

1                    **Process Modeling of Aerosol-cloud Interaction in**  
2                    **Summertime Precipitating Shallow Cumulus over the**  
3                    **Western North Atlantic**

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23                   **Key Points:**

- 24                   • Aerosol-cloud interactions in precipitating shallow cumuli are investigated using  
25                   large-eddy simulations (LES) and observations  
26                   • LES show that aerosol-induced cloud water adjustment is dominated by precip-  
27                   itation and is anticorrelated with cloud-top entrainment  
28                   • A decrease in cloud fraction in response to aerosol increase is shown in the pre-  
29                   cipitating cumuli

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**Abstract**

Process modeling of aerosol-cloud interaction is essential to bridging gaps between observational analysis and climate modeling of aerosol effects in the Earth system and eventually reducing climate projection uncertainties. In this study, we examine aerosol-cloud interaction in summertime precipitating shallow cumuli observed during the Aerosol Cloud meTeorology Interactions oVer the western ATlantic Experiment (ACTIVATE). Aerosols and precipitating shallow cumuli were extensively observed with in-situ and remote-sensing instruments during two research flight cases on 02 June and 07 June, respectively, during the ACTIVATE summer 2021 deployment phase. We perform observational analysis and large-eddy simulation (LES) of aerosol effect on precipitating cumulus in these two cases. Given the measured aerosol size distributions and meteorological conditions, LES is able to reproduce the observed cloud properties by aircraft such as liquid water content (LWC), cloud droplet number concentration ( $N_c$ ) and effective radius  $r_{\text{eff}}$ . However, it produces smaller liquid water path (LWP) and larger  $N_c$  compared to the satellite retrievals. Both 02 and 07 June cases are over warm waters of the Gulf Stream and have a cloud top height over 3 km, but the 07 June case is more polluted and has larger LWC. We find that the aerosol-induced LWP adjustment is dominated by precipitation and is anticorrelated with cloud-top entrainment for both cases. A negative cloud fraction adjustment due to an increase of aerosol number concentration is also shown in the simulations.

**Plain Language Summary**

Aerosol-cloud-interaction (ACI) regulates the energy budget of the Earth and poses the largest uncertainty in climate projection. Particularly, ACI of low clouds is poorly understood and causes the spread of Earth System Models (ESMs) in predicting cloud and climate responses to aerosol changes. Process studies have shown a nonlinear cloud water amount and cloud fraction adjustments due to aerosol changes via precipitation and cloud-top entrainment, which are not often captured correctly in ESMs. This study explores the physical mechanisms of ACI in marine low clouds with a focus on precipitating low clouds using a cloud process model and unprecedented field campaign measurements of meteorology states, cloud properties, and aerosols collected during the Aerosol Cloud meTeorology Interactions oVer the western ATlantic Experiment (ACTIVATE). We show that the aerosol-induced cloud water amount adjustment is dominated by changes in precipitation and is negatively correlated with cloud-top entrainment. Our findings can help improve the representation of ACI within precipitating marine low clouds in ESMs.

## 1 Introduction

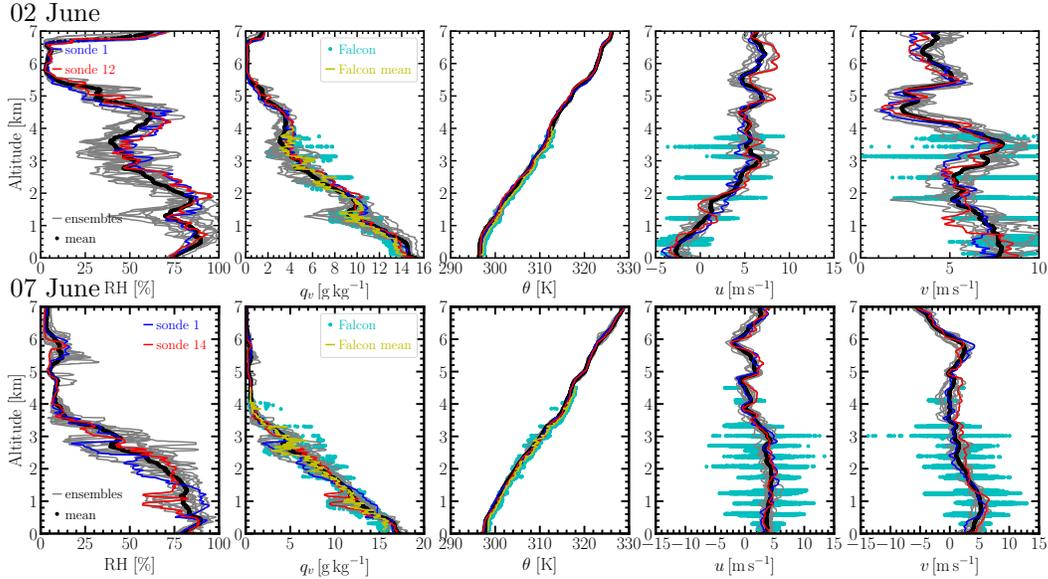
Aerosol-cloud interaction (ACI) poses the largest uncertainty for accurate climate projection (Seinfeld et al., 2016). Assuming fixed liquid water path (LWP) and Cloud Fractional Coverage (CFC) increasing the aerosol number concentration  $N_a$  leads to a larger cloud droplet number concentration  $N_c$ , smaller effective radius  $r_{\text{eff}}$ , and stronger outgoing shortwave radiation (Twomey, 1977). The enhanced shortwave cloud radiative effect can affect the meteorological state and boundary layer structure of clouds (Li et al., 2023). Decreased  $r_{\text{eff}}$  suppresses the precipitation rate by inhibiting the collision-coalescence processes, resulting in a higher liquid water path (LWP) and possibly larger cloud fraction (Albrecht, 1989). This is the conventional wisdom of ACI involving multi-scale, non-linear processes from aerosol/cloud microphysics to large-scale atmospheric circulations, which are fundamentally poorly understood due to the intractable scale range and strong nonlinearity, and therefore, mathematically poorly represented in Earth System Models (ESMs) (Seinfeld et al., 2016; Smith et al., 2020). Besides the strongly nonlinear nature of ACI, it is also in an emergent and non-equilibrium state, which hinders our understanding of processes that determines sinks and sources of aerosol and cloud. One example is the nonlinear interplay among  $N_a$ ,  $N_c$ , and the sink terms of liquid water content (LWC) for summertime shallow cumuli (Seinfeld et al., 2016). The parameterization of these clouds is responsible for the spread of ESMs in estimating the equilibrium climate sensitivity (ECS) (Zhao et al., 2016). In addition, even though ACI of shallow cumuli was shown to affect the ECS (Gettelman et al., 2019), the magnitude is uncertain. Droplet evaporation and precipitation are the two major sinks of LWC. Larger  $N_a$  results in larger  $N_c$  but smaller cloud droplets, which are easier to evaporate.  $N_a$ -induced suppression of the precipitation rate leads to a positive LWP response (Glassmeier et al., 2021). Precipitation removes aerosols from clouds and therefore leads to a negative feedback on  $N_c$  (Radke et al., 1980). This LWP feedback is also influenced by aerosol hygroscopicity  $\kappa$ , which determines the activation rate of aerosols acting as cloud condensation nuclei (CCN) to form cloud droplets.  $\kappa$  poses a great challenge for ESMs due to uncertainties in the detailed composition of particles, aerosol mixing state, and potential nonlinear interactions between them (Petters & Kreidenweis, 2007).

The aforementioned aerosol-cloud-precipitation interaction is very challenging to simulate even using large-eddy simulations (LES), where the relevant large-size turbulent eddies for cloud formation are resolved but droplet-turbulence interactions are ignored. The latter (Li et al., 2020) enhances the collision-coalescence process that determines the precipitation rate. Ackerman et al. (2004) showed that LWP response to aerosol-induced precipitation suppression depends on the competition between moistening due to decreased surface precipitation and drying due to enhanced cloud-top entrainment. Therefore, the LWP adjustment in response to increasing  $N_a$  can be divided into entrainment-dominated and precipitation-dominated regimes. For non-precipitating clouds (typically having the appearance of closed cells), increasing  $N_a$  leads to more abundant, smaller cloud droplets that can evaporate more readily due to entrainment drying because smaller droplets provide a larger surface area for a fixed amount of LWP (Ackerman et al., 2004). This entrainment drying process leads to a negative LWP adjustment to increasing  $N_a$ , indicating less reflective clouds and therefore, a weaker cooling effect. For precipitating clouds (typically having the appearance of open cells), more abundant but smaller cloud droplets increase colloidal stability through the suppression of precipitation rate, and thus yield a larger value of LWP. This positive LWP adjustment to increasing  $N_a$  indicates thicker and more reflective clouds, i.e., a stronger shortwave radiative cooling effect (Albrecht, 1989). Satellite observations have suggested complex LWP adjustments. Gryspeerd et al. (2019) showed a negative LWP adjustment in the majority of the oceanic regions using satellite retrievals, indicating that LWP reductions due to ACI could offset a significant fraction of the indirect aerosol radiative effect related to albedo increase. Diamond et al. (2020) reported significant cloud brightening due to increased  $N_a$  from ship emissions in subtropical low clouds, which is refuted by Glassmeier et al. (2021), who

