

An assessment of representing land-ocean heterogeneity via convective adjustment timescale in the Community Atmospheric Model 6 (CAM6)

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Abstract

The time needed by deep convection to bring the atmosphere back to equilibrium is called convective adjustment timescale or simply adjustment timescale, typically denoted by τ . In the Community Atmospheric Model (CAM), convective adjustment timescale is a tunable parameter with one value, 1 hour, worldwide. Albeit, there is no justified reason why one adjustment timescale value should work over land and ocean both. Continental and oceanic convection are different in terms of the vigor of updrafts and hence can have different longevities. So it is logical to investigate the prescription of two different convective adjustment timescales for land (τ_l) and ocean (τ_o). To understand the impact of representing land-ocean heterogeneity via τ , we investigate CAM climate simulations for two different convective adjustment timescales for land and ocean in contrast to having one value globally.

Following a comparative analysis of 5-year-long climate simulations, we find $\tau_l = 4$ hrs and $\tau_o = 1$ hr to yield the best results. Particularly, we find better MJO simulations. Although these τ values were chosen empirically and require further tuning, the conclusion of our finding remains the same, which is the recommendation to use two different τ values for land and ocean.

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2 **convective adjustment timescale in the Community Atmospheric**
3 **Model 6 (CAM6)**

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6 **Key Points:**

- 7 • Two distinct values of convective adjustment timescale, τ , over land & ocean in the con-
- 8 vective parameterization scheme are prescribed.
- 9 • The mean climate stays qualitatively the same, except for a moister and colder near-surface
- 10 atmosphere for longer τ s over the oceans.
- 11 • A primary gain of using two different τ s for land and ocean is improved simulation of the
- 12 convectively coupled equatorial waves.

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Abstract

The time needed by deep convection to bring the atmosphere back to equilibrium is called convective adjustment timescale or simply adjustment timescale, typically denoted by τ . In the Community Atmospheric Model (CAM), convective adjustment timescale is a tunable parameter with one value, 1 hour, worldwide. Albeit, there is no justified reason why one adjustment timescale value should work over land and ocean both. Continental and oceanic convection are different in terms of the vigor of updrafts and hence can have different longevities. So it is logical to investigate the prescription of two different convective adjustment timescales for land (τ_L) and ocean (τ_O). To understand the impact of representing land-ocean heterogeneity via τ , we investigate CAM climate simulations for two different convective adjustment timescales for land and ocean in contrast to having one value globally.

Following a comparative analysis of 5-year-long climate simulations, we find $\tau_O = 4$ hrs and $\tau_L = 1$ hr to yield the best results. Particularly, we find better MJO simulations. Although these τ values were chosen empirically and require further tuning, the conclusion of our finding remains the same, which is the recommendation to use two different τ values for land and ocean.

1 Introduction

Deep convection is complex to parameterize [Arakawa, 2004]. While the explicit representation of deep convection is becoming a plausible option to navigate this "deadlock" [Randall *et al.*, 2003; Randall, 2013], for long-term projections of our climate, cumulus parameterization is still unavoidable. Hence, amidst the fierce emergence of convection-resolving models [Stevens *et al.*, 2019], various schemes to parameterize convection continue to develop. In particular, the recent decades have witnessed a surge of novel ideas that have accelerated this progress [Rio *et al.*, 2019, and references therein].

The "art" of tuning parameters used in convection parameterization schemes, or simply parameter tuning, plays a vital role in this development process [Hourdin *et al.*, 2017]. While deficiencies of convective parameterization are primary factors for model biases, it alone cannot alleviate all mode biases [Goswami *et al.*, 2017]. Hence, parameter sensitivity investigations are necessary not only to optimize the performance of a scheme but also to understand the extremities to which a scheme can be held responsible for biases in a simulation [Qian *et al.*, 2015; Goswami *et al.*, 2017]. In this study, we aim to contribute to understanding one tunable parameter, the convective adjustment timescale τ , by investigating the sensitivity of climate simulations to two dif-

44 ferent τ values for land and ocean in contrast to having one value globally in the Zhang-McFarlane
45 convective parameterization scheme [Zhang and McFarlane, 1995, ZM95 hereafter] in the Com-
46 munity Atmospheric Model (CAM), the atmospheric model of the Community Earth System Model
47 [Danabasoglu et al., 2020].

48 In CAM, deep convection is represented using the Zhang-McFarlane (ZM) convection pa-
49 rameterization scheme. The ZM is an adjustment-type convective parameterization scheme where
50 the atmospheric instability is removed via an adjustment towards a background state. In ZM, con-
51 vective available potential energy (CAPE) defines atmospheric instability, and τ is the CAPE con-
52 sumption time. In their paper, ZM95 used τ values of 2, 4, and 6 hours. To quote ZM95, "The
53 adjustment time scale determines the intensity and duration of convection for a given CAPE. With
54 small τ the convection is short-lived but intensity is high, on the other hand with larger τ the con-
55 vection is long-lived but of low intensity". ZM95 reported their scheme to be particularly sen-
56 sitive to the choice of τ . Since there is no strict range of τ , several studies investigated the sen-
57 sitivity of CAM simulations to different τ values. For example, Mishra and Srinivasan [2010]
58 used $\tau=[1,\infty]$. Contrasting water-vapor isotope simulations in a suite of CAM single-column sim-
59 ulations with a range of τ values, Lee et al. [2009] found their simulations to match better with
60 satellite observations with $\tau = 8$ hrs. Mishra [2011, 2012] prescribed $\tau = 8$ hrs in global climate
61 simulations and noted improvements in the simulations of tropical climate, especially the con-
62 vectively coupled equatorial waves. Evaluating 22 tunable parameters in CAM, Qian et al. [2015]
63 reported τ as one of the most critical tuning parameters. In all of the above studies, τ has a sin-
64 gles value globally.

65 One value of τ globally is not a logical choice because deep convection exhibits different
66 behaviors over continents and oceans [Hagos et al., 2013; Matsui et al., 2016; Roca et al., 2017;
67 Roca and Fiolleau, 2020]. Since the width of a thermal plume is steered by boundary layer height
68 [Williams and Stanfill, 2002], a deep continental boundary layer generates wider updraft veloc-
69 ities in deep convection [Lucas et al., 1994]. Matsui et al. [2016] provided a climatological view
70 of the contrast between oceanic and continental convective precipitating clouds from long-term
71 TRMM satellite multisensor statistics. They found large proportions of deep clouds over land.
72 Zipser et al. [2006] also found the most intense storms typically over continents. These obser-
73 vations suggest that the atmospheric deep convection over land is wider and stronger than those
74 over the oceans. In other words, atmospheric convection over land is shorter lived than that over
75 ocean [Roca et al., 2017]. It advocates for a shorter convection consumption time scale over land
76 than over oceans which motivated us to address the following question: although two different

77 τ values incorporating land-ocean inhomogeneity are logical, is it fruit-bearing in a model-simulated
78 climate? To answer this question, we investigate,

- 79 • response of the mean climate, and
- 80 • response of large-scale waves,

81 by contrasting 5-year-long climate simulations with and without incorporating land-ocean inho-
82 mogeneity via τ values.

