

1 **Heterogeneous Land-Surface Effects on TKE and Cloud**
2 **Formation: Statistical Insights from LES Cases**

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6 **Key Points:**

- 7 • Analysis of 92 LES cases shows strong statistical correlations between land-surface
8 heterogeneity and mesoscale atmospheric development
9 • Correlation between surface heterogeneity and circulation/cloud production is dom-
10 inated by the largest wavelengths in the land-surface field

Abstract

To aid development of sub-grid scale (SGS) parameterizations for Earth system models which consider heterogeneity in land-surface fields and land-atmosphere coupling, results from large-eddy simulations of 92 shallow convection cases over the Southern Great Plains are presented and analyzed. Each case is simulated with heterogeneous surface fields obtained from an offline field-scale land-surface model, and with spatially homogeneous surface fields with the same domain-wide mean value. By comparing corresponding heterogeneous and homogeneous cases, it is found that turbulent kinetic energy and liquid water path has a high correlation with the spatial variance of the surface heat flux fields. By further comparing the source of this correlation over the range of wavelengths in the surface fields, it is found that the majority of the heterogeneous land-atmosphere coupling is contained in wavelengths of order 10 km and larger, suggesting an encouraging degree of feasibility of including land-surface heterogeneity in global-scale SGS parameterizations.

Plain Language Summary

To help efforts to alleviate some of the issues associated with the relatively low-resolution grids used by modern global weather and climate models, we first created a dataset of 92 high-resolution simulations over the Southern Great Plains region of Oklahoma. All of the cases in the dataset are based on days which were observed to produce shallow clouds, which can have a significant impact on the incoming solar radiation. The high-resolution simulations were designed to cover a region large enough to contain relevant cloud production which is also too small to be represented on a modern global model. The dataset of high-resolution simulations is analyzed to compare the strength of the patterns in the land surface to the associated increase in cloud production. It is hoped that this and similar future studies will provide insights which increase the fidelity of cloud production models which intend to capture effects which are smaller than the grid used for global models.

1 Introduction

Modern coupled Earth system models (ESMs) are run at horizontal resolutions that are $\mathcal{O}(10 - 100 \text{ km})$, which is decided by the balance between computational resources and the demands of the atmospheric component of the coupled model, while the land-surface model (LSM) component could conceivably have an effective horizontal spatial resolution around $\mathcal{O}(10 - 100 \text{ m})$ (e.g., Chaney et al., 2018). This loss of land-surface information is made more significant by the fact that it spans the relevant length scales for many important coupled processes, namely those related to boundary-layer growth and cloud production (Bertoldi et al., 2013; Kang & Bryan, 2011; Ntelekos et al., 2008; Weaver, 2004).

The parameterization associated with sub-grid scale (SGS) cumulus production is very important in contemporary ESMs, by virtue of the importance of cloud production to the Earth system in general. Many modeling and observational studies find that secondary circulations induced by thermal surface heterogeneity can act as sources of convection and significantly alter local cloud production rates and distribution (e.g., Albertson et al., 2001; Dixon et al., 2013; Kang, 2020; Marsham et al., 2008; Mendes & Prevedello, 2020; Taylor et al., 2011; Phillips & Klein, 2014).

While the aforementioned land-surface patterns are SGS on grids used for most modern global models, there is a large amount of information available regarding the characteristics of the land-surface which could potentially be utilized by SGS parameterizations. Towards this effort, we present a large-eddy simulation (LES) study of 92 shallow convection cases over the Southern Great Plains (SGP) site,

60 based on cases developed by the LES ARM Symbiotic Simulation and Observation
61 Workflow (LASSO) campaign (W. Gustafson et al., 2019; W. I. Gustafson et al.,
62 2020). The cases are run using high-resolution spatially-heterogeneous land-surface
63 fields and also using spatially-homogeneous land-surface fields, which match the het-
64 erogeneous cases' domain-wide mean values through time but contain none of the
65 spatial structure.

66 We find that there is a strong correlation between basic metrics of heterogene-
67 ity in the surface heat flux fields and the resultant additional production of liquid
68 water path (LWP) and circulating kinetic energy. We also find that, for the cases
69 considered here, the majority of the relevant information about the heterogeneity of
70 the land-surface is contained in the few Fourier modes of the fields with the largest
71 wavelengths, which is encouraging from the perspective of computational resources
72 potentially required to consider SGS land-surface features.

