

The Impenetrable Barrier: Suppression of Chorus Wave Growth by VLF Transmitters

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Key Points:

- Coherent VLF transmitter signals can entrain resonant electrons at the outer limit of the VLF bubble outside the contracted plasmasphere.
- Nonlinear growth of chorus rising tones and subsequent acceleration of 100s keV electrons to MeV energies is suppressed in that region.
- Suppression of local MeV acceleration contributes to diffusive and loss-related effects in establishing the impenetrable barrier at $L=2.8$.

Abstract

Rapid radiation belt recovery following storm time depletion involves local acceleration of multi-MeV electrons in nonlinear interactions with VLF chorus waves. Previous studies of an apparent impenetrable barrier at $L \sim 2.8$ focused on diffusion and precipitation loss mechanisms for an explanation of the sharp reduction of multi-MeV electron fluxes earthward of $L \sim 3$. Van Allen Probes observations for cases when the plasmasphere is contracted earthward of $L \sim 3$ indicate that strong coherent signals from VLF transmitters can play significant roles in the suppression of nonlinear chorus wave growth earthward of $L \sim 3$. As a result, local nonlinear acceleration of 100s keV electrons to MeV energies does not occur in this region. During the recovery of the outer radiation belt when the plasmasphere is significantly contracted, the suppression of chorus wave growth and local acceleration by the action of the transmitter waves at the outer edge of the VLF bubble contributes to the sharp inner edge of the new MeV electron population and the formation of the impenetrable barrier at $L \sim 2.8$.

Plain Language Summary

The variability and intensity of relativistic electrons in earth's outer radiation belt pose serious space weather concerns. Highly-relativistic (MeV) electrons are accelerated locally in rapid nonlinear interactions with very low frequency (VLF) naturally-generated rising tone chorus waves. Chorus risers, in turn, grow through interactions with 10s keV electrons injected earthward from the outer magnetosphere during storm conditions. Most of the MeV electrons are lost early in such storms and are seen then to recover rapidly through this local acceleration process. Strong man-made signals from VLF transmitters extend outward in space forming a bubble of wave energy surrounding Earth and confined by its magnetic field. At the outer limits of the bubble, the transmitter waves interact strongly with the 10s keV electrons preventing the growth of the chorus waves that act to accelerate electrons to MeV energies. These effects act to establish a sharp earthward boundary of the rapidly recovering radiation belt MeV electron population at a distance ~ 2.8 earth radii from our planet. This persistently observed boundary has been termed the impenetrable barrier.

1 Introduction

Dual toroids of energetic protons and electrons circling the Earth were discovered early in the Space Age - the Van Allen radiation belts. (cf Van Allen et al., 1958). Lying earthward of $L \sim 2$, a relatively stable inner zone consists of multi-MeV protons. While fluxes of 10s - 100s keV electrons can extend into the slot region ($2 < L < 3$), relativistic (> 500 keV) and highly-relativistic (> 3 MeV) electrons in the outer radiation belt are typically confined to $L > 3$. The characteristics and dynamics of the multi-MeV electron fluxes are strongly dependent on storm time processes leading to pronounced loss often rapid recovery of their outer radiation belt populations (e.g. Ukhorskiy et al., 2014; Reeves et al., 2013). As a result, the high energy outer belt electron fluxes vary significantly in their spatial extent and on multiple time scales.

Nonlinear processes play a significant role in determining the dynamic variability of relativistic electrons in Earth's outer radiation belt. Despite drastic changes in the geomagnetic field configuration during solar storms that result in an almost total depletion of the MeV outer belt electrons, a rapid recovery of the outer zone can subsequently take place in a matter of a few hours (e.g. Baker et al. (2014a)). In the following study we use Van Allen Probes (Mauk et al., 2012) observations to examine intercoupled the roles played by cold plasmaspheric electrons

(Electric Field and Waves (EFW) (Wygant et al., 2013), keV electrons (Magnetic Electron Ion Spectrometer (MagEIS) (Blake et al., 2013), and plasma waves (Electric and Magnetic Field Instrument and Integrated Science (EMFISIS) (Kletzing et al., 2012) in the recovery of multi-MeV electrons (Relativistic Electron-Proton Telescope (REPT) (Baker et al., 2012).

Radiation belt recovery in the inner magnetosphere involves local acceleration of 100s keV seed electrons to multi-MeV energies in the low-density region outside the plasmopause (Reeves et al., 2013) through interactions with whistler mode very low frequency (VLF) chorus waves (Thorne et al., 2013; Li et al., 2016). Jaynes et al. (2015) have shown that magnetospheric substorm activity produces two distinct electron populations that are essential to the acceleration of highly relativistic electrons in the outer belt: a source population of 10s keV electrons that drive VLF wave growth, and a seed population at 100s keV that are then accelerated by the VLF waves to MeV energies. The recovery of multi-MeV electrons at $L \sim 3 - 5$ can take place rapidly (30-60 min), indicating the importance of nonlinear processes (Foster et al., 2014). In the following sections we present Van Allen Probes observations that indicate that, for cases where the plasmopause moves strongly inward, strong coherent VLF transmitter signals at the outer edge of a VLF bubble surrounding Earth (Foster et al., 2016) can contribute to the suppression of nonlinear chorus wave growth and local MeV electron acceleration and the formation of a seemingly impenetrable barrier for the occurrence of multi-MeV electrons earthward of $L \sim 2.8$.