83 Convective parameterization schemes, particularly adjustment-type schemes, are based on
84 the idea that convection takes some time to stabilize the atmosphere to a background state. Es-
85 sentially, this time taken is τ in the ZM scheme. Although numerically τ can have almost any value,
86 it is decided based on a scale separation between the convective activity of the individual clouds
87 and large-scale forcing. This concept is nicely depicted in Figure 1.1 of [Davies, 2008]. The graph
88 in that figure is a function of timescales associated with convection, and consists of a turbulent
89 initial segment indicating fluctuation of individual clouds, followed by a flat segment where these
90 fluctuations smooth out, and finally a segment corresponding to longer time-scales that shows
91 the evolution of the large scale forcing field itself. Conceptually, changing τ within a reasonable
92 range (within the flat segment of Figure 1.1 of [Davies, 2008]) should not result in a dramatic change
93 in the mean state of the simulated climate. We shall investigate it in detail in the first part of our
94 results section.

95 Some changes that we expect in our experiments are in the simulated organization of con-
96 vection. The organization of convection comes from the dynamic and thermodynamic impacts
97 of convection on the atmosphere. Simply put, it is the memory of convection [Davies *et al.*, 2009],
98 i.e. the fact that convection changes the large-scale properties, and can make their environment
99 favorable or unfavorable to subsequent convection. Identifying sources of convective memory
100 in cloud-resolving simulations, Colin *et al.* [2019] argued that the persistence of the state of con-
101 vection contributes to convective memory. Colin *et al.* [2019] also suggested that convective mem-
102 ory and organization interact mutually. By altering τ we essentially alter memory associated with
103 convection. Hence, we expect to see changes in convective organization. Taking a cue from Mishra
104 [2011], we anticipate improved convective organization in the tropics for longer τ . However, land-
105 ocean heterogeneity in τ is a unique feature of our experiments that we argue is essential based
106 on heterogeneity in the behavior of convection over land and ocean. As supporting evidence, we

107 shall present an analysis of equatorial waves focusing on the MJO to evaluate the organization
 108 of convection in the second part of our results section.

109 The paper is organized as follows. A brief description of the methodology is provided in
 110 Section 2. Section 3 evaluates the response of the model to different τ values. Finally, a few con-
 111 cluding remarks are provided in Section 4.

112 **2 Model and simulation details**

113 We used the atmospheric model of the Community Earth System Model, version 2.1.3 (CESM
 114 2.1.3) [Danabasoglu *et al.*, 2020], that is the Community Atmosphere Model, version 6 (CAM6),
 115 developed and maintained at the National Center for Atmospheric Research (NCAR), with lon-
 116 gitude and latitude specifications 1.25° and 0.9° , respectively, and 32 vertical levels. We forced
 117 the model by HadISST1 climatological monthly mean SST data provided by the Met Office Hadley
 118 Centre [Rayner, 2003]. In short, we performed CESM “F2000climo” simulations. In general,
 119 these are atmospheric simulations forced by present-day climatology. All simulations are 6 years
 120 long, and we analyzed the last 5 years of each simulation since, for atmosphere-only simulations,
 121 1-year spin-up is enough.

122 We performed 5 simulations. The one with out-of-the-box τ value of 1 hour globally is called
 123 the control (*CTRL*). In the next 3 simulations, we delayed the τ value over ocean (τ_O) to 2, 3 and
 124 4 hours keeping τ over land (τ_L) 1 hour. We called these 3 simulations *EXPT_{2h}*, *EXPT_{3h}* and
 125 *EXPT_{4h}*, respectively. We performed a last 5th experiment, named *EXPT_{slow}*, for which we used
 126 a τ value of 4 hours globally. Before starting our comparative analysis, we rename our first sim-
 127 ulation as *EXPT_{fast}*, which initially we had named CTRL, for clarity and better fluency of nar-
 128 ration of our findings. Table 1 depicts the τ values for different experiments.

129 Our analyses primarily show a comparison between the 5 aforementioned simulations. For
 130 some analyses we have used outgoing long-wave radiation (OLR) from NOAA ($2.5^\circ \times 2.5^\circ$; daily
 131 from 01-Jun-1974 to 12-Dec-2019) [Liebmann and Smith, 1996] as observational benchmark.

133 **3 Results**

134 **3.1 Mean Climate**

135 Since about 75% of the global surface is ocean, in the simulations of the mean climate, we
 136 expect a similar model response in our experiments by delaying τ only over the oceans, as ear-

Experiment Name	τ_L	τ_O
$EXPT_{fast}$	1hr	1hr
$EXPT_{2h}$	1hr	2hr
$EXPT_{3h}$	1hr	3hr
$EXPT_{4h}$	1hr	4hr
$EXPT_{slow}$	4hr	4hr

Table 1. τ values for different experiments

132

137 lier studies did by having a larger τ globally. An evaluation of some of the mean features of sim-
 138 ulated climate in our experiments confirm this. We find an increase in large-scale rainfall and a
 139 decrease in convective rain going from $EXPT_{fast}$ to $EXPT_{slow}$ (Fig 1 and Supplementary Fig
 140 S1). Similarly, we also notice warming in the lower levels, stronger warming in the upper lev-
 141 els, slight cooling in the mid-levels; moistening in the lower levels, and drying in the mid-levels
 142 (Fig 2 and Supplementary Fig S2). These features have been reported in earlier studies [for ex-
 143 ample, Fig 8 in *Mishra and Srinivasan, 2010*].