119 pointed out that the shipping-induced aerosol radiative cooling for non-precipitating strato-  
120 cumuli is overestimated by a factor of up to 200% because of the underestimated nega-  
121 tive LWP adjustment related to current estimates of the average lifetimes of ship tracks.  
122 However, by considering both visible (as in Glassmeier et al. (2021)) and invisible ship  
123 tracks, Manshausen et al. (2022) showed positive LWP adjustment and therefore, a larger  
124 aerosol cooling effect. Aerosol effects on LWP and CFC based on satellite measurements  
125 only use snapshots of aerosol-cloud fields and ignore the temporal nature of cloud ad-  
126 justments, which could lead to inaccurate estimation of aerosol effects (Bellouin et al.,  
127 2020). Recently, Arola et al. (2022) found that a positive LWP adjustment can be eas-  
128 ily misinterpreted as a negative adjustment based on satellite measurements due to satel-  
129 lite retrieval errors (Painemal & Zuidema, 2011) and the propagation and spatial vari-  
130 ability in aerosols and clouds that cannot be captured by satellite instruments. In ad-  
131 dition, Christensen et al. (2022) concluded that these results from natural experiments  
132 cannot be easily scaled to global scales. This is because only shallow clouds are consid-  
133 ered and the effect of emission on deeper clouds is omitted in natural experiments. Mod-  
134 eling studies (Wang et al., 2011; Possner et al., 2018) have shown that cloud brighten-  
135 ing and LWP adjustments in response to aerosol emissions from ships depend strongly  
136 on boundary-layer meteorological conditions and dynamical feedback induced by pre-  
137 cipitation change. This drives the need for an in-depth investigation of ACI in a more  
138 comprehensive meteorological context. The Aerosol Cloud meTeorology Interactions oVer  
139 the western ATlantic Experiment (ACTIVATE) field campaign (2020-2022) has been con-  
140 ducted to bridge such a gap.

141 Many studies focus on stratocumulus-to-cumulus cloud transitions, of which the  
142 physical drivers and feedbacks are still unclear (Sandu & Stevens, 2011). Wang & Fein-  
143 gold (2009) showed that precipitation change can drive the transition. Yamaguchi et al.  
144 (2017) and Wood et al. (2018) found fast transition ( $\sim 10$  h) because of the drizzle ini-  
145 tiation and depletion of aerosols by precipitation change using LES. A larger  $N_a$  elon-  
146 gates the timing of the transition even though it is modulated by the diurnal cycle and  
147 large-scale meteorology, as shown in the LES studies (Goren et al., 2019). Using satel-  
148 lite retrievals, Christensen et al. (2020) showed that aerosols enhance the lifetime of clouds  
149 and increase cloud fraction in stable atmospheric conditions during the stratocumulus-  
150 to-cumulus transition. Erfani et al. (2022) confirmed the delayed stratocumulus-to-cumulus-  
151 transition due to aerosol-cloud-precipitation interactions for initially clean MBL and Twomey  
152 effect for initially polluted MBL using LES with a prognostic aerosol model, where aerosol  
153 life cycle with sources and sinks of aerosols included.

154 In this study, we consider precipitating summertime cumuli observed during AC-  
155 TIVATE since they can rapidly form rain (Rauber et al., 2007) and are an ideal candi-  
156 date to study aerosol-cloud-precipitation interactions. In addition, the cloud fraction of  
157 these clouds are severely under-predicted (few percent) in the ESMs (Rémillard & Tse-  
158 lioudis, 2015; Sorooshian et al., 2019) compared to the satellite observations (15–20%)  
159 in the North Atlantic region. We investigate the ACI of summertime cumuli over the West-  
160 ern North Atlantic Ocean (WNAO) region using LES and measurements during the AC-  
161 TIVATE campaign. The ACTIVATE campaign aims to build unprecedented statistics  
162 to improve process-level understanding of ACI and their representation in ESMs (Sorooshian  
163 et al., 2019). To study aerosol-cloud-precipitation interactions, we select two contrast-  
164 ing cases from the ACTIVATE campaign. The first one is a clean case with heavy pre-  
165 cipitation. The second one is a polluted case with light drizzling conditions. Contrary  
166 to most previous process studies that focused on sensitivity tests of ACI by arbitrarily  
167 perturbing the  $N_a$  or  $N_c$  (Wang & Feingold, 2009; Chen et al., 2011; Yamaguchi et al.,  
168 2017; Goren et al., 2019), we utilize measured  $N_a$  and  $N_c$  from ACTIVATE to under-  
169 stand ACI and its impact on LWP and CFC adjustments.



**Figure 1.** Profiles from dropsonde and Falcon measurements up to 7 km (same as the LES vertical domain size) for the 02 June (upper row) and 07 June (lower row) 2021 cases. The blue and red curves represent the first and last dropsonde, respectively, released at about the same location but one hour apart. The gray lines represent dropsondes in between, and the thick black lines represent the corresponding mean profile. The cyan dots show all the data points from the Falcon measurement (up to  $\sim 4$  km) during the dropsonde measurement time. The yellow lines represent the averaged Falcon measurement every 10 m vertically to approximately match the vertical spacing of dropsonde profiles.

## 2 Data and methods

### 2.1 Observations and reanalysis data

#### 2.1.1 Two precipitating cases

We select two contrasting process-study cases during the ACTIVATE 2021 summer field campaign. The case on 02 June 2021 is a heavily precipitating case (see the satellite visible image in Figure S1(a)) with the highest rain rate of  $23 \text{ mm h}^{-1}$  (the FCDP sampling frequency of 1 Hz) while the one on 07 June 2021 (Figure S1(b)) is a drizzling case (up to  $5 \text{ mm h}^{-1}$ ). The mean precipitation rate over the dropsonde circle is not provided due to the sample issue of the two-dimensional stereo (2DS) probe during the flight. However, the Fast Cloud Droplet Probe (FCDP)- $r_{\text{eff}}$ , rain water path (RWP) produced by LES, and Falcon forward camera records support our categorization of the precipitation for these two cases. The ACTIVATE campaign employed a dual-aircraft strategy to provide spatially coordinated measurements of meteorology states, trace gases, aerosol, and cloud properties (Sorooshian et al., 2019, 2023). The high-flying ( $\sim 9$  km in altitude) King Air measures meteorology states using dropsondes (Li et al., 2022) as well as aerosol and cloud retrievals using remote sensing instruments. The low-flying Falcon conducts in-situ measurements of water vapor (Diskin et al., 2002), trace gases, aerosol, and cloud properties. Figure S2 and Figure S3 shows the vertical profiles of water vapor mixing ratio  $q_v$  at 12 dropsonde locations and the simultaneously measured  $N_c$  from the FCDP for the 02 and 07 June 2021 cases, respectively. The measurements took place between 18:29:20 to 19:46:16 UTC and 18:25:54 to 19:45:37 UTC for the 02 and 07 June 2021 cases, respectively.

### 192 **2.1.2 Measured aerosol size distribution**

193 A Scanning Mobility Particle Sizer (SMPS, TSI model 3085 differential mobility  
 194 analyzer and TSI model 3776 condensation particle counter, 1/60 Hz) and a Laser Aerosol  
 195 Spectrometer (LAS, TSI model 3340) were used to measure aerosol particles with diam-  
 196 eter  $d$  between 3 – 100 nm and larger than 100 nm below the cloud base, respectively.  
 197 Their uncertainty is within  $\pm 10 - 20\%$  over the submicron aerosol size range (Moore  
 198 et al., 2021). The measured aerosol size distributions are fitted with lognormal modes  
 199 as shown in Figure S4(a) for the 02 June 2021 case and in Figure S4(b) for the 07 June  
 200 2021 case. The corresponding fitted parameters are listed in Table S1. The vertical struc-  
 201 ture of  $N_a$  is derived using combined polarimetric and lidar remote sensing observations  
 202 (Schlosser et al., 2022). The retrieved vertical structure of  $N_a$  exhibits exponential decay  
 203 with height (Figure S5). Our LES takes the lognormal distributions as aerosol in-  
 204 put, which follows this exponential decay with height in the simulation domain.

### 205 **2.1.3 Estimated hygroscopicity**

206 The bulk hygroscopicity ( $\bar{\kappa}$ ) of aerosol particles for each lognormal size mode is es-  
 207 timated from  $\kappa$  and mass of each chemical component  $m_i$  following the volume mixing  
 208 rule (Petters & Kreidenweis, 2007). The  $m_i$  (listed in Table S2) is measured by an Aero-  
 209 dyne High Resolution Time-of-Flight Aerosol Mass Spectrometer (HR-ToF-AMS) (De-  
 210 Carlo et al., 2008) with an uncertainty up to 50%. The estimated bulk hygroscopicity  
 211  $\bar{\kappa}$  for aerosol particles larger than 60 nm in diameter  $d$  is listed in Table 1 and Table S3,  
 212 which is used for the second and third mode of the lognormal distribution. For aerosol  
 213 particles with  $d \leq 60$  nm (first mode of the lognormal distribution) that lack valid mea-  
 214 surements, we use the smallest value of the organic component  $\kappa = 0.014$  and the mean  
 215 value  $\kappa = 0.1$  for the 02 and 07 June 2021 cases, respectively. We adopt such treatment  
 216 of estimating  $\bar{\kappa}$  for two reasons. First, measuring mass fraction of aerosol particles with  
 217  $d \leq 60$  nm is very challenging with high uncertainties. Therefore, we use the estimated  
 218  $\kappa$  of the organic components from existing literature. Second, the smallest and mean  $\kappa$   
 219 value of the organic component as input yields the best matching cloud microphysical  
 220 properties to the in-situ measurements for the 02 June and 07 June cases, respectively.