2 The Impenetrable Barrier

Baker et al., (2014b) found that highly relativistic (>2 MeV) and ultra-relativistic (>5 MeV) electrons in the outer radiation belt were found at $L \gtrsim 3$, but were not observed at lower L values. Considering that diffusive radial transport and pitch angle diffusion into the loss cone proceed very slowly for relativistic electrons at $2.5 < L < 3$ (Abel & Thorne, 1998), Baker et al. (2014b) theorized that very slow inward radial diffusion combined with loss through weak wave-particle pitch angle scattering were the most likely causes of an apparent barrier for relativistic electrons at the earthward edge of the outer radiation belt near $L=2.8$.

Ozeke et al. [2018] presented simulations of the action of ultra-low frequency wave radial diffusion on the formation of the apparent impenetrable barrier during the recovery phase of geomagnetic storms. Their simulations result in the creation of a very steep radial gradient in flux at the inner edge of the outer ultra-relativistic radiation belt, and find that the location of this gradient changes with time and from storm to storm. However, for the period between September 2012 and May 2014 there was no evidence of penetration of ultra-relativistic electron flux inward of $L \sim 2.8$ in the simulation results. They conclude that the apparently impenetrable barrier can be explained naturally by inwards ULF wave radial transport. The overall results of Ozeke et al. [2018] are similar to the initial conclusion reached by Baker et al. [2014] that the apparent barrier at $L \sim 2.8$ resulted from very slow inward diffusive transport rates around $L \sim 2.8$ in the inner magnetosphere combined with slow wave-particle scattering losses of MeV electrons to the atmosphere due to chorus and plasmaspheric hiss waves.

There are exceptions to the impenetrability of the barrier near $L = 2.8$. In particular, interplanetary shocks can accelerate electrons adiabatically to multi-MeV energies in the inner magnetosphere on a 1-min time scale (Foster et al., 2015), and well inside $L=2.8$ in extreme cases (Blake et al., 1992). In the absence of subsequent strong disturbances, the initially sharp multi-MeV electron flux boundaries that formed outside $L \sim 2.8$ during the radiation belt recovery from the 17-18 March 2015 storm were observed to diffuse gradually earthward across the

barrier to $L \sim 2.5$ over the ensuing months (Baker et al., 2016). The evolution of this boundary for 3.4 MeV REPT electrons is shown in Figure 1b, eventually forming an inner long-lived "storage belt" such as described by Baker et al. (2013).

3 Effects of Anthropogenic VLF Signals

Strong signals at 20 kHz – 25 kHz from naval VLF frequency radio stations provide one-way communication to submerged submarines. In addressing a possible cause for the impenetrable barrier, Baker et al. (2014b) considered precipitation of energetic electrons through interactions with these transmitter signals. Following on the work of Abel and Thorne (1998), Baker et al. concluded that precipitation loss would only be significant for electrons with energies < 0.5 MeV and furthermore would only be important at mid-latitude locations where $L < 2$.

For the 17-18 March 2015 storm event discussed above (cf Figure 1b), Foster et al. (2016) addressed characteristics of the magnetically confined bubble of very low frequency (VLF) wave emissions of terrestrial, human-produced origin that surrounds the Earth. VLF whistler-mode signals are confined to L shells such that the wave frequency remains $\leq \frac{1}{2}$ the minimum electron gyrofrequency (f_{ce}) along the field line. The outer limit of the VLF bubble at the strong 25 kHz VLF transmitter frequency closely matches the position of the barrier to the inward extent of multi-MeV radiation belt electrons described by Baker et al. (2014b). Foster et al. (2016) showed an observed 1000x increase in VLF electric field amplitude at the transmitter frequency seen outside both the $\frac{1}{2} f_{ce}$ boundary and the plasmapause, and attributed this effect to wave growth stimulated by the presence of the transmitter signal. An ample energy source for the amplification of VLF emissions near the transmitter frequency was provided by abundant fluxes of 200 keV electrons that extended well earthward of the MeV electrons. Cyclotron ($n = 1$) interactions leading to VLF wave growth take perpendicular energy from the resonant electrons leading to their precipitation loss to the atmosphere (cf. Foster and Rosenberg (1976); Inan et al. (2007)). A recent statistical and modeling study (Ma et al., 2017) has addressed the crucial role played by VLF waves from transmitters in energetic electron loss at $L < 2.5$. Direct evidence for the precipitation loss of the ~ 100 keV electrons resonant with the VLF transmitter frequency during the 18 March 2015 event reported by Foster et al. (2016) is presented in the Supplementary Material to this paper. Foster et al. (2016) re-examined the effects of observed strong VLF transmitter signals and evidence of wave-wave coupling near $L = 2.8$ on the formation of the impenetrable barrier. The processes identified were calculated to be sufficient to produce enhanced MeV electron precipitation loss in the region beyond the eroded plasmapause significantly greater than that estimated by Baker et al. (2014b).

Both the Baker et al. (2014b) and the Foster et al. (2016) discussions of the impenetrable barrier focused on particle loss mechanisms for an explanation of the sharp reduction of multi-MeV electron fluxes near $L = 2.8$. Baker et al. (2014b) noted that the inner edge of the ultra-relativistic electron population at $L \sim 2.8$ seemed to require that electron acceleration occur just outside that location, since the radial transport of such electrons to $L \sim 2.8$, well inside the nominal position of the plasmapause at $L \sim 4-5$, is usually very slow (years). They concluded that the appearance of robust populations of ultra-relativistic electrons at $L \sim 3$ would require a local wave acceleration process occurring just outside the plasmapause during conditions when the plasmapause was contracted to $L < 3$. In addition, the Foster et al. (2016) study pointed out that, for the March 2015 event, natural chorus band emissions appeared to have been suppressed at the

position where the transmitter frequency approached $\frac{1}{2} f_{ce}$ at the edge of the VLF bubble ($L \sim 2.8$).