144 Investigating the mean features for land and ocean separately, we notice in addition, lower
 145 level (upper level) warming (cooling) is more (less) over land than over oceans (Fig 2). In the case
 146 of moisture, the letter "S" patterned vertical structure over the ocean is more curvy and squeezed
 147 down meaning lower level (middle level) moistening (drying) is stronger over oceans than over
 148 land and the respective peaks are vertically closer to the sea surface. These profiles, all together,
 149 indicate a model response to changes in τ in terms of the distribution of atmospheric convection
 150 and clouds, which impacts heating/cooling and moistening/drying of the air column (Supplemen-
 151 tary Fig S2). Essentially these responses indicate an accumulation of convective instability in the
 152 atmosphere with delaying of convective adjustment time scale. It is attributable to more low-level
 153 warming over the continents and more low-level moistening over the oceans. More moistening
 154 near the ocean surface is relatively straightforwardly understandable, and it is a consequence of
 155 the atmosphere taking longer to convect with larger τ . To a zero-order approximation, as a re-
 156 sult of the near-surface moisture pile-up in the oceanic regions, there is a moisture deficit in the
 157 lower levels over the continental regions (Fig 3 and Supplementary Fig S3a and S3b). Indeed it
 158 is apparent, in relative sense, in Fig 3. Although q_O does not exhibit a clear moistening signal,

159 the land drying in q_L is profound. The consequences are reflected in terms of changes in cloud
 160 cover. In an overall declining tendency of cloud cover, from $EXPT_{fast}$ to $EXPT_{slow}$, over the
 161 tropics high clouds decrease more steeply than low clouds. Low clouds decrease less rapidly over
 162 the ocean compared to those over land (Fig 4). It should be noted that cloud categories are ob-
 163 jectively defined in CESM. For example, low-level clouds are the ones below 700 hPa and high
 164 clouds are between 400 and 50 hPa. Cloud covers are integrated for each model level correspond-
 165 ing to respective cloud categories. In that regard, going from $EXPT_{fast}$ to $EXPT_{slow}$, low-cloud
 166 cover changes (Fig 4) are consistent with relative surface moistness over land and ocean (Fig 3).

167 Taken together, the altered vertical profiles of moisture and temperature, distribution of con-
 168 vective and large-scale rainfall, and associated clouds are consistent with the idea that convec-
 169 tion is short-lived and stronger for smaller τ values and long-lived and weaker for longer τ value.
 170 It is also evident from the solution of the CAPE equation in the ZM scheme, which can be ex-
 171 pressed as $CAPE(t) = CAPE_o \exp(-\frac{t}{\tau})$ in the absence of large-scale CAPE generation, where
 172 $CAPE_o$ is the values of CAPE at $t = 0$. A larger τ in this expression means a slower decay of
 173 CAPE. The duration of convection is essentially linked with its persistence and hence "mem-
 174 ory". We discuss its impact on the simulation of the equatorial waves in the following section.

175 3.2 Simulation of MJO variance and propagation

176 Organization is a primary feature of tropical convection. It essentially means a cluster of
 177 deep precipitating clouds tied together. An important question is, what brings these clouds to-
 178 gether? In other words, what causes convection to organize? One idea to see the organization of
 179 convection is through superpositions of convectively coupled equatorial waves (CCEWs). These
 180 atmospheric waves and tropical convection are entangled. In the tropics, the atmosphere responds
 181 to convective heating in terms of waves that, in turn, organize convection. Therefore, the fidelity
 182 of a model in simulating tropical climate is essentially its ability to simulate the CCEWs. A stan-
 183 dard metric to analyze CCEWs is the Takayabu-Wheeler-Kiladis (TWK) spectra [*Takayabu, 1994a,b;*
 184 *Wheeler and Kiladis, 1999*]. Figure 5 depicts the symmetric and asymmetric TWK-spectra for
 185 the observed and simulated outgoing long-wave radiation (OLR). Understandably, a striking fea-
 186 ture of the TWK-spectra of observed OLR shown in Fig 5a and b is the spectral power near the
 187 origin of the plots in the wavenumber range 1-5 and frequency 20-100 days, well known as the
 188 MJO. The MJO is a combination of or envelope of other waves in the equatorial atmosphere. Hence,
 189 the accuracy of MJO simulation is arguably a measure of the fidelity of accurate simulation of

190 waves in the atmosphere [Zhang *et al.*, 2020]. Guo *et al.* [2015] showed in detail that the accu-
 191 racy of CCEW simulation is critical for a realistic MJO simulation.

192 A comprehensive review of the science of MJO is available in Zhang *et al.* [2020]. Promi-
 193 nent observed features of MJO suggest that they are most active in the Indo-Pacific warm pool
 194 with an eastward propagation. An interesting fact, along its path from the Indian to the Pacific
 195 Ocean, is that an MJO passes over the Indonesian maritime continent (IMC). During this pas-
 196 sage, MJO and the prominent diurnal variabilities in the meteorology over the IMC islands in-
 197 teract and mutually influence each other. So much so that nearly half of the MJOs fail to prop-
 198 agate into the Pacific. It is critical, therefore, to represent the land-ocean heterogeneity as real-
 199 istically as possible in climate models. Hence, we expect our experiments with logically defined
 200 different values of τ for land and ocean to improve simulated MJO features. Here, we shall present
 201 analyses evaluating the simulation of MJO variance and propagation. We can draw some idea
 202 of MJO simulation in different experiments from Fig 5. In Fig 5, the foremost remarkable fea-
 203 ture is the increase in spectral power in the MJO wave number and frequency range for experi-
 204 ments with a longer τ . A closer visual inspection reveals that the MJO spectral power does not
 205 dramatically change from $EXPT_{2h}$ to $EXPT_{slow}$. For other waves, no one simulation is remark-
 206 ably better than the rest. Fig 5 loosely suggests that overall the symmetric signal waves are im-
 207 proved for longer time scales, but there are no clear improvement for the antisymmetric part.

208 To bring out the active region of MJO we applied space-time filtering on OLR data con-
 209 taining the signal corresponding to wavenumbers 1-5 and a period of 20–100 days. In Fig 6 the
 210 variance of the MJO-filtered daily OLR anomalies is shown. In observations (Fig 6a), the peak
 211 variance is over the Indo-Pacific warm pool. Feeble variance peaks are noted in the eastern sides
 212 of the Pacific (off the Gulf of California) and Atlantic (around the western coast of Sierra Leone).
 213 It is consistent with the fact that although MJO is most active in the Indo-Pacific warm pool re-
 214 gion, it has considerable influence modulating the convective activity over the eastern equato-
 215 rial Pacific [Maloney and Hartmann, 2000a,b; Maloney and Kiehl, 2002] and Atlantic [Klotzbach,
 216 2014]. For $EXPT_{fast}$ high variance is noted around the warm-pool region but widely spread and
 217 has multiple peaks. The strongest variance is around Northern Australia and the south-western
 218 Pacific region. The other secondary maxima are over the southern Bay of Bengal, the central equa-
 219 torial Indian Ocean, and the central Pacific regions.