73 **2 Model description**

74 Large-eddy simulations are conducted using version 3.8.1 of the WRF
75 model (Skamarock et al., 2008) with modifications as described by J. S. Simon
76 et al. (2021). Cases here use a horizontal resolution of 250 m and a domain of
77 130×130 km² laterally. The land-surface fields in the outer 15 km of the domain
78 are tapered to linearly approach their domain-wide mean on each boundary to elim-
79 inate discontinuities in the land-surface that may otherwise be introduced by the
80 periodic boundary conditions. Each domain is also rotated to closer align the bulk
81 liquid-water flux normally to the boundaries, based on results from an initial simu-
82 lation using the unrotated land-surface, to limit artificial spreading of liquid water
83 caused by the fluxes through the boundaries not aligning with the periodicity of
84 the domain. The model configuration is otherwise the same as in J. S. Simon et al.
85 (2021).

86 Each case is run with heterogeneous and homogeneous land-surface fields (sen-
87 sible heat flux, latent heat flux, skin temperature, albedo, and momentum drag co-
88 efficient), where homogeneous cases specify a uniform (in space) surface of each field
89 to match the time-evolving domain-wide mean of the corresponding heterogeneous
90 case. There is no feedback from the atmosphere to the land surface in the LES; the
91 HydroBlocks LSM is run offline and the output surface fields are specified as the
92 bottom boundary in the WRF model. Further details of the HydroBlocks LSM and
93 its coupling to the WRF model can be found in the Supporting Information.

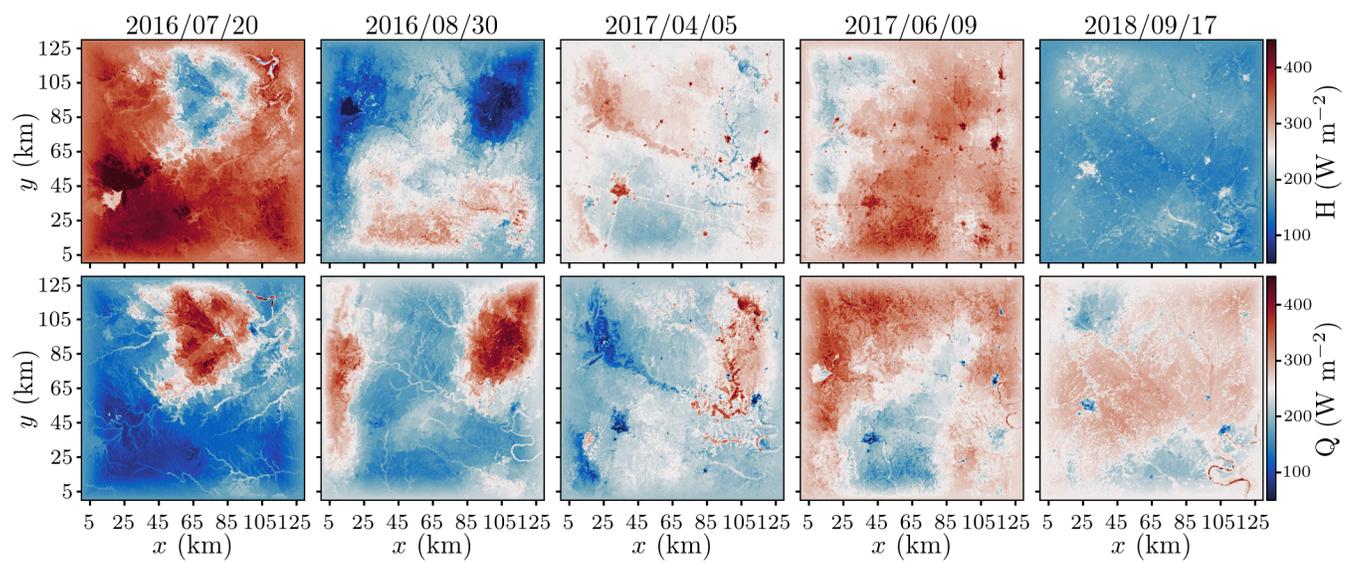


Figure 1. Example maps of sensible (H) and latent (Q) heat flux fields.