In this study, we investigate observations for a moderate storm in 2017 when the plasmapause was eroded to $L < 2$. We reconsider the strong VLF transmitter signals observed near $\frac{1}{2} f_{ce}$ at the outer limits of the VLF bubble. Observations and theoretical analysis described below indicate that phase perturbation and subsequent precipitation loss of 10 - 100 keV electrons in resonance with the growing coherent transmitter signal can result in the suppression of the nonlinear growth of VLF chorus rising tone elements. Discrete chorus rising tones are essential for the rapid acceleration of MeV electrons (e.g. Omura et al., 2015a). We suggest that the suppression of local MeV acceleration through the action of the VLF transmitter signal at the outer edge of the VLF bubble can contribute significantly to diffusive and loss-related effects in establishing the impenetrable barrier near $L = 2.8$.

4 Inner Magnetosphere Plasma and Wave Observations

The general characteristics of the 8-9 September 2017 event have been described in detail in the recent literature. A double phase geomagnetic storm on 8 September 2017 (Dst max of -142 nT and -122 nT respectively), combined with multiple solar flare impacts (e.g. Yamauchi et al., 2018), led to a contraction of the plasmasphere, large variations in global total electron content, and unusual mid-latitude effects in the plasmasphere boundary layer including deep plasma depletions (Aa et al., 2018). In the outer radiation belt, the event featured a sharp dropout of MeV electron fluxes and their subsequent rapid recovery at L values > 2.8 . This is seen in the REPT observations shown in Figure 1a and Figure 2 (red curve) that show that the peak REPT 2.6 MeV electron flux between $L = 3$ and $L = 5$ decreased by a factor of 30 at the beginning of Sept. 8th and then increased by 300x by the beginning of the 9th. Whereas the plasmapause receded to $L \sim 2$ during the strong MeV electron recovery (blue curve), the inward extent of the sharp gradients at the earthward edge of the recovering fluxes formed only outside $L = 2.8$ (black curve). Plasmapause locations at ~ 18 MLT and 08 MLT (density $< 100 \text{ cm}^{-3}$; blue curves) were derived from spacecraft potential (EFW) and upper hybrid frequency measurements (EMFISIS).

The propagation of terrestrial VLF signals in the inner magnetosphere has been described by Starks et al. (2009). The primary mode of propagation for $L > \sim 2$ is ducted and restricted to a magnetically confined bubble surrounding Earth at L shells such that the wave frequency is less than $\frac{1}{2}$ the electron cyclotron frequency (f_{ce}). Figure 3a presents the EMFISIS high frequency (HFR) electric field spectrogram for the inbound pass of RBSP-A at 03:00 UT on 09 September 2017. Strong signals from VLF frequency transmitters (horizontal lines) extend to the outer regions of the VLF bubble. Wave power at the strongest 24.9 kHz transmitter frequency increased by a factor of 10^6 between $L = 3.2$ and $L = 2.7$. In that region, natural whistler mode emissions are suppressed by $\sim 1000x$ in a band from 20-24 kHz, immediately below the enhanced transmitter frequency. The sharp plasmapause is at $L=2.1$ at this time and the ratio f_{pe}/f_{ce} is < 2 earthward of $L=2.8$. Chorus rising tones near $\frac{1}{2} f_{ce}$ end abruptly inside $L = 2.8$.

The observations presented in Figure 3a for the 09 September 2017 event are qualitatively similar to those discussed by Foster et al. (2016) for the 18 March 2015 storm. In each case natural whistler mode emissions and chorus rising tones were suppressed immediately

below a strong coherent VLF transmitter signal at the outer extent of the VLF bubble at and around $L=2.8$.

Naturally generated whistler mode waves experience growth in ($n = 1$) cyclotron interactions with injected electrons (e.g. Foster et al., (1976)) through transfer of perpendicular energy from resonant electrons to the waves, driving the electrons toward the precipitation loss cone (e.g. Foster and Rosenberg (1976)). The presence of an anthropogenic, coherent VLF signal in this region also can interact with the same lower energy electron population. Evidence of such direct interaction of the transmitter signal with lower energy electrons in the chorus generation region outside the plasmopause is seen in the MagEIS observations across the outer edge of the VLF bubble shown in Figures 3b and Figure 4. The upper panel of Figure 4 presents the variation of the electric field wave power at the 24.9 kHz transmitter frequency across the region. Resonant electron energy at this frequency for the observed conditions was calculated to be ~ 60 keV. Wave power was enhanced $>10\times$ suggesting local amplification of the transmitter signal across the region between the plasmopause and the outer edge of the VLF bubble, and was decreased by a factor $> 10^5$ outside that region. In the region of transmitter wave power enhancement, the flux resonant of low energy electrons between 30 keV and 90 keV was decreased by $\sim 3\times$. These observations are consistent with loss cone scattering of the 10 - 100 keV electrons resonant with the enhanced transmitter signal. In the same region, sub-relativistic electron flux (235 keV) was increased $3\times$ at perpendicular pitch angles. The local acceleration of relativistic electrons (597 keV and 1079 keV) exhibited a sharp flux decrease by factors of 100 inside $L = 3.2$. These inter-related observations indicate that in the region at the outer extent of the VLF bubble where lower energy injected electrons interacted directly with the transmitter signal, the transmitter wave power grew while rising tone chorus development and MeV acceleration were suppressed.