220 The simulated MJO variance strength and pattern experience some changes with changes
 221 in τ values. In general, a slower τ_O keeping τ_L same yields more variance. In other words, it in-

222 creases convective activity in MJO space and time scales. In $EXPT_{2h}$ a pronounced peak is
 223 located over the western-central equatorial Pacific with two secondary maxima near the south-western
 224 equatorial Pacific and eastern equatorial Indian Ocean. In $EXPT_{3h}$ the variance is more concen-
 225 trated over the western equatorial Pacific, with a secondary peak south of the central equatorial
 226 Indian Ocean. With larger values of τ_L , the maximum variance gets more and more focused over
 227 the warm pool region, from $EXPT_{fast}$ to $EXPT_{3h}$ (comparing Fig 6b-d). It is noteworthy, that
 228 all the pronounced peaks for $EXPT_{2h}$ and $EXPT_{3h}$ are over oceans, in and around the Indo-Pacific
 229 warm pool region, but split unlike observations (Fig 6a). The model simulated MJO variance fur-
 230 ther slowing τ_O to 4 hours ($EXPT_{4h}$ shown in Fig 6e) suggests that MJO variance does not nec-
 231 essarily increase with increasing τ_O . The variance peak intensities are visibly weaker in $EXPT_{04}$
 232 compared to that in $EXPT_{2h}$ and $EXPT_{3h}$ and more only than that in $EXPT_{fast}$. However, a note-
 233 worthy feature of $EXPT_{4h}$, a fine detail missing in all other simulations, is the variance peaks
 234 near the eastern side of the equatorial Pacific and Atlantic oceans. Baring these subtle variance
 235 peaks, $EXPT_{slow}$ looks the best, although still a considerably weaker variance peak compared
 236 to observations. The variance fields normalized by the respective domain means are available
 237 in Supplementary Fig S4, which depicts a better visual illustration of the variance peaks.

238 A prominent feature of MJOs is eastward propagation. The propagation features of the MJO
 239 are arguably better characterized by Hovmöller plots averaged over the latitude band between 10°S
 240 and 10°N, shown in Fig 7. Each frame in Fig 7 depicts 10°S-10°N averaged cross-correlations
 241 of OLR anomalies with MJO-index. The MJO-index is defined as the 20-100-day filtered OLR
 242 anomalies averaged over 5°S-5°N, 75°E-85°E following *Guo et al. [2015]*. It is noteworthy to
 243 mention, reiterating *Guo et al. [2015]*, the philosophy behind using such an MJO index. An in-
 244 dex based on a 20-100 day filter brings out the dominant intraseasonal signal in the data that ide-
 245 ally should be an MJO signal. The eastward propagating red and blue patches of correlation val-
 246 ues in observations (Fig 7a) confirm it. We note the phase speed is faster over the west Pacific
 247 (east of $\sim 120^\circ\text{E}$) than that over the Indian Ocean (west of $\sim 100^\circ\text{E}$). The relatively slow phase speed
 248 in the longitude range $\sim 100^\circ\text{-}120^\circ\text{E}$ is collocated with the Indonesian archipelago. These dif-
 249 ferent phase speeds over land and oceanic regions are consistent with MJO interaction with the
 250 profound diurnal variations of meteorology over the MC. It furthermore emphasizes the need to
 251 mimic land-ocean heterogeneity realistically in climate models.

252 To assess the performance of our different experiments in simulating MJO propagation fea-
 253 tures, we recall the "good" and "bad" models of *Guo et al. [2015]*. In Figure 2, *Guo et al. [2015]*
 254 showed that the "good" models simulated more realistic eastward propagation than the "bad" mod-

255 els. In Fig 7, $EXPT_{4h}$ is the only experiment with an eastward propagation and exhibits some
 256 resemblance with observations and the only "good" model, albeit with some key caveats. The
 257 positive anomalies almost abruptly died over the MC and reappeared over the western Pacific.
 258 Nonetheless, an intriguing observation, that contains the novelty of our research, is the more re-
 259 alistic eastward propagation simulated in $EXPT_{4h}$ than in $EXPT_{slow}$. An improved simulation
 260 of eastward propagation in $EXPT_{4h}$ supports our argument that using two τ s for land and ocean
 261 is a logical choice. It reconfirms our anticipation that representing land-ocean heterogeneity via
 262 τ in ZM in CAM alters convective memory and affects the organization of convection. A larger
 263 τ_O than τ_L , although reasonable, is only based on intuition. Detailed sensitivity analysis would
 264 be needed to investigate and pin down the best pair of τ values.

265 **4 Discussion and Conclusion**

266 Climate models continue to grow, fueled by a growing understanding of the earth system.
 267 Hence, it is only logical to include a fairly well-recognized and relatively old knowledge about
 268 land and ocean heterogeneity of atmospheric convection in the parameterization of convection.
 269 We argue that using two different τ in ZM in CAM can be one simple yet fruit-bearing way. In
 270 our experiments to investigate the model response to land-ocean heterogeneity in τ values, we
 271 used $\tau_L = 1$ hr, and $\tau_O = 2$ hrs, 3 hrs, 4 hrs. In two additional experiments, $EXPT_{fast}$ and $EXPT_{slow}$,
 272 we used $\tau_L = \tau_O = 1$ hr and $\tau_L = \tau_O = 4$ hrs, respectively, to complement the previous group
 273 of experiments. The τ values that we have used are informed by our knowledge of frequency, life-
 274 cycle, and behavior of atmospheric convection over land and ocean learned from previous stud-
 275 ies [Lucas *et al.*, 1994; Williams and Stanfill, 2002; Zipser *et al.*, 2006; Hagos *et al.*, 2013; Mat-
 276 sui *et al.*, 2016; Roca *et al.*, 2017; Roca and Fiolleau, 2020] and inspired by results of relevant
 277 model sensitivity experiments [Zhang and McFarlane, 1995; Lee *et al.*, 2009; Mishra and Srini-
 278 vasan, 2010; Mishra, 2011; Misra *et al.*, 2012].