3 Results

3.1 Evaluation Metrics

The domain-wide measure of vertically-integrated, mass-coupled turbulent kinetic energy (TKE) is compared between cases, serving as a metric for general activity in ABL development. For brevity, “TKE” will refer to the vertically-integrated, mass-coupled form unless otherwise stated. On the discretized WRF grid, the TKE is found as

$$\text{TKE} = \sum_z \rho_a \left[\frac{1}{2} (u'^2 + v'^2 + w'^2) \right] \Delta_z, \quad (1)$$

where ρ_a is air density, (u, v, w) are the velocity components in the (x, y, z) directions, Δ_z is the grid spacing in the vertical direction, and a primed variable indicates deviation from the mean value in the (x, y) plane. For illustration, a time series of TKE for heterogeneous and homogeneous simulations of an example case is shown in Fig. 2a. Cases are also compared by their domain-wide LWP signal, which serves as a proxy for overall cloud production. On the discretized WRF grid, our measure of LWP is found as

$$\text{LWP} = \sum_z \rho_a q_l \Delta_z, \quad (2)$$

where q_l is liquid water mixing ratio.

Part of the LASSO modification to the WRF code is the addition of output solution fields as average values over a given interval of time, in addition to the standard instantaneous output fields. Here, time-averaged fields are found over 10 min intervals from samples taken internally every 30 s. Notationally, we will use $\mu(\phi)$ and $\sigma(\phi)$ to indicate the spatial mean and standard deviation, respectively, of a field $\phi = \phi(x, y)$ at a point in time. For temporal averages, we will use the notation $\text{mean}[\vartheta]$, found as

$$\text{mean}[\vartheta] = \frac{\sum_t g_s \vartheta}{\sum_t g_s}, \quad (3)$$

where $\vartheta = \vartheta(t)$ is a domain-wide scalar, with $g_s = g_s(t)$ defined as

$$g_s(t) = \begin{cases} 1 & : s(t) > 0.05 \max(s), \\ 0 & : s(t) \leq 0.05 \max(s), \end{cases} \quad (4)$$

where $s(t)$ is the surface downward clear-sky shortwave radiation at time t , and $\max(s)$ is the maximum value of s over the given simulation. The averaging procedure in (3) is used for both the atmospheric fields (Fig. 2c) and the land-surface statistics (Fig. 2e).

The heterogeneous vs. homogeneous statistics for TKE and LWP are compared using the metric $\gamma(\vartheta)$, defined as

$$\gamma(\vartheta) = \text{mean} [\log \gamma_t(\vartheta)], \quad (5)$$

where

$$\gamma_t(\vartheta) = \frac{g_s \vartheta_{\text{heterogeneous}} + 1}{g_s \vartheta_{\text{homogeneous}} + 1}. \quad (6)$$

Equations (6) and (5) are demonstrated visually in Fig. 2c and d, respectively. The form of (6) is motivated as a ratio of ϑ between heterogeneous and homogeneous cases, which is weighted by g_s to isolate daytime values. The addition of 1 to both

104 terms is included to limit the influence of very small values which are effectively neg-
 105 ligible, as well as to avoid the edge cases of $\gamma_t = 0$ or $\gamma_t = \infty$ when $\vartheta_{\text{heterogeneous}} = 0$
 106 or $\vartheta_{\text{homogeneous}} = 0$, respectively. When $\vartheta \approx 0$ for both the heterogeneous and ho-
 107 mogeneous cases, $\gamma_t \approx 1$, indicating the two cases have approximately equal mea-
 108 sures of ϑ , as intended. In (6), TKE is given in units of kg s^{-2} and LWP in units of
 109 g m^{-2} .

In addition to $\gamma(\text{TKE})$ and $\gamma(\text{LWP})$, we also compare corresponding heteroge-
 neous and homogeneous cases by only the circulating portion of kinetic energy. This
 is found from the turbulence spectra of 10-minute averaged u and v fields, where
 only the energy from modes that are in the lowest 5 km of the atmosphere and
 longer than 10 km laterally are included. The ratio of circulating energy between the
 heterogeneous and homogeneous cases, which we denote χ , is found as similarly to
 the TKE and LWP fields, as

$$\chi = \text{mean} \left[\log \sqrt{\gamma_t(E_u)\gamma_t(E_v)} \right], \quad (7)$$

where

$$E_\varphi = E_\varphi(t) = \sum_{z < 5 \text{ km}} \left[\sum_{\ell > 10 \text{ km}} |\widehat{f}(\varphi')|^2 \right], \quad (8)$$

110 and $|\widehat{f}(\varphi')|$ is absolute value of the normalized two-dimensional discrete Fourier
 111 transform of φ' , ℓ is the component of the Fourier mode's wavelength, λ , in the
 112 direction aligned with φ (e.g., $\ell = \lambda_x$ for $\varphi = u$).