For the 09 Sept. 2017 event, the MagEIS electron energy range and pitch angle coverage needed to show precipitation loss of the resonant electrons was not available in either the RBSP-A or B observations. However, for the very similar barrier event on 18 March 2015, Figure 4b of the earlier Foster et al (2016) showed that 50-300 keV electrons were resonant with the transmitter frequency at low pitch angles. For that event, RBSP-A MagEIS directly observed loss-cone fluxes of the resonant electrons. Those data are presented in the Supplementary Material.

5 Suppression of Nonlinear Chorus Wave Growth at $L < 2.8$

Initially, whistler waves grow linearly, becoming coherent as the wave grows at the frequency of the largest linear growth rate while also suppressing the growth of other waves around that frequency. When the wave amplitude exceeds a threshold amplitude for nonlinear instability, the wave amplitude grows rapidly as the frequency increases monotonically (Omura et al., 2009). Nonlinear growth stops near the optimum wave amplitude (Omura et al., 2015b) and then decreases gradually, forming a discrete subpacket within the chorus rising tone. However, this natural sequence of chorus wave development and subsequent nonlinear relativistic acceleration in the inner magnetosphere can be broken by several interrelated effects which can occur during unusual stormtime plasmasphere configurations.

In the inner magnetosphere the magnetic field strength and f_{ce} increase rapidly with decreasing L . At times when the plasmopause contracts earthward of $L \sim 3$, as in the 9 September 2017 case discussed above, the unusually low plasma density in the inner

magnetosphere can alter the conditions needed for nonlinear chorus wave growth, as these are highly dependent on the ratio f_{pe}/f_{ce} of cold plasma frequency to electron cyclotron frequency. Figure 5 presents the results of theoretical calculations (Omura et al., 2015b) that describe the dependencies of the nonlinear threshold frequency and optimal wave amplitude on f_{pe}/f_{ce} . Dashed curves indicate the threshold amplitude for the onset of nonlinear growth as a function of normalized wave frequency ($f_{wave}/f_{ce} = 0.5$ and 0.25 are indicated by a dotted black line). Solid curves indicate the optimal amplitude for the nonlinear wave growth. Threshold and optimal amplitude curves are shown for three values of f_{pe}/f_{ce} . In the case we are describing, $f_{pe}/f_{ce} \sim 1$ at the location where the VLF transmitter frequency matches $1/2 f_{ce}$ (cf. Figure 4c).

For the results of the Omura et al. (2015b) calculations shown in Figure 5a, f_{pe}/f_{ce} varies as $[3, 2, 1]$ (magenta, blue, green). The density of hot to cold electrons (n_h/n_c) = 4×10^{-4} , and the velocity distribution function has been set so that the resonance energy is 40 - 60 keV, (this is specified by parameters in Equation (1) in Hikishima et al. (2010) $U_{t, para} = 2.3$, $U_{t, perp} = 2.5$, $\beta = 0.3$, $Q = 0.5$). These calculations show that in a region where the cold plasma density is low enough such $f_{pe}/f_{ce} \sim 1$, the nonlinear threshold amplitude exceeds the optimal (maximum) amplitude for wave growth at $L \sim 2.8$, total suppressing the development of chorus rising tones.

A complementary effect is the suppression of the chorus wave growth in the presence of a strong coherent transmitter signal near the chorus wave starting frequency. The strong constant frequency transmitter signal can entrain the resonant electrons suppressing the development of rising tone emissions. Wave power grows at the transmitter frequency and resonant electrons in the 50-100 keV range are lost by precipitation. Strong evidence of the direct interaction of the transmitter signal with these lower energy electrons has been presented in the flux ‘bite-out’ seen in the MagEIS observations as shown in Figures 3b and 4.

Precipitation loss of 10s - 100 keV electrons resonant with the enhanced transmitter waves was seen in MagEIS pitch angle observations during the 18 March 2015 event (cf. Figure 1b) and is included here in the Supplementary Material. Figure 5b demonstrates the effect of the depletion of the resonant electrons (n_h) by the transmitter wave. We vary n_h/n_c as 1×10^{-4} , 2×10^{-4} , 4×10^{-4} (red, black, blue curves), while other parameters are the same ($f_{pe}/f_{ce} = 2$). As the density of hot resonant electrons decreases (blue \rightarrow black \rightarrow red), the overlap of the optimum and threshold amplitudes decreases and disappears. The decrease in density of the hot electrons around the resonant energy caused by their interactions with the transmitter signal prevents the development of the chorus waves responsible for local MeV electron acceleration. The resonant electron energy is determined from the resonance condition which depends on the ratio of f_{pe}/f_{ce} especially in the low frequency range. As shown in Figure 5d, the resonance energy becomes higher with lower f_{pe}/f_{ce} . Since the flux of the resonant electrons decreases for higher energy, generation of chorus emissions becomes unlikely at lower L . Linear growth rates for $f_{pe}/f_{ce} = 1$, 2, and 3 are plotted in Figure 5c. The linear growth rate for $f_{pe}/f_{ce} = 1$ (green) is zero or negative and does not appear in the logarithmic scale of Figure 5c.