279 Our findings regarding the model simulated mean state in different experiments are con-
 280 sistent with earlier studies [Lee *et al.*, 2009; Mishra and Srinivasan, 2010; Mishra, 2011; Misra
 281 *et al.*, 2012]. For example, total rainfall remained approximately the same while large-scale rain-
 282 fall increased and convective rain decreased for longer τ_L s. Consistency of the model response
 283 for a slow τ only over the oceans with slowing down τ globally is most likely a result of 75% of
 284 the global surface being ocean. However, since there is no physical barrier between the atmospheric
 285 columns over continents and oceans, having two τ values in our experiments, which essentially
 286 are prescribed to represent heterogeneity in the persistence of convection over the two different

287 surfaces, created a distinction between the intensities with which the model responses are felt over
288 land and ocean. For example, the oceanic boundary layer is moister and warmer than the con-
289 tinental boundary layer (Fig 3). Furthermore, the mid-troposphere is drier and cooler over oceans
290 than over the continents (Fig 2). These land-ocean heterogeneities inevitably create differences
291 in atmospheric instabilities. These instabilities are essentially realized in the form of atmospheric
292 convection that, by design in our experiments with slower τ , takes longer to bring the atmosphere
293 back to a background state. It is suggestive of a longer persistence of convective instability over
294 the ocean than that over the continents which essentially can be linked with memory of convec-
295 tion [Davies et al., 2009; Colin et al., 2019; Hwong et al., 2023].

296 The conclusion that the model simulated better convectively coupled equatorial waves in
297 $EXPT_{2h}$ than in $EXPT_{slow}$ is a key. We conclude this based on our finding of a better MJO sim-
298 ulation in $EXPT_{2h}$, consistent with improved symmetric waves. Scientists had advocated in fa-
299 vor of a slower τ in earlier studies [Mishra, 2011; Misra et al., 2012]. We also noted a signifi-
300 cant increase in MJO power for $\tau = 4$ hrs than $\tau = 1$ hr (comparing Fig 5b and Fig 5f). However,
301 an evaluation of the model simulated intraseasonal zonal propagation reveals that $EXPT_{4h}$ per-
302 forms considerably better than $EXPT_{slow}$. This confirms that having one τ globally is not only
303 unphysical but also slowing down tinkering persistence of convection to improve simulation of
304 equatorial waves, and may result in model responses that might look improved, but only super-
305 ficially.

306 Our results, in general, serve as proof of concept that a realistic representation of convec-
307 tive adjustment time scale over land and ocean is a logical requirement that properly implemented
308 shall lead to improvements in climate model simulations. In specific, we advocate at least two
309 τ values, one for the continents and one relatively slower for the oceans in ZM in CAM. The fact
310 that we did not perform a rigorous model sensitivity analysis [e.g., Qian et al., 2015; Lin et al.,
311 2016; Goswami et al., 2017] nor did we perform any cloud-resolving simulation targeting the
312 life-cycle of atmospheric convection [Davies et al., 2013; Colin et al., 2019; Daleu et al., 2020,
313 e.g.,] leaves a scope as well as the requirement for future research to determine the best values
314 of τ_L and τ_O for ZM in CAM. It will hopefully guide convection parameterization schemes, es-
315 pecially the adjustment types, to address land-ocean heterogeneity. Specifically, we recommend
316 that future developments of CAM should consider prescribing different τ_L and τ_O in ZM in CAM.

5 Open Research

- 318 • Model : We used the atmospheric model of the Community Earth System Model, version
319 2.1.3 (CESM 2.1.3) [Danabasoglu *et al.*, 2020]
- 320 • Description of the model simulations is provided in Section 2 of the manuscript. A source
321 file of CESM 2.1.3, zm_conv.F90, modified for our experiments is provided in https://github.com/bidyutbg/CESM_Tau_experiment.git.
322
- 323 • Data analysis software: Figures 1-5 are produced in Python and the details of the method-
324 ology is provided in the relevant sections of the text. Figure 5 is produced using script avail-
325 able at https://github.com/bidyutbg/CESM_Tau_experiment/blob/main/WK_spectra_FINAL-NEW.ipynb. Figure 6 is produced using script available at https://github.com/bidyutbg/CESM_Tau_experiment/blob/main/CCEW_variance-compare_FINAL.ipynb. Figure 7 is produced using script available at https://www.ncl.ucar.edu/Applications/Scripts/mjoclivar_9.ncl.
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329
- 330 • Model Output Data: Data archival is underway in Zenodo. Archival will be completed
331 soon. A sample of the data is provided as Supporting Information for review purposes.