### 221 **2.1.4 Measured cloud microphysical properties**

222 The cloud droplet size distribution,  $N_c$ , and  $r_{\text{eff}}$  and LWC were measured by FCDP.  
 223 The FCDP can measure cloud droplets with diameter ranging from 3–50  $\mu\text{m}$  with an  
 224 uncertainty of less than 20% (Baumgardner et al., 2017; Knop et al., 2021). Cloud par-  
 225 ticles larger than 50  $\mu\text{m}$  are measured by the 2DS probe (Lawson et al., 2006) with a spa-  
 226 tial resolution of 11.4  $\mu\text{m}/\text{pixel}$  (Voigt et al., 2010; Bansmer et al., 2018). The 2DS cov-  
 227 ers a size range of 28.5 – 1464.9  $\mu\text{m}$  in this study.

### 228 **2.1.5 Reanalysis and satellite data**

229 Since the idealized LES cannot capture the large-scale motions of the atmospheric  
 230 flow, we use the fifth generation of European Centre for Medium-Range Weather Fore-  
 231 casts’s Integrated Forecast System (ERA5) reanalysis (hourly model-level and single-level  
 232 with a mesh grid-size of 31 km) large-scale forcings (i.e., moisture and temperature ad-  
 233 vective tendencies and wind profiles) and surface heat fluxes to drive the LES (Li et al.,  
 234 2022, 2023). LWP retrieved from hourly single-level (quantities obtained from the model  
 235 level) ERA5 and the Modern-Era Retrospective analysis for Research and Applications  
 236 version 2 (MERRA-2) (starting from 00:30 UTC) is used for comparison with WRF-LES  
 237 results and observations. The mean ERA5 (MERRA-2) LWP is calculated by averag-  
 238 ing model grids over the dropsonde-covered area. Both ERA5 (hourly) and MERRA-  
 239 2 (3-hourly) provide the CFC field at individual model levels, from which the time evo-  
 240 lution of CFC is obtained by averaging the maximum values of the CFC vertical-profiles

obtained by sampling each layer conditionally with a threshold of  $LWC = 0.02 \text{ g cm}^{-3}$  for clouds below 7 km. Both LES and ERA5/MERRA-2 reanalysis results are compared to the GOES-16 product, the first of the GOES-R series of the Geostationary Operational Environmental Satellites (GOES). The GOES-16 cloud retrievals we use in this study have a pixel size of 2 km and a time interval of 20 minutes.

## 2.2 LES numerical experiment design

The Weather Research and Forecasting (WRF) model (Skamarock et al., 2019) in the idealized LES mode (WRF-LES), i.e., periodic boundary condition in horizontal directions (Wang et al., 2009), is used in this study. The LES domain has a lateral size of  $L_x = L_y = 20 \text{ km}$  with a grid spacing of  $dx = dy = 100 \text{ m}$  and a vertical extent of  $z_{\text{top}} = 7 \text{ km}$  with 153 vertical layers. Although our sensitivity tests with  $dx = dy = 300 \text{ m}$  produce deeper clouds (closer to the measurements) than the ones with  $dx = dy = 100 \text{ m}$  (Figure S8 and Figure S11 in the supplement), we use a 100 m horizontal grid spacing to resolve smaller turbulent eddies that are important for the formation and evolution of shallow cumuli. Time-varying, area-averaged temperature and moisture advective tendencies ( $\partial_t \bar{\theta}$  &  $\partial_t \bar{q}_v$ ), divergence ( $\bar{D}$ ), and surface turbulent heat fluxes are obtained from ERA5 for both cases except that the largest hourly surface heat fluxes among all the ERA5 grids within the dropsonde circle area are used for the 07 June 2021 case. A relaxation time scale of 3 hours is applied to nudge  $\theta$  and  $q_v$  above 3 km and 1 hour for  $u$  and  $v$  in the entire domain to ERA5 for the 02 June 2021 case. This nudging strategy produces the best matching meteorology state (Figure S6 and Figure S7) to dropsonde measurements and observed cloud properties (Figure S8) to the FCDP measurements. For the 07 June 2021 case, only the  $u$  and  $v$  profiles are nudged to ERA5 with a time scale of 1 hour above 400 m and a 200 m transition depth to best reproduce the observed cloud properties. We adopt the Eulerian forcing instead of the Lagrangian one (forcing derived following the Lagrangian trajectory of the air mass) because the former leads to more comparable clouds to the observations (see the comparison between them in Figure S12–S14 in the supplement). The CAM radiative transfer model and a constant sea surface albedo of 0.06 are used. The Coriolis force corresponding to the center location of model domain is applied to all simulations.

The two-moment Morrison cloud microphysics scheme (Morrison et al., 2009) with prescribed aerosol size modes (see section 2.1.2) and hygroscopicity (see section 2.1.3) is employed, as initially implemented by Endo et al. (2015). Simulations with prescribed aerosol size distributions derived from the ACTIVATE campaign measurements, as described in section 2.1.2, are performed for both cases. We use prescribed aerosol size distribution instead of the prognostic one as in Erfani et al. (2022) because a prognostic aerosol model requires accurate information about particle and gas emissions to reproduce the observed aerosol size distributions. All simulations start at 06:00 UTC and end at 21:00 UTC with a fixed time step of 1 s. Initial profiles of temperature, humidity, and winds for all simulations are obtained from the corresponding ERA5 profiles averaged over the targeted case domain at 06:00 UTC. We refer to Table 1 for the input  $N_a$ ,  $N_c$ , and  $\bar{\kappa}$  of simulations.

## 3 Results

### 3.1 Aerosol effect on heavily precipitating cumuli: 02 June 2021 case

The 02 June 2021 case is characterized by heavy precipitation. The meteorology state from the dropsonde measurements exhibits strong spatial variation of RH ( $q_v$ ) by comparing the grey curves (individual dropsondes) and black curve (mean profile) as shown in Figure 1(a). The strong spatial variation of  $q_v$  makes the simulation of this case challenging. The mean  $q_v$ -profile from the Falcon measurement (yellow curve) agrees with the dropsonde measurement. The instantaneous Falcon measurements (cyan dots) within

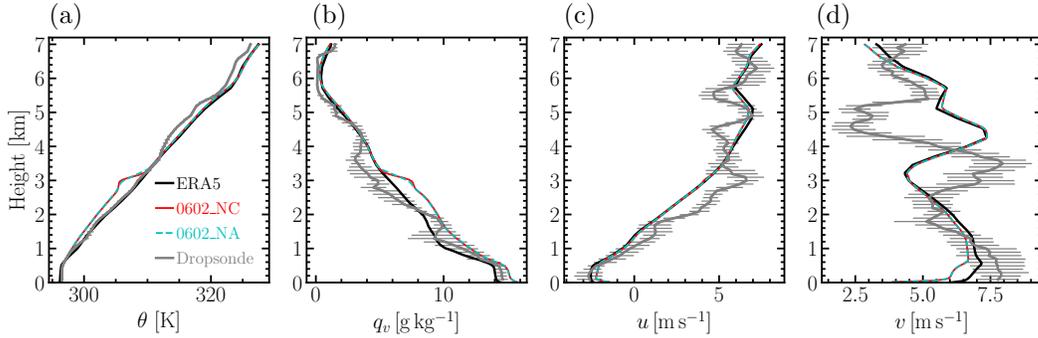
**Table 1.** List of simulations. “NC” denotes prescribed cloud droplet number concentration and “NA” denotes prescribed aerosol number concentration measured below cloud base.