The absence of chorus development prevents the local acceleration of 100s keV seed electrons to relativistic energies. During the recovery of the outer radiation belt when the plasmasphere is significantly contracted, the suppression of chorus wave growth and local acceleration by the rapid decrease of f_{pe}/f_{ce} combined with the action of the transmitter waves on resonant electrons at the outer edge of the VLF bubble can contribute significantly to the

formation of a sharp inner edge of the new MeV electron population and the appearance of a seemingly impenetrable barrier at $L \sim 2.8$.

6 Discussion

Foster et al. (2017); Hsieh & Omura (2018); and Omura et al. (2019) have presented the theory for nonlinear wave-particle electron acceleration in interactions with strong chorus rising tones outside the plasmopause, and have applied these calculations to Van Allen Probes observations during rapid radiation belt recovery events. Nonlinear acceleration by strong VLF chorus rising tones is sufficient to explain the prompt (30-60 min) increase of multi-MeV electron fluxes in the inner magnetosphere by factors of 10s to 100 as reported by Foster et al. (2014). A localized set of processes near $L \sim 3$ that suppresses the development of chorus rising tones would produce a sharp earthward boundary of the region of multi-MeV local electron acceleration and flux enhancement that Baker et al. (2014b) have described as the impenetrable barrier. Near $L \sim 3$ in the low-density region outside the plasmopause the ratio f_{pe}/f_{ce} is decreasing rapidly and the nonlinear starting frequency of the chorus emission is shifted to the ~ 25 kHz frequency of strong VLF transmitter signals. Interaction with the strong coherent transmitter wave leads to the precipitation of resonant electrons and a reduction of the hot/cold electron density ratio, leading to the suppression of nonlinear chorus development. In the absence of chorus rising tones, the local acceleration of electrons to MeV energies is shut down at the outer edge of the VLF bubble creating an apparent impenetrable barrier to their earthward extent. Because of the very slow cross-L diffusion for MeV electrons inside $L = 3$ (e.g. Baker et al. (2014); Ozeke et al (2018)), the sharp earthward flux gradients formed outside the region of chorus suppression can persist for days to weeks without extending significantly earthward of $L \sim 2.8$.

7 Summary

For conditions when the plasmasphere is contracted inside $L \sim 2.8$, two factors combine to suppress local acceleration of 100s keV seed electrons to relativistic energies at $L < 3$. (1) The ratio of f_{pe}/f_{ce} decreases to values < 2 in the low-density region outside the plasmopause where the magnetic field strength increases rapidly with decreasing L . (2) Interaction of resonant electrons with the VLF transmitter waves decreases the hot 10s - 100 keV electron population (n_h) needed to support nonlinear chorus wave growth. The calculations of Figure 5 indicate that, for the circumstances observed, a suppression of the nonlinear growth of chorus waves near $L \sim 2.8$ would result both from the localized decrease in f_{pe}/f_{ce} and the reduction of the density of the gyro-resonant electron population. In the absence of chorus rising tones, the rapid nonlinear acceleration of electrons to MeV energies earthward of $L \sim 3$ is suppressed. The result is an abrupt, sharp inner edge of the newly created MeV electron population and the appearance of an impenetrable barrier at the fixed location where the strongest transmitter signals near 25 kHz cross $0.5 f_{ce}$ at the outer extent of the VLF bubble at $L \sim 2.8$.

Acknowledgments and Data

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- Van Allen Probes observations used in this study can be obtained through instrument websites (EMFISIS wave data: <http://emfisis.physics.uiowa.edu>; MagEIS and REPT particle data: https://rbsp-ect.lanl.gov/rbsp_ect.php; EFW data : <http://www.space.umn.edu/rbspew-data/>). The values of the optimum and threshold amplitudes in Figures 5a and 5b are calculated by Equations (29) and (35) in Omura et al. (2015b). The nonlinear growth rates are calculated by Equation (22) in Omura et al. (2009), and the nonlinear growth rates are obtained by KUPDAP (Sugiyama et al., 2015; <http://space.rish.kyoto-u.ac.jp/software/>).

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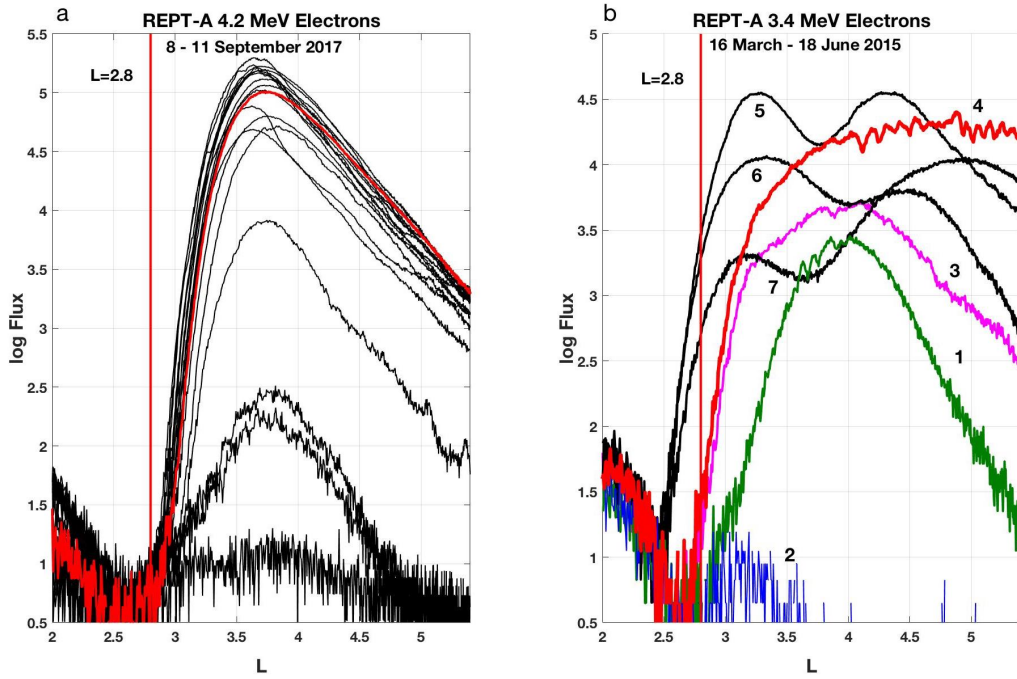