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336

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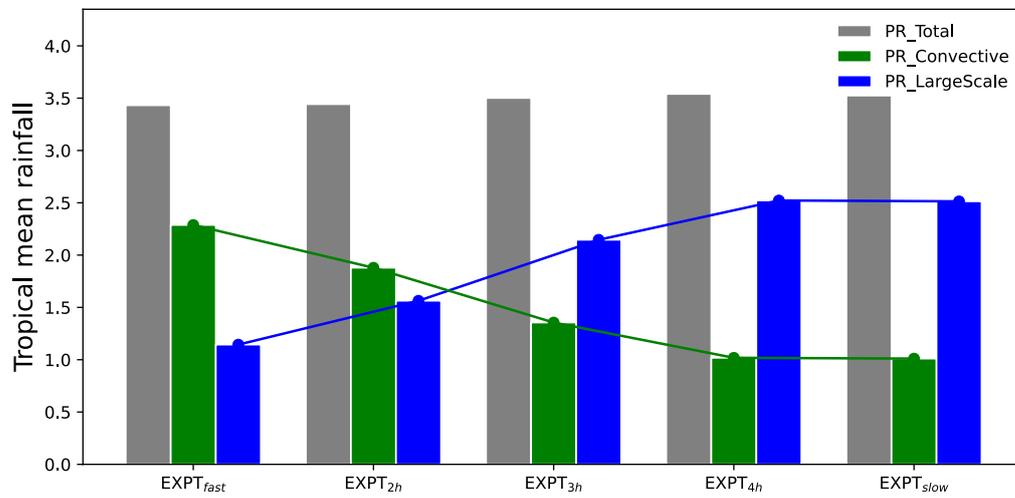
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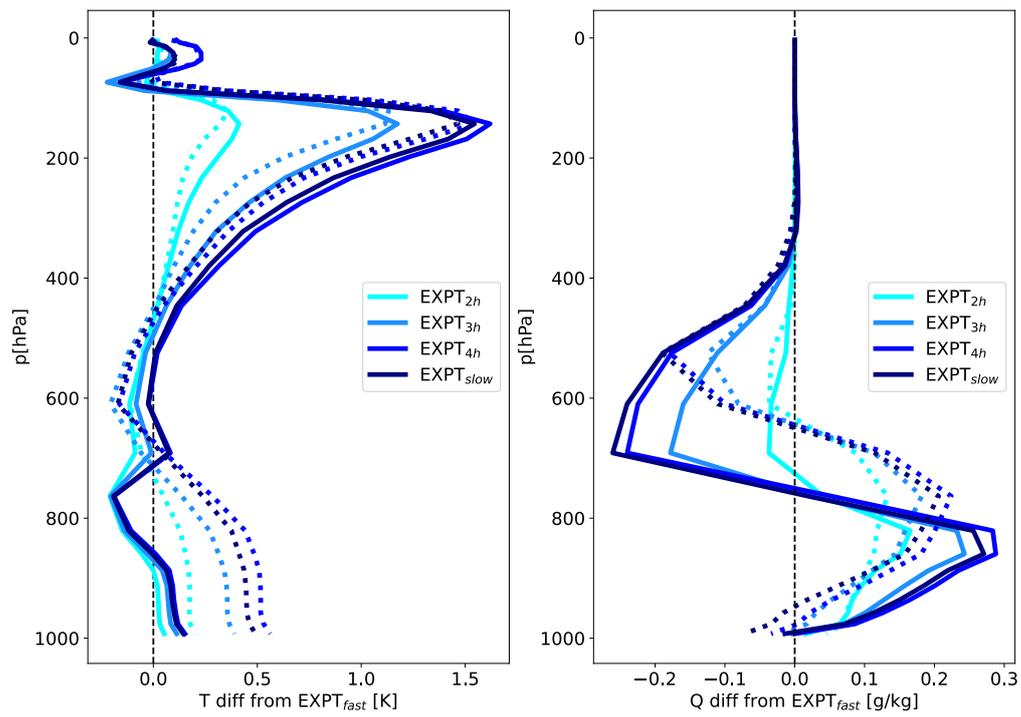
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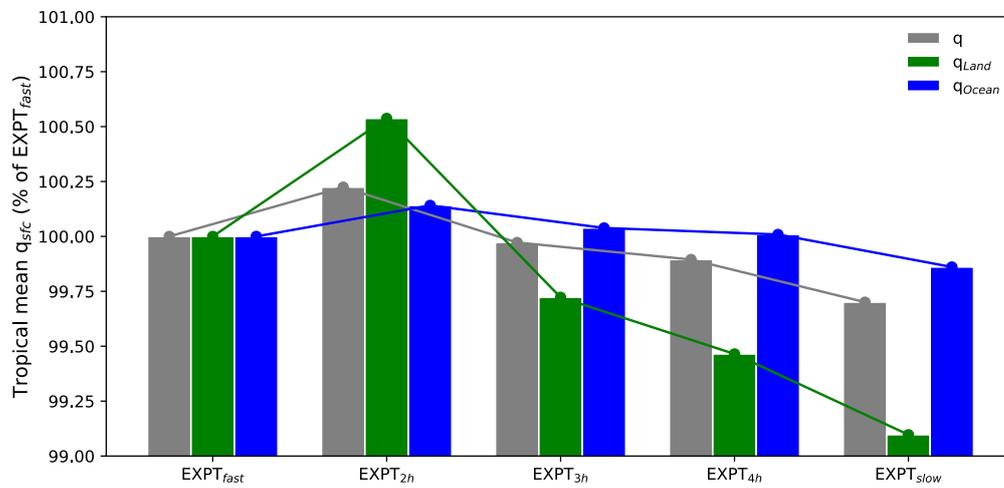
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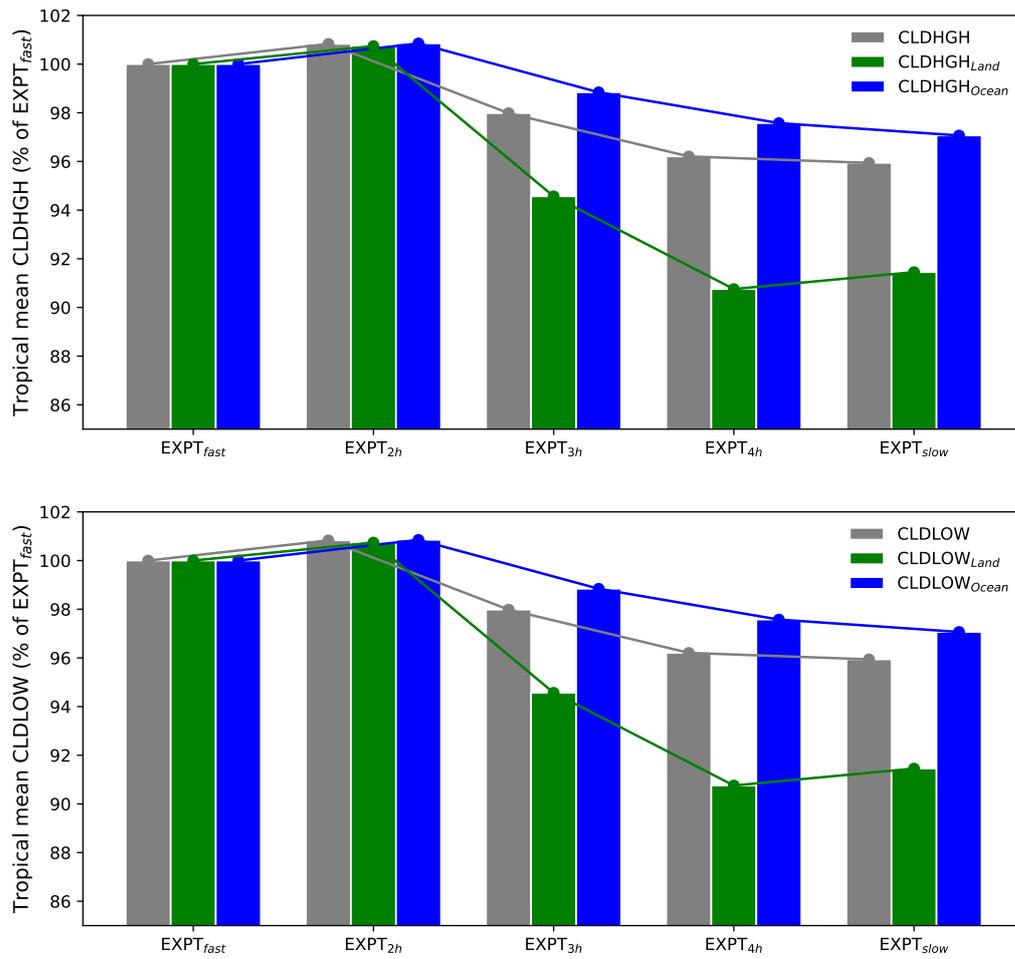
472 **Figure 1.** Tropical (tropics defined as the zonal belt between 30°S-30°N) annual mean daily rainfall
 473 (mm/day) for different experiments mentioned in Table 1.