Simulations	$N_a^{\text{input}} [\text{cm}^{-3}]$	$N_c^{\text{input}} [\text{cm}^{-3}]$	$\bar{\kappa}$
0602_NC	–	93	–
0602_NA	707	–	0.55
0607_NC	–	267	–
0607_NA	2073	–	0.35

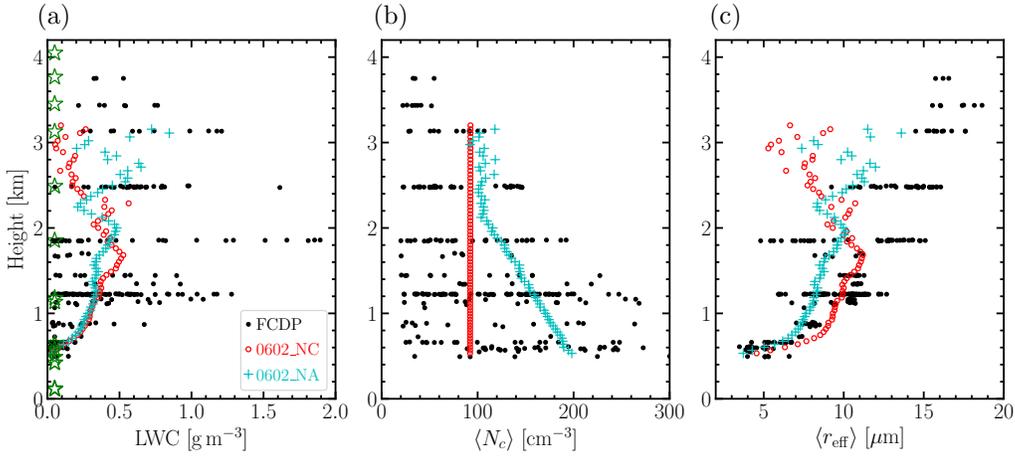
291 the dropsonde circle for each case show strong spatiotemporal variations as well. The  
 292 meteorology state from LES is evaluated against the dropsonde measurements. Com-  
 293 pared to dropsonde measurements, both the 0602\_NC and 0602\_NA simulations yield colder  
 294  $\theta$  and larger  $q_v$  below about 3.5 km and vice versa above, as shown in Figure 2(a) and  
 295 (b). We then compare the cloud properties between LES and the FCDP measurements.  
 296 Both simulations capture the measured LWC as shown in Figure 3(a) and the correspond-  
 297 ing statistics in Figure S15(a). Simulation 0602\_NA slightly overestimates (underesti-  
 298 mates)  $N_c$  ( $r_{\text{eff}}$ ) as shown in Figure 3(b) and (c) by comparing the red circles and black  
 299 dots. Overall, our simulations capture the observed cloud properties reasonably well de-  
 300 spite the aforementioned challenges.

301 To quantify the aerosol effect on precipitating cumuli, we adopt the metric of per-  
 302 centage difference (PD), defined as  $\text{PD} = (\mathcal{Q}_{\text{NA}} - \mathcal{Q}_{\text{NC}}) / \mathcal{Q}_{\text{NC}} \times 100\%$  with  $\mathcal{Q}_{\text{NC}}$  and  
 303  $\mathcal{Q}_{\text{NA}}$  representing quantities from the NC (baseline) and NA simulations, respectively  
 304 (Li et al., 2023).  $\mathcal{Q}$  is averaged between 08:00 and 20:00 UTC (Table S4). Simulation  
 305 0602\_NA yields 52.7% larger (–6.6% smaller)  $N_c$  ( $r_{\text{eff}}$ ) compared to 0602\_NC as shown  
 306 in Figure S16(e) and (f), suggesting a significant Twomey effect (Twomey, 1977). An  
 307 increased  $N_c$  due to the aerosol loading suppresses the precipitation (the RWP is reduced  
 308 by 38.9%) and leads to larger LWP (5.8%) as shown in Figure 4(a). This LWP adjust-  
 309 ment is consistent with the findings in Albrecht (1989), which suggested that increas-  
 310 ing cloud condensation nuclei (CCN) decreases the drizzle production and therefore in-  
 311 creases the LWP for shallow marine clouds. CFC from simulation 0602\_NA is larger until  
 312 16:00 UTC and then becomes smaller compared to 0602\_NC (Figure 4(b)). In total,  
 313 the aerosol loading leads to a 6.7% decrease of CFC, which is contrary to an increased  
 314 CFC due to aerosol loading as suggested in Albrecht (1989). CFC increases and then de-  
 315 creases monotonically with RWP. The timing of the CFC peak is consistent with the one  
 316 of RWP evolution (Figure S16). The shortwave cloud forcing at the top of the model de-  
 317 creases by  $3.2 \text{ W m}^{-2}$  (a 4.4% decrease), suggesting a net aerosol cooling effect due to  
 318 aerosol-cloud-precipitation interactions.

319 To quantify the aerosol effect on precipitation change, we examine the precipita-  
 320 tion susceptibility defined as  $S_o = -\Delta \ln R_p / \Delta \ln N_c$ , where  $R_p$  is the precipitation rate.  
 321 The rain frequency ( $E_p$ ) susceptibility  $-\Delta \ln \bar{E}_p / \Delta \ln \langle N_c \rangle$  is also examined. The pre-  
 322 cipitation susceptibility to aerosol perturbation depends on the LWP threshold (Sorooshian  
 323 et al., 2009) and cloud thickness (Jung et al., 2016).  $R_p$  is least susceptible to  $N_c$  for weakly  
 324 precipitating shallow MBL clouds because of low LWP ( $\leq 500 \text{ g m}^{-2}$ ), is less suscepti-  
 325 ble to deeper convective BL clouds because of the large abundance ( $\geq 1000 \text{ g m}^{-2}$ ) of  
 326 LWP, and is most susceptible to MBL clouds with intermediate LWP ( $\sim 500\text{--}1000 \text{ g m}^{-2}$ ).  
 327  $R_p$  from simulation 0602\_NC (red dots with error bars) is larger than that from 0602\_NA  
 328 (cyan dots), as shown in Figure 5, because the prescribed aerosol loading leads to a larger  
 329  $N_c$ , smaller droplet size, and weaker precipitation rate. Consistent with the aerosol ef-  
 330 fect on  $R_p$ , the precipitation event  $E_p$  is also reduced in 0602\_NA, as shown by the dashed  
 331 curves in Figure 5. This aerosol-induced suppression of precipitation is further quanti-  
 332 fied by a positive value of  $S_o$  and  $-\Delta \ln \bar{E}_p / \Delta \ln \langle N_c \rangle$  between NA and NC simulations

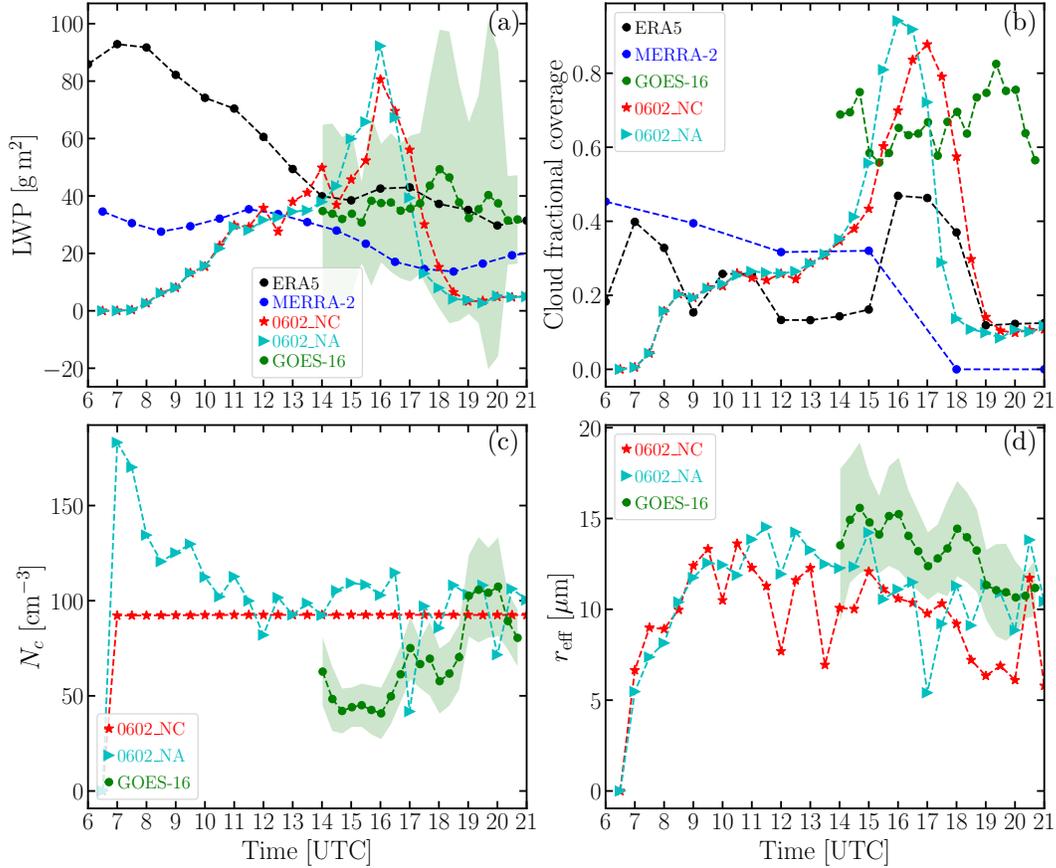


**Figure 2.** Domain-averaged vertical profiles from the WRF-LES simulation with the corresponding input forcings shown in Figure S23 at the measurement time for the 02 June 2021 case. The black line represents the ERA5 reanalysis data and the red (0602\_NC) and cyan (0602\_NA) ones represent the WRF-LES averaged values during the measurement time. The grey curves represent the dropsonde measurement with  $\pm\sigma$  error bars.