Figure 1. a) The pass by pass evolution of 3.4 MeV electron flux beginning 08 September 2017 demonstrates the rapid recovery of highly-relativistic electrons by > 4 orders of magnitude inside $L \sim 4$. Steep gradients were formed outside $L = 2.8$. The red curve highlights the pass at 04:40 UT on 09 September 2017, immediately following the data discussed in Figure 3 (below). b) The initially sharp multi-MeV electron flux boundaries that formed outside $L = 2.8$ during the 17-18 March 2015 storm gradually diffused earthward across the barrier to $L \sim 2.5$ over the ensuing months. Flux profiles 1-4 show the evolution of 4.2 MeV fluxes over 24-hour intervals from March 16th (1, pre-storm) to March 19th (4). Following storm onset, outer zone fluxes were decreased by $> 1000\times$ (2) at $L = 4$. As the outer zone fluxes recovered (3,4), sharp gradients were formed outside $L = 2.8$. Profiles 5 (18 April), 6 (18 May), 7 (18 June) show the evolution of the residual "storage belt" fluxes over the ensuing 3 months.

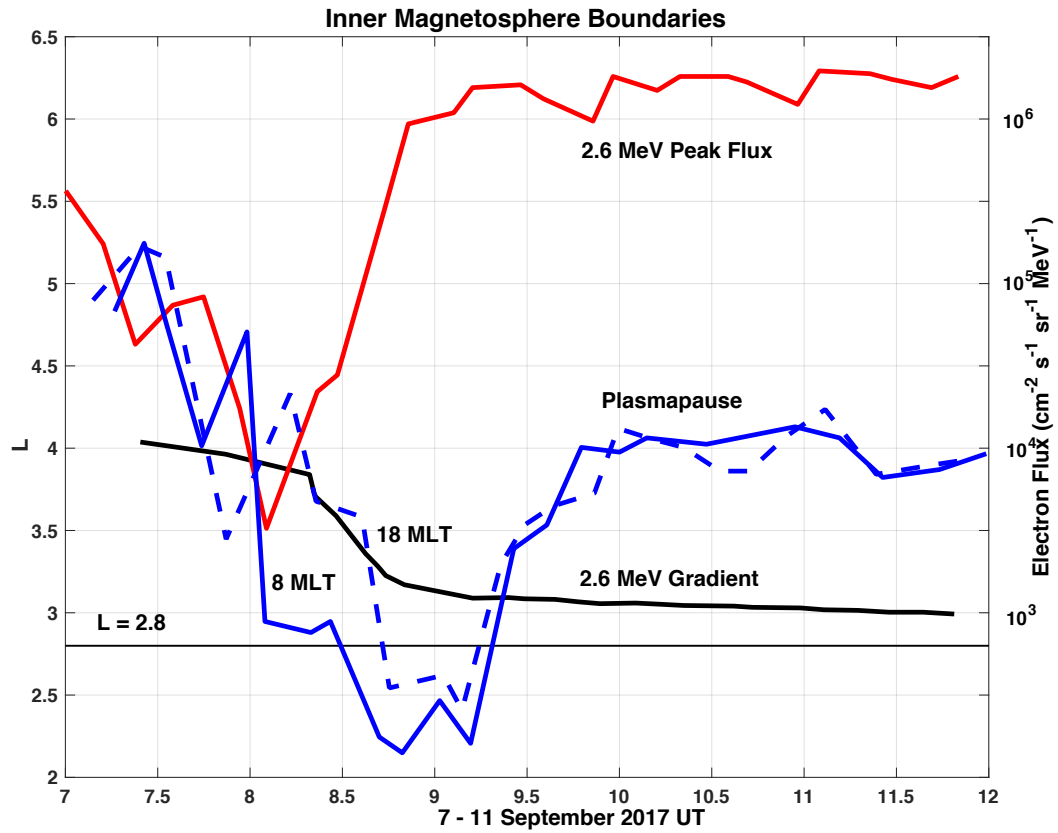


Figure 2. The peak REPT 2.6 MeV electron flux between $L = 3$ and $L = 5$ is indicated by the red curve. The plasmopause (blue curves; electron density $< 100 \text{ cm}^{-3}$) was earthward of $L = 3$ during the entire interval of MeV electron recovery. The midpoint of the sharp earthward gradient of 2.6 MeV fluxes (black dashed curve, log flux < 4.3) moved inward to $L \sim 3.3$ during the flux recovery. As seen in Figure 1a, MeV electrons were completely excluded inside $L \sim 2.8$.