Two length-scale metrics, L_Δ and L_2 , are presented for the land-surface fields,
 based on their Fourier spectra (the relaxation to the mean value on the outer 15 km
 of the land-surface fields render their boundaries as effectively periodic). The L_Δ
 length scale gives the approximate scale of the largest coherent structures in the
 field, and is found as

$$L_\Delta(\phi) = \frac{\sum_\lambda \lambda \Delta_\lambda \sqrt{|\widehat{f}(\phi)|}}{\sum_\lambda \Delta_\lambda \sqrt{|\widehat{f}(\phi)|}}, \quad (9)$$

where $\phi = \phi(x, y)$ is a heterogeneous surface field, and Δ_λ is the difference between
 λ and the next (smaller) wavelength in the discrete spectrum. The L_2 length scale
 gives the approximate scale of the smallest coherent structures in the field, and is
 found as

$$[L_2(\phi)]^2 = \frac{\sum_\lambda \lambda^2 \sqrt{|\widehat{f}(\phi)|}}{\sum_\lambda \sqrt{|\widehat{f}(\phi)|}}. \quad (10)$$

113 Correlations between atmosphere and land-surface fields are evaluated by the
 114 Pearson (ρ_p) and Spearman (ρ_s) correlation coefficients, as implemented by Virtanen
 115 et al. (2020) (e.g., Fig. 2f).