**Figure 3.** Comparison of vertical profiles of LWC,  $\langle N_c \rangle$ , and  $\langle r_{\text{eff}} \rangle$  between the WRF-LES (0602\_NC and 0602\_NA listed in Table 1) and the FCDP measurements (black dots). A threshold of  $\text{LWC} = 0.02 \text{ g m}^{-3}$ , effective diameter  $d_{\text{eff}} = 3.5 \mu\text{m}$  and  $N_c = 20 \text{ cm}^{-3}$  is applied to both the WRF-LES and the FCDP data. For the WRF-LES, only grid cells within clouds are averaged to obtain the vertical profile. The corresponding mean vertical profile of LWC,  $\langle N_c \rangle$ , and  $\langle r_{\text{eff}} \rangle$  is obtained by averaging three snapshots of WRF-LES output (30 minutes apart). The green stars mark all flight legs above cloud base (ACB) and below cloud top (BCT).

333 shown in the penultimate column of Table 2, respectively.  $R_p$  from simulation 0602\_NC  
 334 (red dots with error bars) is larger than that from 0602\_NA (cyan dots), as shown in Fig-  
 335 ure 5(a), because the prescribed aerosol loading leads to a larger  $N_c$ , smaller droplet size,  
 336 and weaker precipitation rate. Consistent with the aerosol effect on  $R_p$ , the precipita-  
 337 tion event  $E_p$  is also reduced in 0602\_NA, as shown by the dashed curves in Figure 5(b).  
 338 This aerosol-induced suppression of precipitation is further quantified by a positive value  
 339 of  $S_o$  and  $-\Delta \ln \bar{E}_p / \Delta \ln \langle N_c \rangle$  between NA and NC simulations shown in the penultimate  
 340 column of Table 2, respectively. The domain averaged LWP from our LES is less than  
 341  $100 \text{ g m}^{-2}$  for both cases, which leads to a small mean rain rate and  $S_o$  (0.86 for the 02  
 342 June case). This is consistent with the findings in Sorooshian et al. (2009).



**Figure 4.** Comparison of cloud properties between the WRF-LES (0602\_NC and 0602\_NA in red and cyan, respectively), ERA5 (black), MERRA-2 (blue), and GOES-16 (green) for the 02 June 2021 case. The domain averaged LWP includes both cloud and rain water.  $N_c$  and  $r_{\text{eff}}$  are averaged (cloudy grid with  $\text{LWC} \geq 0.02 \text{ g kg}^{-1}$ ) over the cloud top (200 – 300 m) from the WRF-LES output. ERA5, MERRA-2, and GOES-16 data are averaged over the dropsonde area. The GOES-16 LWP,  $N_c$ , and  $r_{\text{eff}}$  data are filtered by a cloud optical depth threshold  $\geq 3$  to limit the systematic biases in LWP and  $r_{\text{eff}}$  following the procedure described in Painemal & Zuidema (2011) and Painemal et al. (2021).

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### 3.2 Aerosol effect on drizzling cumuli: 07 June 2021 case

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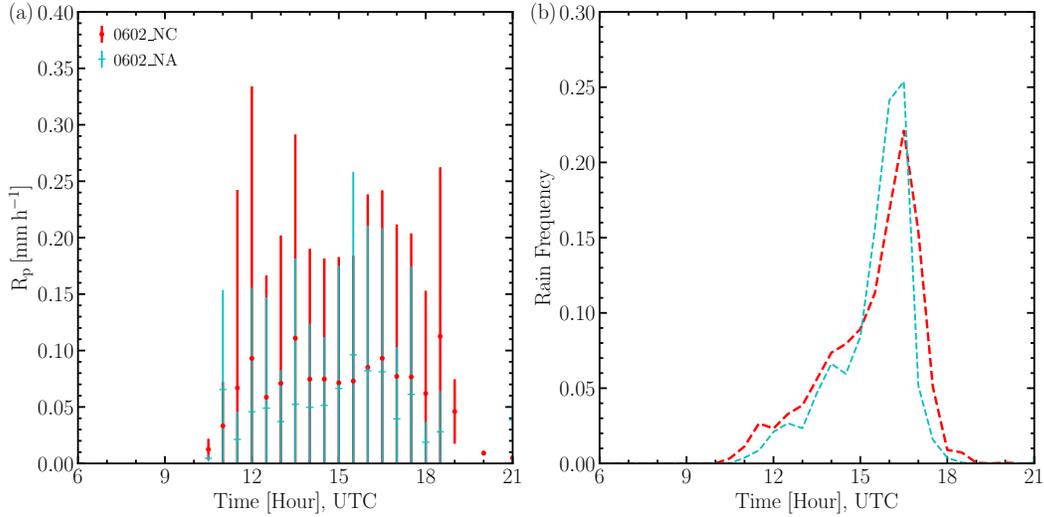
353

Compared to the 02 June 2021 case, the 07 June 2021 one is initially more polluted, in which  $N_a$  is about three times larger, as shown in the second column of Table 1. As a result, only drizzle was observed for the 07 June case. We first compare the vertical profiles (meteorology states) amongst the dropsonde measurement, LES, and ERA5 datasets. Both simulations, 0607\_NC and 0607\_NA, can reproduce the  $\theta$ -profile compared to the dropsonde measurement, except for a warmer free troposphere between 4-6 km as shown in Figure 6(a). LES produces a more humid boundary layer until 2 km compared to the dropsonde measurements and ERA5, as shown in Figure 6(b). The LES horizontal wind components agree with the dropsonde measurements (Figure 6(e) and (f)). Overall, the LES captures the observed MBL meteorology states.

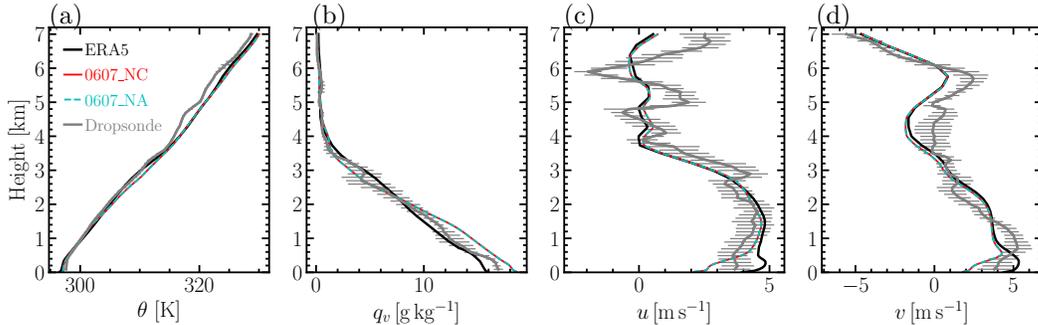
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LES cloud microphysical properties for this case are also evaluated against the FCDP measurements. Figure 7 shows that the measured LWC and  $N_c$  are very scattered, in-



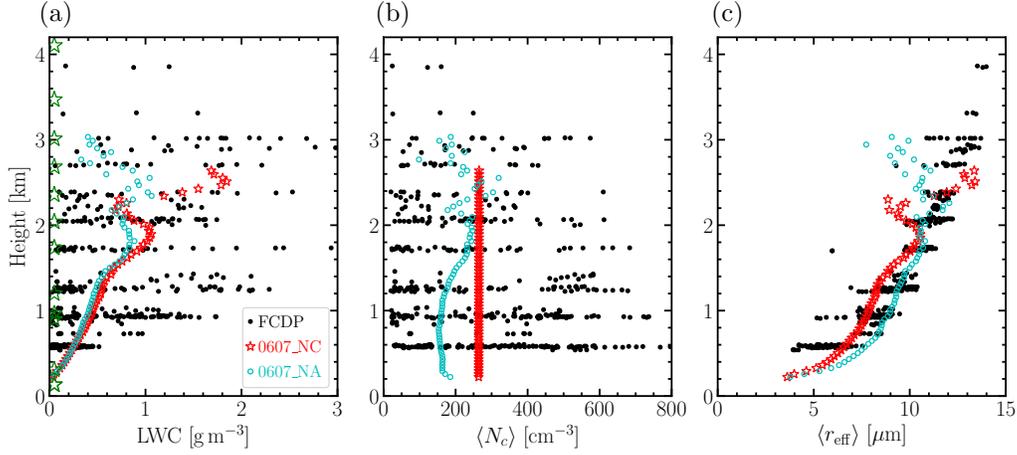
**Figure 5.** Precipitation rate (a)  $R_p$  (dots and pluses with error bars) and rain event frequency (b)  $E_p$  (dashed lines) for the 02 June 2021 case. A threshold of  $LWP > 50 \text{ g m}^{-2}$  and  $R_p > 0.004 \text{ mm h}^{-1}$  is applied to each grid to define rain events following Table 5 of Jiang et al. (2010). See Table 2 for the  $S_o$  calculation.