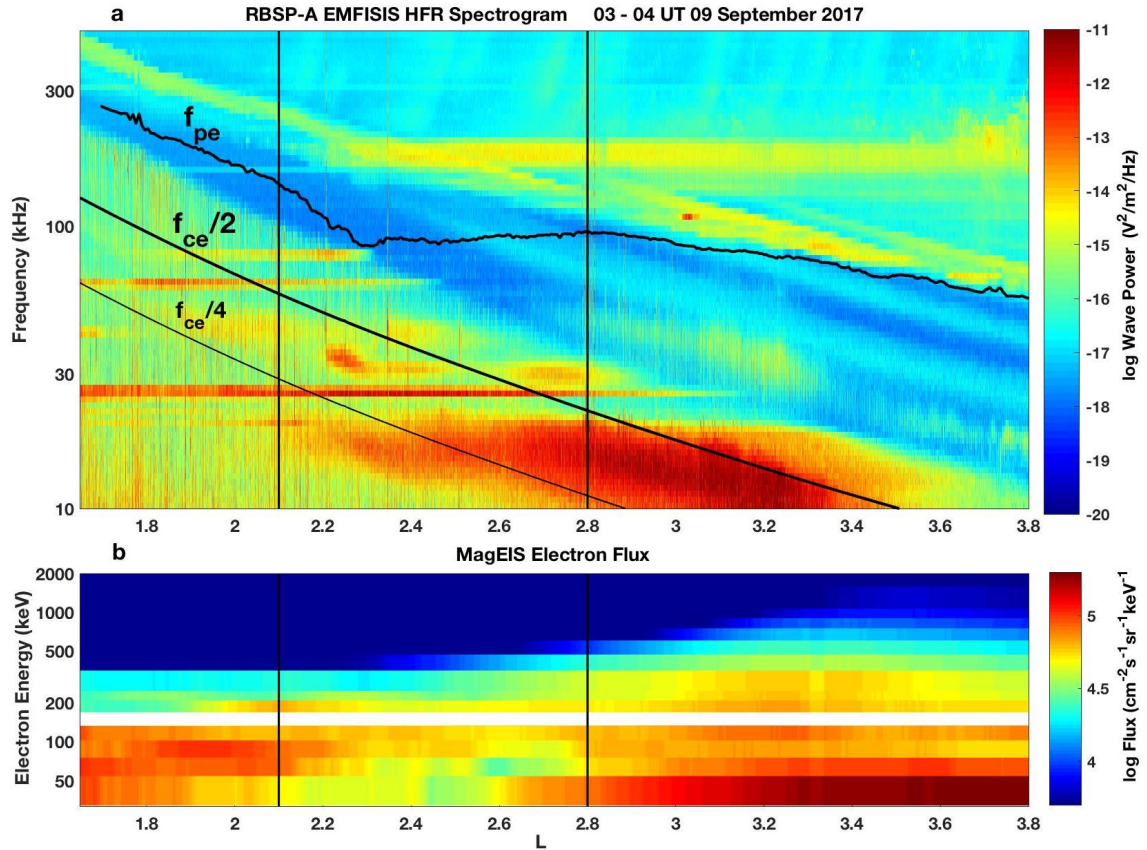


Figure 3. (a) EMFISIS HFR electric field spectrogram for the inbound pass of RBSP-A at 03:00 UT on 09 September 2017. Black vertical lines indicate $L = 2.8$ and $L = 2.1$ (the position of the plasmopause at this time), and heavy black curves show the electron plasma frequency (f_{pe}) and $1/2 f_{ce}$. Natural whistler mode emissions are suppressed by $\sim 1000\times$ in a band from 20-24 kHz immediately below the strong VLF transmitter signal at 24.9 kHz. (b) Low energy (< 100 keV) MagEIS electron fluxes, measured on a preceding pass of RBSP-B, were depleted between $L=2.1$ and $L=2.8$. Electron energies in this range are in cyclotron resonance ($n = 1$) at low pitch angles with VLF waves at the 24.9 kHz transmitter frequency.