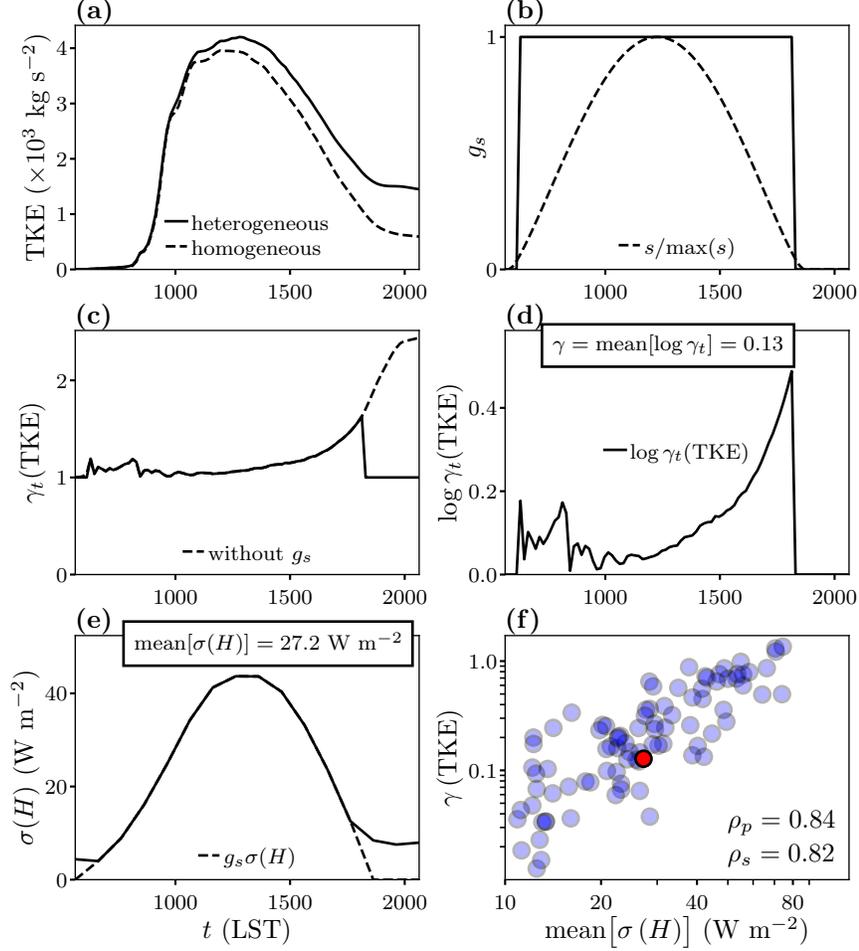


Figure 2. Demonstration of the comparison process for heterogeneous and homogeneous cases: (a) the domain-wide time series of TKE for the two simulations of 2017/08/30; (b) the time filter, g_s , used by the mean[ϑ] function as described by (4); (c) the calculation of γ_t , as described by (6); (d) the calculation of γ as described by (5); (e) the application of the time filter, g_s , and the mean[ϑ] function to the $\sigma(H)$ time series; (f) an example scatter plot of $\text{mean}[\sigma(H)]$ vs. $\gamma(\text{TKE})$ for all 92 days with the datapoint for 2017/08/30 shown in red.

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3.2 Land-Atmosphere Correlation

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The emergent secondary circulations driven by land-surface heterogeneity are analyzed by the correlation between χ and each of $\mu(H)$, $\sigma(H)$, $L_{\Delta}(H)$, and $L_2(H)$ (Figs. 3a–d, respectively). There is a strong positive relationship between χ and all of $\sigma(H)$, $L_{\Delta}(H)$, and $L_2(H)$, but only a trivial correlation with $\mu(H)$. The same presentation is repeated for $\gamma(\text{LWP})$ in Fig. 3e–h. The $\gamma(\text{LWP})$ data is very similar to that of the χ metric, but with $\sim 20\%$ smaller magnitudes. Of the 92 cases, 4 have more liquid water production in the homogeneous simulation, indicated by a negative value of $\gamma(\text{LWP})$; these datapoints are not shown in Fig. 3 but are included in the calculation of the correlation coefficients.

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Visually, the data for $\gamma(\text{LWP})$ compared to $\sigma(H)$, $L_{\Delta}(H)$, and $L_2(H)$ show a very similar pattern as χ but with a broader spread, suggesting from that LWP production is statistically driven similarly to circulation production but with additional considerations which are not captured by the land-surface heterogeneity, which is certainly in agreement with the physical perspective of ABL development. The same analysis considering $\gamma(\text{TKE})$, or using statistics from the latent heat flux or skin temperature fields gives very similar results, which is presented in the Supporting Information.