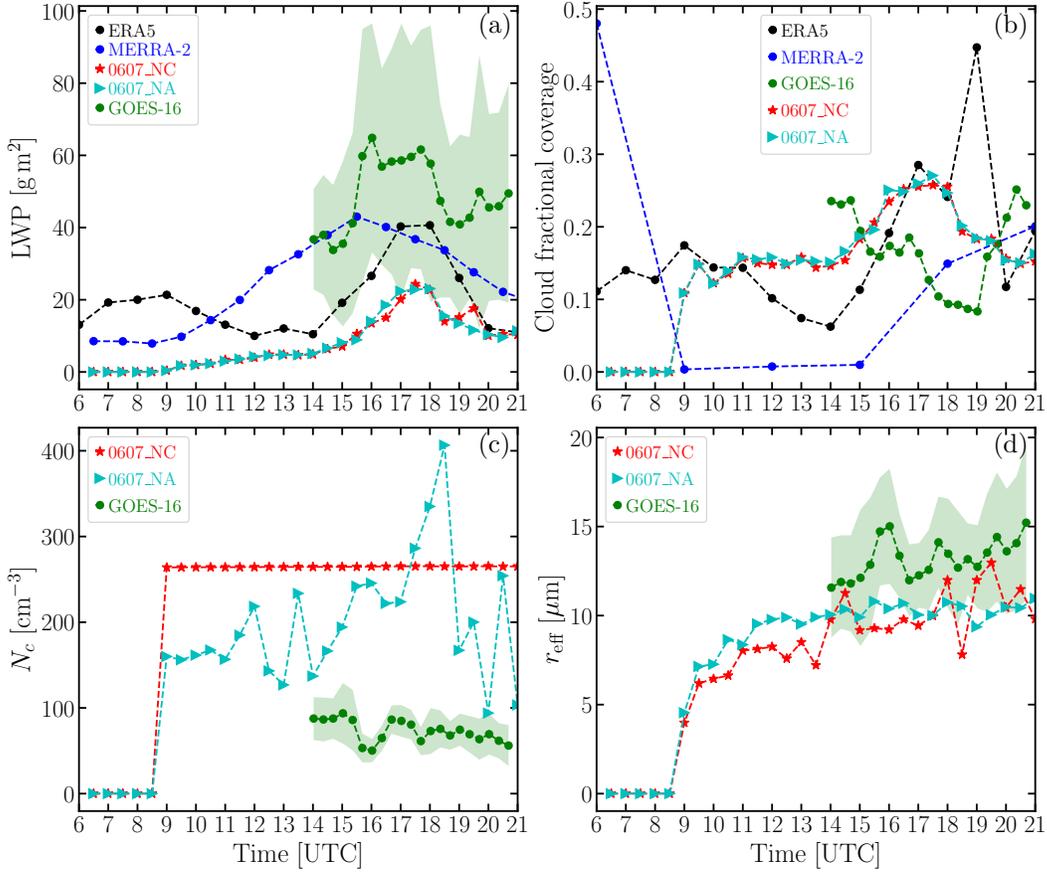
**Figure 6.** Same as Figure 2 but for the 07 June case.

356 indicating a large spatial variability, but both simulations can reproduce the LWC and  $r_{\text{eff}}$   
 357 from the FCDP measurements well (see statistics in Figure S17).

358 After the evaluation of modeled meteorology state and cloud properties against the  
 359 measurements, we now investigate the aerosol effect on drizzling (i.e., weakly precipitat-  
 360 ating) cumuli. Aerosol effect on LWP (PD =  $-0.7\%$ ) and CFC (PD =  $1.6\%$ ) are almost  
 361 negligible by comparing the time evolution of LWP and CFC for simulation 0607\_NC (red  
 362 stars) and 0607\_NA (cyan triangles) as shown in Figure 8(a) and (b). This is likely be-  
 363 cause of the light precipitation, which is consistent with the wintertime ACTIVATE stra-  
 364 tocumulus cases described in (Li et al., 2023).  $N_c$  ( $r_{\text{eff}}$ ) from simulation 0607\_NA is close  
 365 to that from 0607\_NC as shown in Figure S26(e) and (f). The impact of prescribed aerosols  
 366 on the decrease in  $N_c$  (PD =  $-31.6\%$ ) and the increase in  $r_{\text{eff}}$  (PD =  $12.9\%$ ) mostly re-  
 367 flects the Twomey effect. The aerosol induced RWP-reduction is  $17.4\%$ . The overall net  
 368 cooling effect is  $0.8 \text{ W m}^{-2}$  ( $2.7\%$ ) in terms of short-wave (SW) cloud forcing at the top  
 369 of the model.



**Figure 7.** Same as Figure 3 but for the 07 June case.



**Figure 8.** Same as Figure 4 but for the 07 June 2021 case. A comparison of the vertical profiles is shown in Figure 11.

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### 3.3 Aerosol-induced LWP and CFC adjustment: entrainment and precipitation

As discussed in section 1, the LWP adjustment to aerosol-induced  $N_c$ ,  $\Delta \ln \overline{\text{LWP}} / \Delta \ln \overline{\langle N_c \rangle}$ , is nonlinear and depends on cloud regimes. In this section, we examine  $\Delta \ln \overline{\text{LWP}} / \Delta \ln \overline{\langle N_c \rangle}$

for the two cases studied here and contributing factors, i.e., precipitation and cloud-top entrainment.  $\Delta \ln \overline{\text{LWP}} / \Delta \ln \langle N_c \rangle$  is calculated by averaging the time series of LWP and  $\langle N_c \rangle$  between 08:00 and 20:00 UTC. The positive value, 0.13 and 0.02 for the 02 June and 07 June case, respectively (Table 2), indicates precipitation-dominated LWP adjustments. The positive LWP adjustment leads to thicker and more reflective clouds. This is consistent with previous LES (Glassmeier et al., 2021) and satellite (Christensen et al., 2022) studies. We note that even though the 02 June case is clean and heavily precipitating and the 07 June one is polluted and lightly drizzling, the LWP adjustment due to aerosols are quite similar, suggesting that precipitation-dominated LWP adjustment, in response to the small aerosol or  $N_c$  perturbations, may not depend on the precipitation strength. The aerosol impact on cloud radiative effect can be quantified by the perturbation of cloud optical depth  $\tau_c$  to  $N_c$  (Ghan et al., 2016),

$$\frac{\Delta \ln \overline{\tau_c}}{\Delta \ln \langle N_c \rangle} = \frac{\Delta \ln \overline{\text{LWP}}}{\Delta \ln \langle N_c \rangle} - \frac{\Delta \ln \langle r_{\text{eff}} \rangle}{\Delta \ln \langle N_c \rangle}. \quad (1)$$

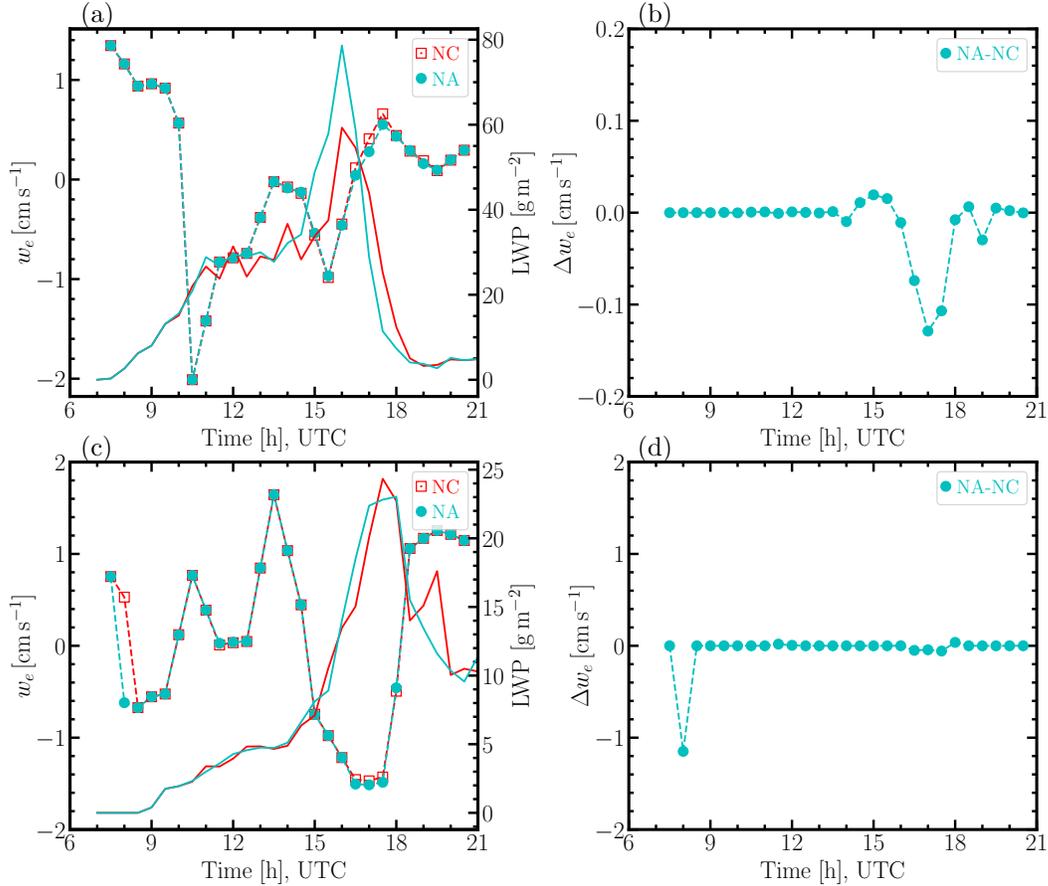
Equation (1) shows that both the Twomey effect (second term) and the cloud macro-physical adjustment (first term) contribute to  $\tau_c$ . For the 02 June case, the Twomey effect (0.2) and LWP adjustment (0.1) terms are comparable, which leads to a positive  $\Delta \ln \overline{\tau_c} / \Delta \ln \langle N_c \rangle$  of 0.3 (Table 1). For the 07 June case, the Twomey effect determines  $\Delta \ln \overline{\tau_c} / \Delta \ln \langle N_c \rangle$  and the LWP adjustment effect is negligible.