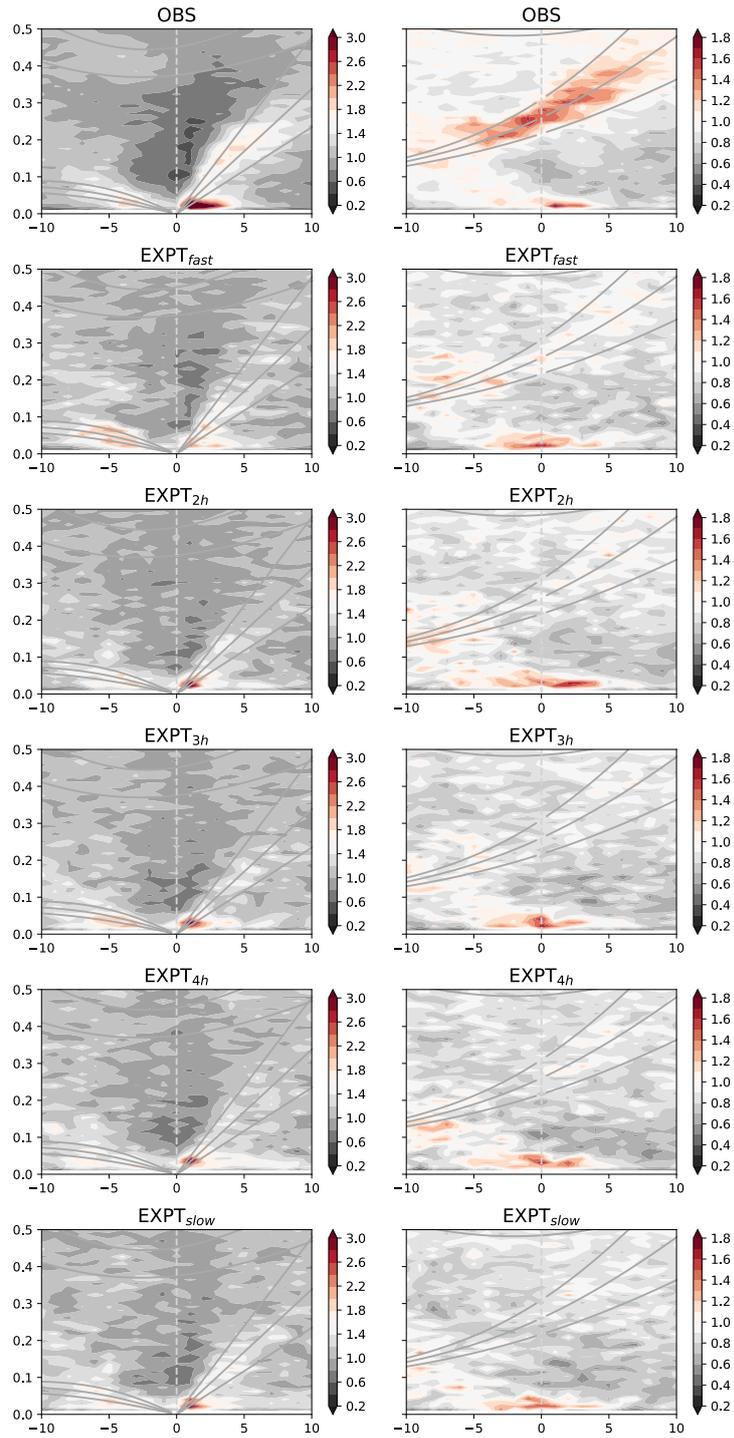
474 **Figure 2.** Tropical (tropics defined as the zonal belt between 30°S-30°N) mean vertical profiles of tem-
 475 perature (T) and specific humidity (Q). Departures of different experiments, as indicated in the legends, from
 476 $EXPT_{fast}$ (Land: Dotted, Ocean: Solid). The vertical dashed line indicate the zero line.



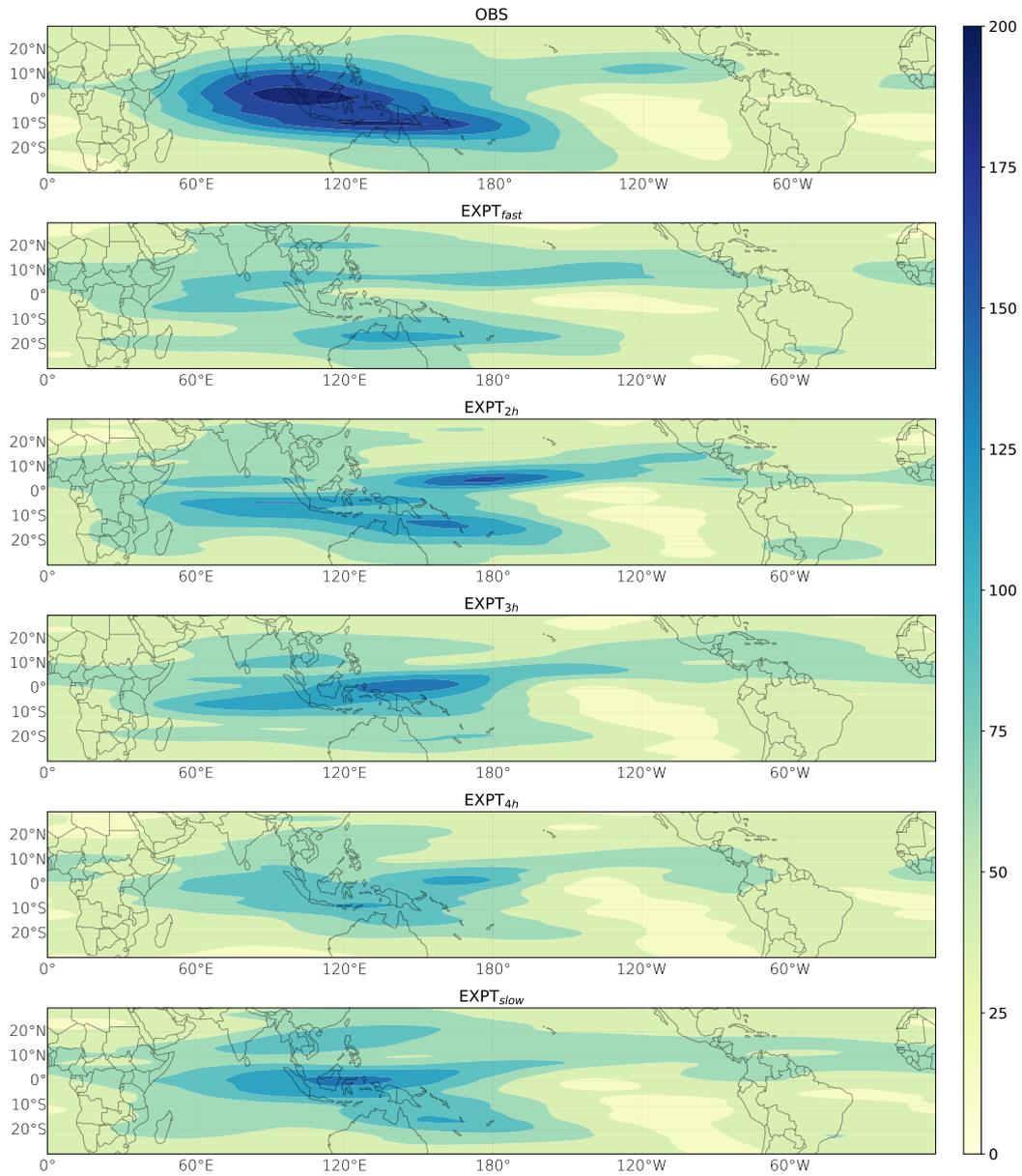
477 **Figure 3.** Tropical (tropics defined as the zonal belt between 30°S-30°N) annual daily mean specific hu-
 478 midity as surface depicted as % of $EXPT_{fast}$.



