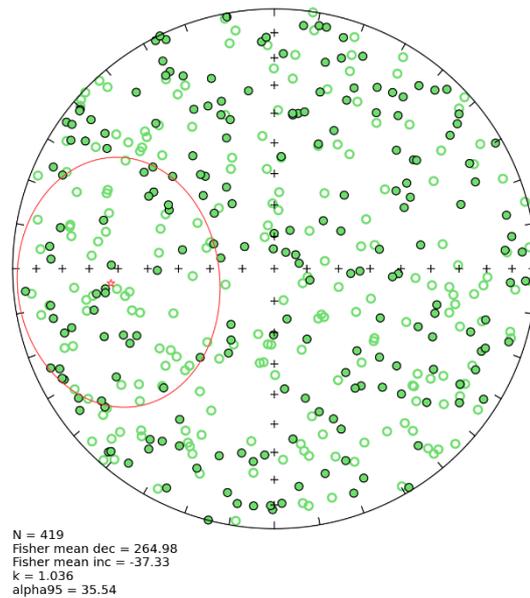


Figure 8 - The mean directions of 419 grains with a SSR  $\geq 0.02$  in an equal area projection; closed symbols are in the lower hemisphere corresponding to positive inclinations, open symbols are in the upper hemisphere with negative inclinations. The Fisher mean of these grains is indicated by a red star, together with its 95% confidence interval as red line.

483 paleomagnetic interpretation of MMT results will. be possible. Technically, resolutions  $< 500 \mu\text{m}$  are  
 484 already achievable in NanoCT/MicroCT scans, but such scans have a limited field-of-view, lowering the  
 485 amount of grains in the scan. Therefore the 'sweet-spot' in the trade-off between field-of-view and  
 486 MicroCT resolution should be determined by the research question and the physical characteristics of  
 487 the sample material to optimize the paleomagnetic interpretation of MMT results.

488

489 Although an interpretation in paleomagnetic terms is not possible, we can use the MMT solutions of  
 490 the grains that have a SSR  $\geq 0.02$  to calculate their dispersion in directions. Since all grains analyzed in  
 491 our study underwent the same magnetic history since solidification, their magnetizations are expected  
 492 to be the result of the same external magnetic field(s). Their magnetic dispersion therefore gives a  
 493 first-order estimate of how many large, MD, grains would be necessary to produce a consistent  
 494 paleomagnetic direction. The Fisher mean for the 419 grains with a SSR  $\geq 0.02$  gives a declination of  
 495  $265.0^\circ$  and an inclination of  $-37.3^\circ$ , with a precision parameter  $k$  of 1.036 and confidence interval  $\alpha_{95}$   
 496 of  $35.5^\circ$  (Fig. 8). As the number of grains, the precision parameter and the confidence interval are  
 497 related (e.g. Butler 1992), the number of grains required to get a smaller confidence interval with the  
 498 current precision parameter can be calculated. An uncertainty of  $10^\circ$  or less is often considered  
 499 reliable (Berndt et al. 2016); this implies that given on our  $k$  of 1.036, 5,430 grains would be required  
 500 to attain a confidence interval of  $10^\circ$ . This is 12.9 times more than the current dataset of 419 grains.  
 501 Our 419 grains are produced by a scanned surface of  $1.52 \text{ mm}^2$ . To acquire data from 5,430 grains  
 502 after applying our MMT inversion and SSR cut-off a scan surface of  $19.6 \text{ mm}^2$  would be required. This

503 is in remarkable agreement with the prediction by de Groot et al. (2021) that scanning an area of 20  
504 mm<sup>2</sup> would be necessary for a 'reliable' result in a volcanic sample, although it must be emphasized  
505 again that the MD grains in our analyses are most likely not reliable paleomagnetic recorders, at  
506 least not on geological timescales.

507

## 508 **6. Conclusion**

509 We presented a large-scale application of MMT on a natural sample producing magnetic moments of  
510 1,646 individual grains in our sample. Due to tiling in our inversion routine, we obtained on average  
511 159 magnetic moments for each grain, so our study produced 261,305 individual magnetic moments  
512 in total. This enabled a statistical assessment of the results using the recently proposed signal strength  
513 ratio. After selecting the most reliable MMT results using a SSR cut-off of 0.02, we obtained robust  
514 results for 419 rather large and/or very shallow grains in our sample. Previous MMT studies produced  
515 magnetic moments for <150 grains. We therefore gain one order of magnitude in the number of grains  
516 that were analyzed and more than three orders of magnitude in the amount of unique magnetic  
517 moments that were obtained in a single MMT study. The most important recommendation that arises  
518 from our findings is that samples for MMT studies should be smaller than the sample with a diameter  
519 of 3 mm that we used here. If the diameter of the sample would be in the order of 1 mm it would be  
520 possible to (1) measure the magnetization of the entire surface in one QDM scan, and (2) fit the sample  
521 in the field-of-view of MicroCT scanners with resolutions <500 nm. This would imply that we could  
522 determine the magnetic moments of smaller (< 1 μm) grains that may exhibit PSD behavior and are  
523 therefore better paleomagnetic recorders than the MD grains we currently can analyze. Also, if a  
524 better optimization of the MMT inversion would allow to invert the magnetic scan of such a 1 mm  
525 sample at once, we would satisfy the boundary condition of the MMT inversion theory that all, but  
526 only, magnetizations arising from the sources in the sample must be measured. This would remove  
527 the need for tiling during the MMT analyses and undoubtedly lead to better MMT results.  
528 Nevertheless, our study is an important step towards making MMT a useful paleomagnetic and rock-  
529 magnetic technique.

## 530 **Data statement**

531 The data used in this study has been uploaded to the Pangaea.de repository and will be available soon.  
532 Pending the FAIR data check by Pangaea we uploaded the data for peer review.

## 533 **Acknowledgements**

534 The authors declare no competing interests. This project has received funding from the European  
535 Research Council (ERC) under the European Union's Horizon 2020 research and innovation

536 programme (Grant agreement No. 851460 'MIMATOM' to LVdG). LVdG acknowledges funding from  
537 the Dutch Science Foundation (NWO) grant ALWOP.641.

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