Entrainment is another important process contributing to cloud macrophysical adjustments. We first examine the 02 June 2021 case.  $w_e$  and LWP are anti-correlated with a Pearson correlation efficient of -0.39 (p-value=0.11) for simulation 0602\_NC between 12:00-20:30 UTC (17 snapshots are used for the statistics) as shown in Figure 9(a). The same conclusion can be drawn for simulation 0602\_NA but with a Pearson correlation efficient of -0.64 (p-value=0.005).  $w_e$  from simulation 0602\_NC is slightly larger than that from 0602\_NA as shown in Figure 9(b) from 16:30-17:30 UTC.

For the 07 June case,  $w_e$  is anti-correlated with LWP (Figure 9(c)) with a Pearson correlation coefficient of -0.45 (p-value = 0.06) and -0.56 (p-value = 0.01) for simulation 0607\_NC and 0607\_NA, respectively. This indicates that the cloud-top entrainment process has a pronounced effect on LWP for the drizzling cumuli, consistent with non-precipitating marine stratocumuli (Ackerman et al., 2004), where entrainment plays a significant role. The net shortwave radiative flux at the model top (not shown) shows moderate correlation with the  $w_e$  with a Pearson correlation coefficient of 0.66 (p-value = 0.003) and 0.69 (p-value = 0.001) for simulation 0607\_NC and 0607\_NA, respectively. The 0607\_NC simulation yields slightly larger  $w_e$  from 16:30-17:30 UTC as can be seen from the time evolution of  $\Delta w_e$  (Figure 9(d)). Since the time-varying large-scale vertical velocity profile (based on ERA5 forcing)  $\langle w \rangle_{z_i}$  is the same for the two simulations,  $\Delta w_e$  is due to the  $dz_i/dt$ , which is caused by the difference in cloud properties and consequent radiative impact on boundary layer structure for both cases.

### 3.4 Evaluation of large-scale models using LES

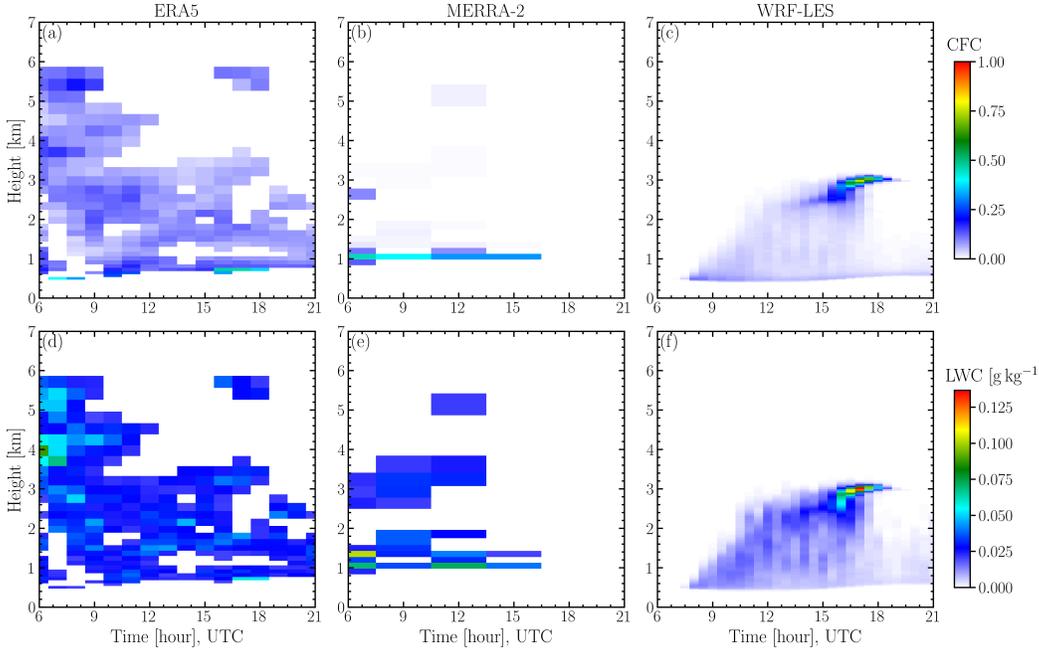
One of the goals of the present study is to evaluate the representation of cloud micro/macro-physics in large-scale models using LES and observations. ERA5 (black dots) agrees well with the GOES-16 measurements (green dots) in LWP while MERRA-2 (blue dots) shows smaller LWP in 14:00-21:00 UTC, as shown in Figure 4(a) for the 02 June 2021 case. CFC from ERA5 and MERRA-2 is smaller compared to GOES-16 (Figure 4(b)). The LES does not capture the spatial structure of LWP (Figure S18) or CFC (Figure 4(b)) compared to GOES-16. The LES  $N_c$  ( $r_{\text{eff}}$ ) is larger (smaller) than GOES-16 as shown in Figure 4(c) and (d). However, we note that the GOES-16  $N_c$  is smaller than FCDP- $N_c$  during the FCDP measurement time (18:30-19:12 UTC). The time evolution of the domain averaged CFC from ERA5 (black dots) exhibits the same diurnal cycle as the LES (red stars and cyan triangles) (Figure 4(c)). However, compared to LES, ERA5 data exhibit



**Figure 9.** Entrainment rate  $w_e = dz_i/dt - \langle w \rangle_{z_i}$  and the corresponding difference between NA and NC simulations  $\Delta w_e$  (squares) for the 02 ((a) and (b)) and 07 ((c) and (d)) June 2021 cases, where the cloud top height  $z_i$  is determined by the threshold  $LWC \geq 0.02 \text{ g kg}^{-1}$ .  $\langle w \rangle_{z_i}$  is the ERA5 large-scale vertical velocity at  $z_i$ . Solid lines in (a) and (c) represent the corresponding LWP.

409 higher clouds while MERRA-2 produces too low and little clouds as shown by the time  
 410 evolution of CFC and LWC vertical profiles from the reanalysis data and LES in Fig-  
 411 ure 10.

412 For the 07 June 2021 case, the LES (red stars and cyan triangles) produces  $\sim 1/3$   
 413 LWP of GOES-16 (green symbols) during 14:00-21:00 UTC, as shown in Figure 8(a). The  
 414 LES does not reproduce the GOES-16 LWP as shown in Figure S19. The CFC from LES  
 415 is larger than that from GOES-16 (Figure 8(b)). The ERA5 LWP (black symbols) fol-  
 416 lows the same diurnal cycle as LES (red stars and cyan triangles) even though the mag-  
 417 nitude is 2 times larger in 12:00-21:00 UTC (Figure 8(b)). This is remarkable consid-  
 418 ering the fact that cumuli hardly reach any steady state compared to the stratocumuli  
 419 and that the ERA5 grid-spacing (30 km) is 300 times coarser than the LES (100 m).  
 420 MERRA-2 (blue dots) has higher LWP than LES and ERA5. Neither the vertical struc-  
 421 ture of the ERA5 nor the MERRA2 LWC resembles the ones from LES (Figure 11(d)-  
 422 (e)). The cloud vertical extent from ERA5 (Figure 11(d)) reaches 3 km between 09:00-  
 423 12:00 UTC and 6 km around 18:00 UTC compared to 2 km from the LES (Figure 11(e)).  
 424 The ERA5 CFC agrees reasonably with the LES while the MERRA-2 CFC is smaller  
 425 than LES as shown in Figure 8(b). However, neither the ERA5 nor MERRA2 capture



**Figure 10.** Evolution of vertical profile of CFC and LWC for the 02 June 2021 case. They are obtained by sampling each layer conditionally with a threshold of  $LWC = 0.02 \text{ g cm}^{-3}$ . The 0602\_NC vertical profiles are calculated by normalizing the cloudy grids with the total number of grids ( $200 \times 200$ ) at each model level. ERA5 and MERRA-2 data are averaged over the dropsonde circle area.