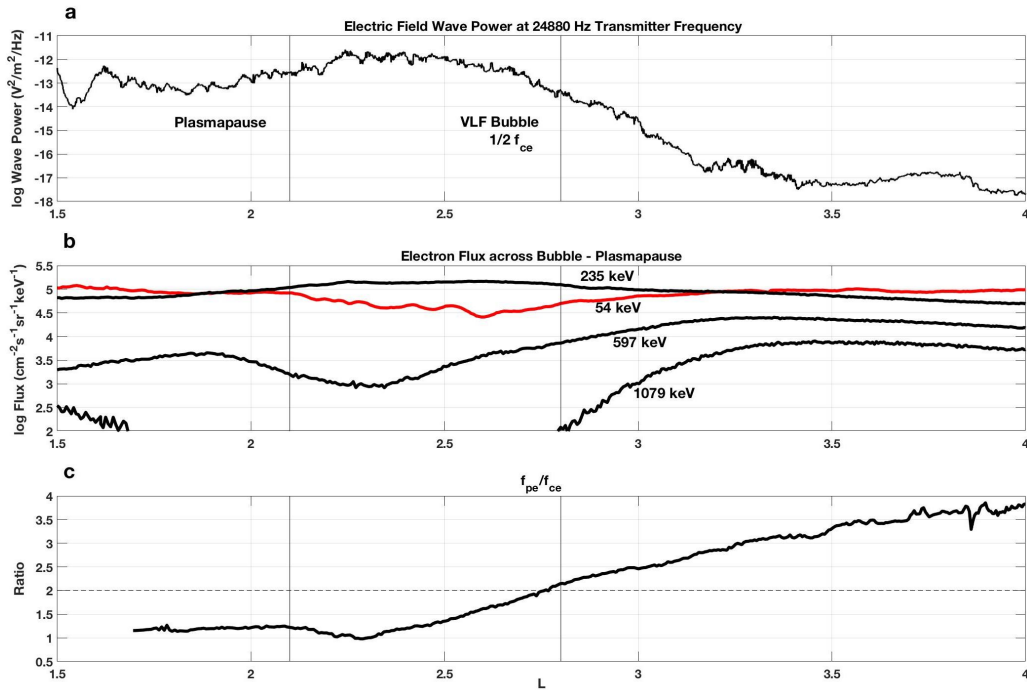


Figure 4. (a) Electric field wave power at the 24.9 kHz transmitter frequency increased by 10x outside the plasmopause. Black vertical lines indicate $L = 2.1$ and $L = 2.8$. (b) For the pass at $\sim 03:00$ on 09 September 2017, the MagEIS 235 keV electron flux (the lowest energy available from RBSP-A at the time) is enhanced between $L = 2.8$ and $L = 2.1$ across the region of the transmitter wave enhancement. 54 keV electron flux observed by RBSP-B ~ 3 hr earlier (red curve, cf. Figure 3b) decreased by ~ 3 x across the region of transmitter wave enhancement. (c) For the event shown in Figure 3, the observed ratio of the plasma frequency to the cyclotron frequency varied from 1 to 3 across the region of transmitter wave enhancement.

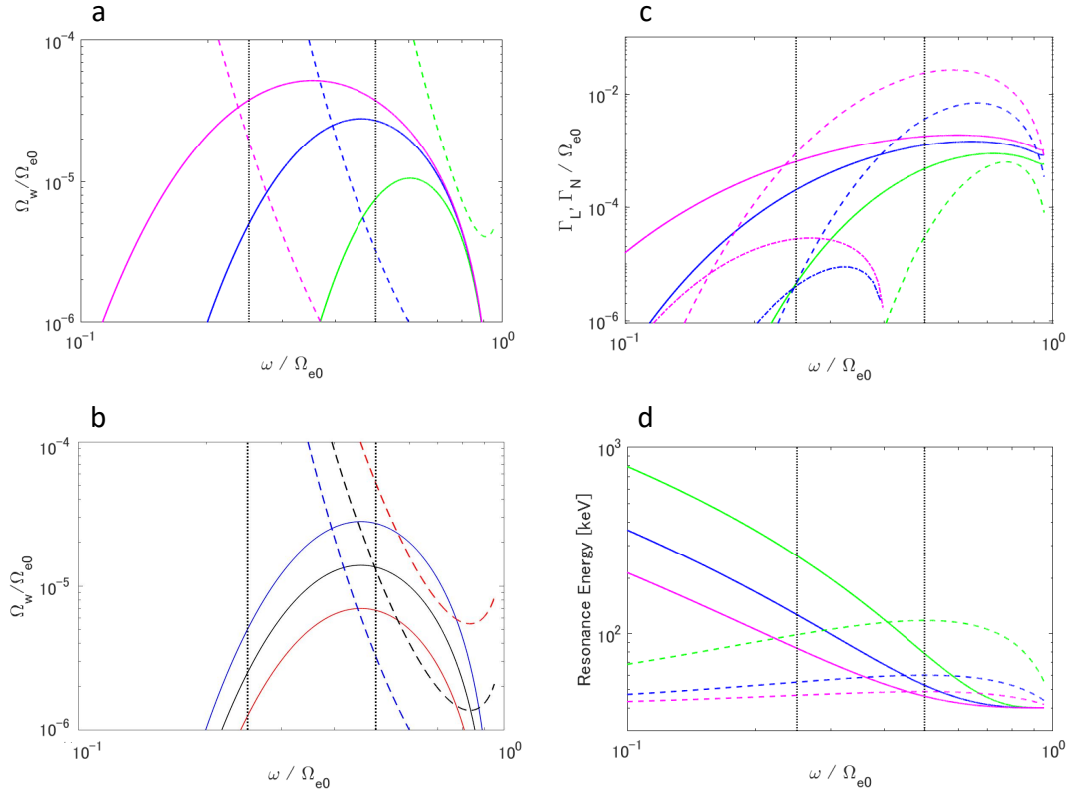


Figure 5. (a) Optimum amplitudes (solid line) and threshold amplitudes (dashed line) for different ratios $f_{pe}/f_{ce} = [3, 2, 1]$. (magenta, blue, green) with $nh/nc = 4 \times 10^{-4}$. (b) Optimum amplitudes (solid line) and threshold amplitudes (dashed line) for different $nh/nc = [4, 2, 1] \times 10^{-4}$ (blue, black, red) with $f_{pe}/f_{ce} = 2$. (c) Linear growth rates (dash-dot line), nonlinear growth rates at threshold amplitudes (dashed line), and nonlinear growth rates at optimum amplitudes (solid line) for different $f_{pe}/f_{ce} = [3, 2, 1]$. (magenta, blue, green) with $nh/nc = 4 \times 10^{-4}$. (d) Resonance energies of cyclotron resonance (solid line) and Landau resonance at quasi-parallel propagation (dashed line) for different $f_{pe}/f_{ce} = [3, 2, 1]$. (magenta, blue, green) with $nh/nc = 4 \times 10^{-4}$.