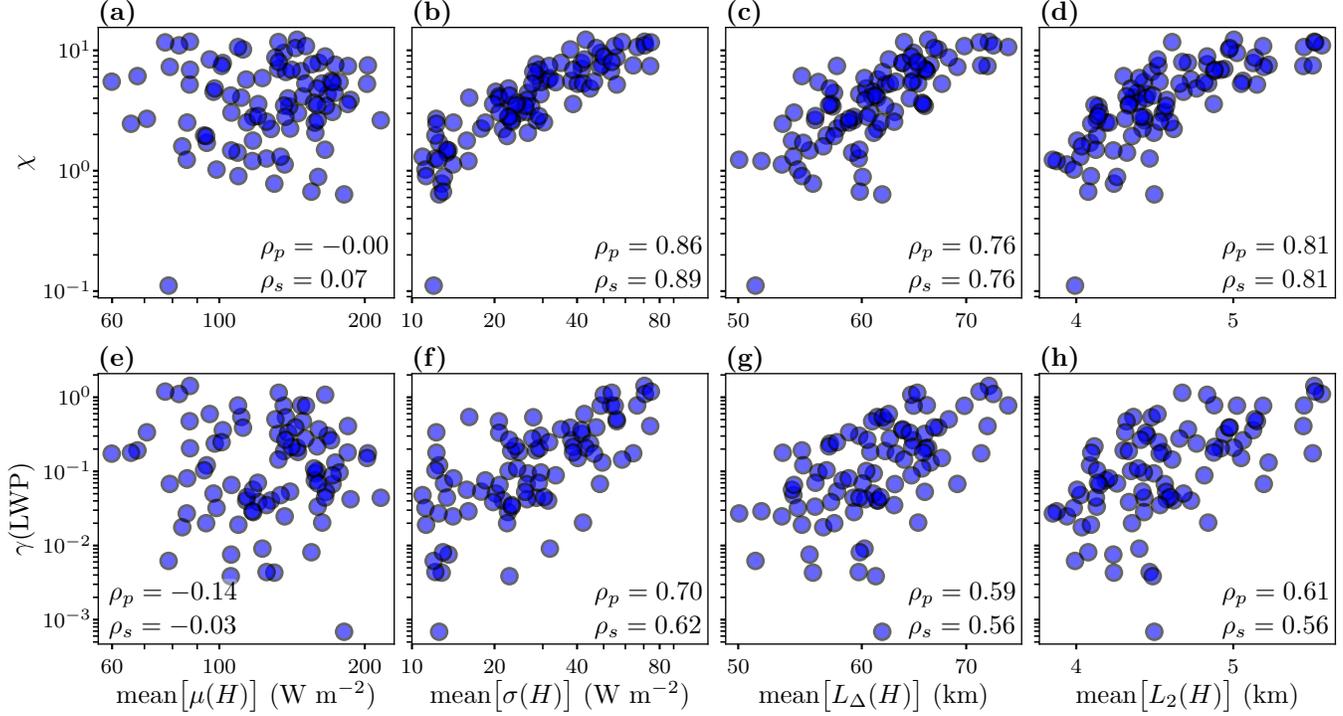


Figure 3. The χ (a – d) and $\gamma(\text{LWP})$ (e – h) metrics as functions of statistics of the surface sensible heat flux field, H . Four negative-valued data points for $\gamma(\text{LWP})$ with magnitudes $\mathcal{O}(10^{-4})$ are not shown, but are included in the calculation of ρ_p and ρ_s .

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3.3 Dominant Length Scales

The results presented in Sec. 3.2 bolster the motivation to include information about the underlying land-surface heterogeneity in global SGS boundary-layer parameterizations. The best methodology to make such considerations in either existing or new parameterization models is not immediately obvious, and potential solutions must add value on a level that is commensurate with their computational and implementation costs. To evaluate the relevance of the different ranges of length scales present in the land surface on the present dataset, the land surface fields are filtered over a range of length scales and compared. The filter, F , is applied in Fourier space as

$$F = 1 - \exp \left[-2\pi^2 \left(\frac{\Delta_{\text{filter}}}{\lambda} \right)^2 \right], \quad (11)$$

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where Δ_{filter} is the nominal filter length. To avoid ambiguity in the discussion, the operation of F is referred to as passing $\lambda < \Delta_{\text{filter}}$. An example of a mid-day sensible heat flux field from the dataset for different filter lengths is shown in the Supporting Information.

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For each filter length, the heterogeneous land-atmosphere coupling is reevaluated following the same procedure as in Sec. 3.2. The average value over the dataset of $\text{mean}[\sigma(H)]$ as a function of filter length is shown in Fig. 4a. Correlation coefficients for $\text{mean}[\sigma(H)]$ of the filtered dataset with $\gamma(\text{TKE})$, $\gamma(\text{LWP})$, and χ are shown as a function of filter length in Fig. 4b, c, and d, respectively.

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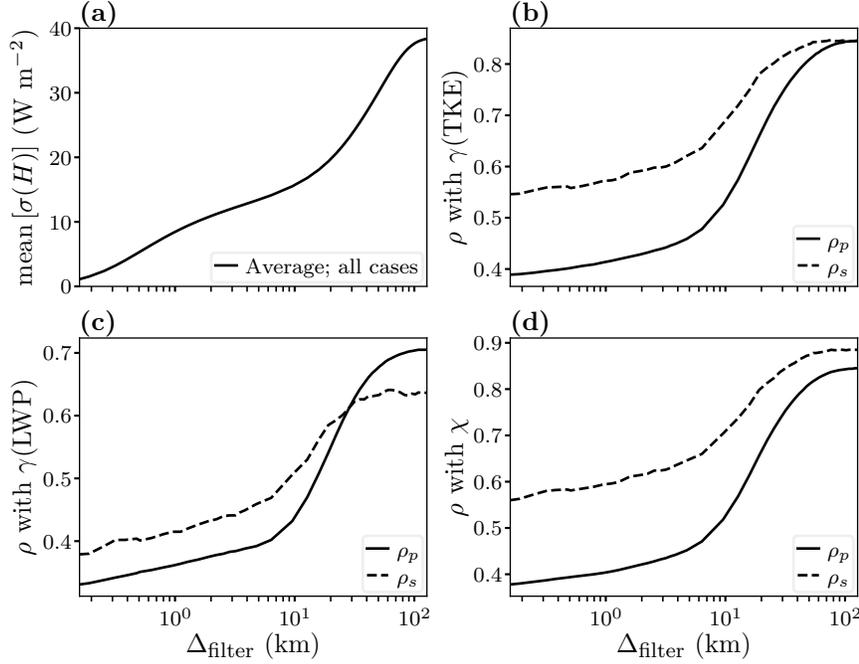


Figure 4. Average value of $\text{mean}[\sigma(H)]$ over the 92 days after filtering (a) and correlation coefficients of $\text{mean}[\sigma(H)]$ after filtering with $\gamma(\text{TKE})$ (b), $\gamma(\text{LWP})$ (c), and χ (d).