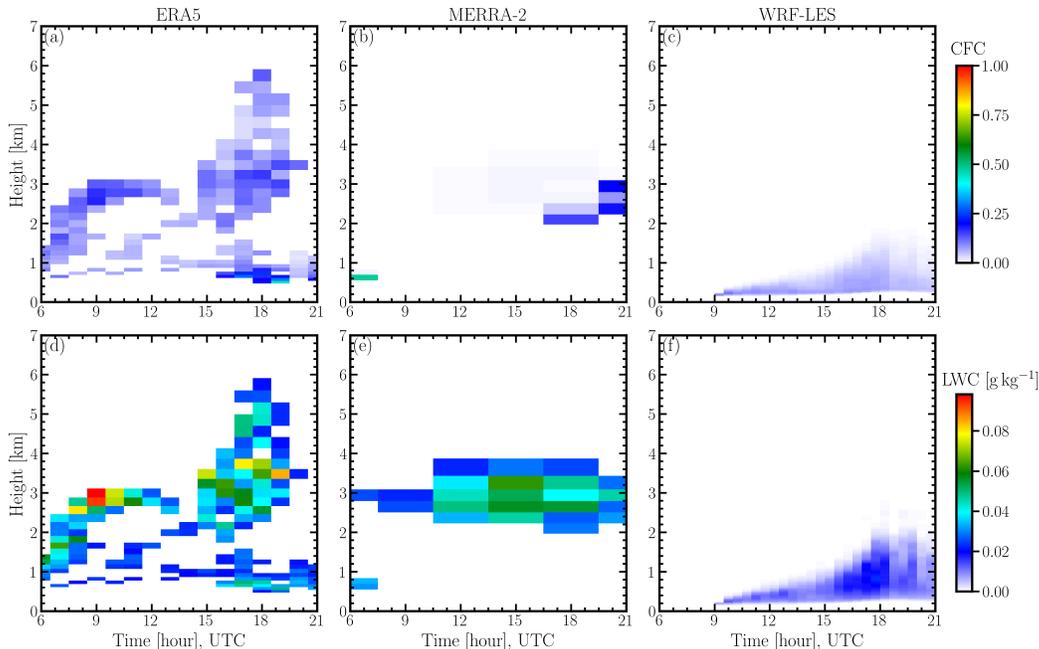
**Table 2.** Aerosol perturbation induced susceptibility of LWP and  $r_{\text{eff}}$  to  $N_c$  (cloudy average) between 08:00 and 20:00 UTC for the 02 and 07 June 2021 cases.

Case	$\Delta \ln \overline{LWP} / \Delta \ln \langle N_c \rangle$	$-\Delta \ln \langle r_{\text{eff}} \rangle / \Delta \ln \langle N_c \rangle$	$\Delta \ln \overline{\tau_c} / \Delta \ln \langle N_c \rangle$	$-\Delta \ln \overline{R_p} / \Delta \ln \langle N_c \rangle$	$-\Delta \ln \overline{E_p} / \Delta \ln \langle N_c \rangle$
02-06-2021	0.13	0.16	0.29	0.86	0.20
07-06-2021	0.02	0.32	0.34	–	–

426 the vertical structure of CFC compared to the LES (Figure 11(a)-(c)). We note that LES  
 427 underestimates the observed cloud top height by about 1 km compared to in-situ mea-  
 428 surements. This again demonstrates the challenge in simulating precipitating cumulus  
 429 even using LES.

#### 430 4 Discussions, conclusions, and outlook

431 We study aerosol-cloud-precipitation interactions in summertime precipitating shal-  
 432 low cumuli observed over the WNAO during the ACTIVATE campaign using LES. Two  
 433 contrasting observational cases are selected. The 02 June 2021 case is a cleaner case fea-  
 434 turing heavier precipitation, while the 07 June 2021 case is a more polluted one with lightly  
 435 drizzling conditions. Both cases are very challenging to simulate due to the strong spa-  
 436 tial variation of humidity and relatively deep boundary layer. For each case, the base-  
 437 line LES is initiated with a constant droplet number concentration  $N_c$  from the ACTI-  
 438 VATE in-situ (FCDP) measurement. To perturb the LES clouds, we performed a sensi-  
 439 tivity experiment with prescribed aerosol size distributions derived from SMPS/LAS



**Figure 11.** Same as Figure 10 but for the 07 June 2021 case.

440 measurements and the hygroscopicity  $\kappa$  derived from the AMS measurements. The LES  
 441 experiments are forced by large-scale forcings, i.e., advective tendencies of  $\theta$  and  $q_v$  and  
 442 surface heat fluxes from ERA5 reanalysis data. The simultaneous measurements of the  
 443 meteorology state and cloud properties allow us to evaluate our LES results at both small  
 444 and large scales that are essential for understanding ACI. For the 02 June 2021 case, LES  
 445 yields a slightly colder and more humid MBL compared to the dropsonde measurements.  
 446 LES can reproduce the FCDP measurements of cloud microphysical properties, which  
 447 agree reasonably well with the satellite retrievals. For the 07 June 2021 case, LES cap-  
 448 tures the  $\theta$  profile well but produces a more humid MBL. The cloud microphysical prop-  
 449 erties (LWC,  $N_c$ , and  $r_{\text{eff}}$ ) from LES agree with FCDP measurements. Overall, the LES  
 450 is able to reproduce the measured cloud microphysics although the spatial variability is  
 451 challenging to simulate. To capture the spatial variability, we perform simulations using  
 452 ERA5 large-scale forcings at the location of individual dropsondes (see Figure S20–  
 453 S22 in the supplement), none of which reproduce the observed clouds. This shows the  
 454 challenge in simulating fast-evolving marine cumuli. The LES fails to reproduce the spa-  
 455 tial structure of LWP compared to GOES-16 and does not agree well with the satellite  
 456 microphysics retrievals, even though the LES reproduces the LWP compared to the Re-  
 457 search Scanning Polarimeter (RSP) retrievals (Figure S25). The former could be because  
 458 of the spatially uniform boundary conditions adopted in our LES that lack the mesoscale  
 459 organizational structures shown in GOES-16 cloud field. The latter is likely due to both  
 460 the idealized boundary conditions of the LES and uncertainties from the GOES-16 re-  
 461 trievals. The time evolution of LWP and CFC from ERA5 shows the same diurnal vari-  
 462 ation as LES for both cases although the fast-evolving subgrid shallow cumuli are chal-  
 463 lenging to simulate in ERA5 with much coarser spatio-temporal resolution than in LES.

464 For the clean and heavily precipitating case, LES predicts a larger  $N_c$ , based on  
 465 the observed aerosol size distribution and hygroscopicity, than the observed one, result-  
 466 ing in a suppression of precipitation and a larger LWP. This mechanism is consistent with  
 467 many previous sensitivity studies of ACI. The CFC decreases as  $N_c$  increases despite an

468 increased LWP. To the best of our knowledge, it is the first time that this mechanism  
469 is tested in LES driven by measured cloud microphysics in the WNAO region. For the  
470 more polluted, lightly drizzling case, the aerosol loading predominately affects  $N_c$  and  
471  $r_{\text{eff}}$  and has negligible effect on LWP and CFC, reflecting the Twomey effect alone. The  
472 LWP adjustment is dominated by precipitation change and is anti-correlated with the  
473 cloud-top entrainment rate for both cases.

474 The aerosol effect on precipitation rate is strongly nonlinear. The precipitation rate  
475  $R_p$  has been argued to decrease with increasing aerosol number concentration  $N_a$  due  
476 to the suppression of collision-coalescence processes at fixed LWP in warm MBL (Albrecht,  
477 1989). The assumption of a statistically steady LWP largely holds for stratocumuli (Glass-  
478 meier et al., 2021) but fails for the cumuli. This makes it challenging to quantify the pre-  
479 cipitation susceptibility in cumuli. Positive precipitation susceptibility is observed for  
480 the heavily precipitating 02 June 2021 case due to the aerosol input that suppresses the  
481 precipitation.  $R_p$  is less susceptible to  $N_c$  for the lightly drizzling 07 June 2021 case. Our  
482 finding is consistent with previous studies of the  $R_p$ — $N_a$  relationship for warm MBL clouds  
483 (Jung et al., 2016). Whether the aerosol effect on precipitation rate observed in the two  
484 cases here can be generalized to global scales remains to be investigated.

**Acknowledgments**

This work was supported through the ACTIVATE Earth Venture Suborbital-3 (EVS-3) investigation, which is funded by NASA's Earth Science Division and managed through the Earth System Science Pathfinder Program Office. The Pacific Northwest National Laboratory (PNNL) is operated for the U.S. Department of Energy by Battelle Memorial Institute under contract DE-AC05-76RLO1830. S.K. and C.V. were funded by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) – TRR 301 – Project-ID 428312742.

**Availability Statement**

The source code used for the simulations of this study, the Weather Research and Forecasting (WRF) model, is freely available on <https://github.com/wrf-model/WRF>. The simulations were performed using resources available through Research Computing at PNNL. Model input and output files are available at <https://doi.org/10.5281/zenodo.8034149>. ACTIVATE observational data are publicly available at <https://asdc.larc.nasa.gov/project/ACTIVATE>. GOES-16 data can be obtained at <https://satcorps.larc.nasa.gov/prod/exp/activate/visst-pixel-netcdf/g16-sd/2021/>. ERA5 reanalysis data are available at <https://doi.org/10.24381/cds.adbb2d47>. MERRA-2 reanalysis data can be obtained at <https://disc.gsfc.nasa.gov/datasets?project=MERRA-2>.

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