144 Figure 4a demonstrates that the majority of the standard deviation in the sensible
 145 heat flux field over the dataset is contained in length scales 10 km and larger.
 146 Figures 4b – d show the same concentration at length scales larger than 10 km
 147 for the Pearson and Spearman correlation coefficients between $\sigma(H)$ and $\gamma(\text{TKE})$,
 148 $\gamma(\text{LWP})$, χ . Because the wavelengths of Fourier modes grow geometrically, the re-
 149 sults seen in Fig. 4 suggest that the bulk of the correlation between $\sigma(H)$ and the
 150 atmospheric metrics is contained in the longest few modes.

151 4 Discussion and Conclusions

152 We have presented a statistical analysis of the TKE and cloud production
 153 caused by land-surface heterogeneity for 92 LES cases representing different summer
 154 days from 2015 – 2019 over the SGP site by comparing simulations using heteroge-
 155 neous and homogeneous land-surface fields. In Sec. 3.2 it is found that, despite all 92
 156 days having unique initial profiles and large-scale tendencies, there is a strong cor-
 157 relation between the production of circulating TKE (measured as the metric χ) over a
 158 diurnal cycle and land-surface heterogeneity. The correlation between cloud produc-
 159 tion, as measured by LWP, is $\sim 20\%$ smaller but is also significant. It is also seen in
 160 Sec. 3.3 that a large portion of the correlation between the atmosphere and hetero-
 161 geneous land-surfaces is concentrated in a relatively small number of the largest (by
 162 wavelength) modes in the land-surface fields.

163 The results in Sec. 3.2 demonstrate a strong, but incomplete, correlation be-
 164 tween heterogeneous land surface fluxes and secondary circulations. The land-surface
 165 heterogeneity is more strongly related to χ than LWP, which was expected: while
 166 TKE production does depend on the temperature and stability of the initial atmo-
 167 spheric profile, liquid water production is additionally constrained by condensation
 168 conditions. Still, the correlation coefficient values seen between $\gamma(\text{LWP})$ and $\sigma(H)$,

169 even without considerations for the state of the atmosphere, are quite strong with
170 $\rho_p = 0.70$ (Fig. 3f). The concentration of relevant land-surface heterogeneity in
171 structures with length scales of $\mathcal{O}(10 \text{ km})$ and larger seen in Sec. 3.3 is easily un-
172 derstood in the context of heterogeneous land-atmosphere coupling being largely
173 driven by emergent mesoscale circulations. That there is such a sharp increase in
174 correlation contained in the longest few modes of the land surface does have the
175 encouraging implication that the level of detail necessary for the successful devel-
176 opment of global-scale SGS parameterizations of heterogeneous land-atmosphere
177 coupling may not be overwhelming.

178 While a large amount of additional work is necessary before the realization
179 of an effective parameterization, the results seen here are encouraging. The most
180 immediate future work is a detailed analysis of the relationship between initial and
181 large-scale atmospheric conditions and land-surface heterogeneity on the atmospheric
182 response. The necessary increase in cases to realize such an experiment would also
183 enable the use of more sophisticated methods for analysis, perhaps eventually in-
184 cluding machine learning, which itself has the potential to provide a huge value to
185 parameterization development efforts.

186 **Acknowledgments**

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188 **Data Availability**

189 Simulations here use a modification of WRF version 3.8.1 developed and main-
190 tained by the LASSO team. The base WRF code, initial sounding files, and large-
191 scale forcing files are available from W. Gustafson et al. (2019). Additional modi-
192 fications to the WRF code to specify heterogeneous surfaces, data files for surface
193 fields for each simulation, and model control files for each simulation are available
194 at J. Simon et al. (2023a). Relevant model output and scripts used for analysis and
195 plotting are available J. Simon et al. (2023b).

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