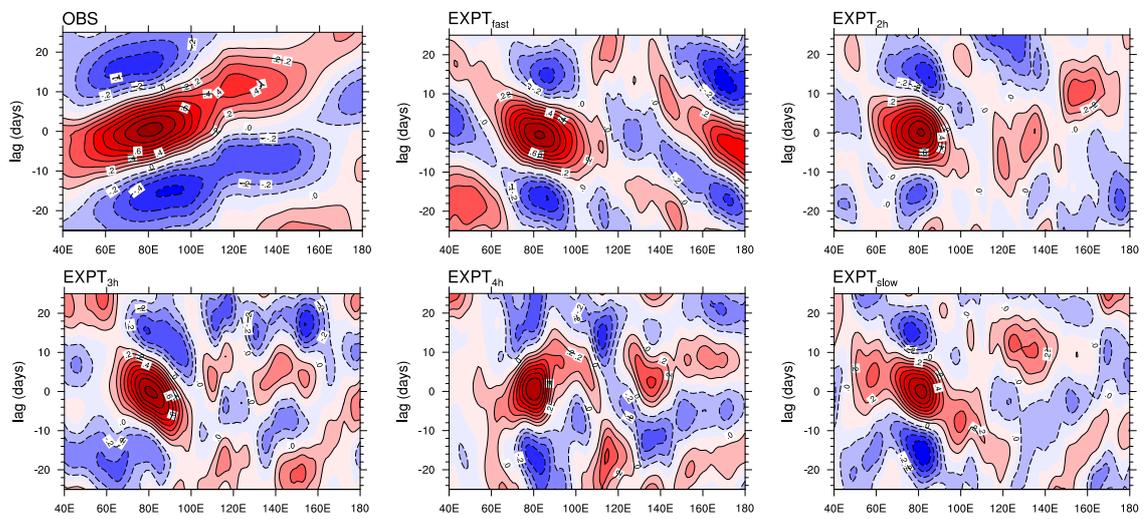
479 **Figure 4.** Tropical (tropics defined as the zonal belt between 30°S-30°N) annual daily mean High and Low
 480 cloud cover depicted as % of $EXPT_{fast}$.



481 **Figure 5.** Takayabu-Wheeler-Kiladis spectra of OLR for OBS (from NOAA) and different experiments (as
 482 named above each panel), for the symmetric component (left-hand side panels) and antisymmetric component
 483 (right-hand side panels).



484 **Figure 6.** MJO variance computed as the daily variance of OLR data filtered for 1-5 wavenumber and 20-
 485 100 day frequency, for OBS (from NOAA) and different experiments (as named above each panel).



486 **Figure 7.** MJO propagation: Hovmoller (averaged from 10°S to 10°N) plots of MJO-filtered OLR (W m⁻²)
487 anomalies (Winter), for OBS (from NOAA) and different experiments (as named above each panel).

1 **An assessment of representing land-ocean heterogeneity via**
2 **convective adjustment timescale in the Community Atmospheric**
3 **Model 6 (CAM6)**

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6 **Key Points:**

- 7 • Two distinct values of convective adjustment timescale, τ , over land & ocean in the con-
- 8 vective parameterization scheme are prescribed.
- 9 • The mean climate stays qualitatively the same, except for a moister and colder near-surface
- 10 atmosphere for longer τ s over the oceans.
- 11 • A primary gain of using two different τ s for land and ocean is improved simulation of the
- 12 convectively coupled equatorial waves.

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Abstract

The time needed by deep convection to bring the atmosphere back to equilibrium is called convective adjustment timescale or simply adjustment timescale, typically denoted by τ . In the Community Atmospheric Model (CAM), convective adjustment timescale is a tunable parameter with one value, 1 hour, worldwide. Albeit, there is no justified reason why one adjustment timescale value should work over land and ocean both. Continental and oceanic convection are different in terms of the vigor of updrafts and hence can have different longevities. So it is logical to investigate the prescription of two different convective adjustment timescales for land (τ_L) and ocean (τ_O). To understand the impact of representing land-ocean heterogeneity via τ , we investigate CAM climate simulations for two different convective adjustment timescales for land and ocean in contrast to having one value globally.

Following a comparative analysis of 5-year-long climate simulations, we find $\tau_O = 4$ hrs and $\tau_L = 1$ hr to yield the best results. Particularly, we find better MJO simulations. Although these τ values were chosen empirically and require further tuning, the conclusion of our finding remains the same, which is the recommendation to use two different τ values for land and ocean.

1 Introduction

Deep convection is complex to parameterize [Arakawa, 2004]. While the explicit representation of deep convection is becoming a plausible option to navigate this "deadlock" [Randall *et al.*, 2003; Randall, 2013], for long-term projections of our climate, cumulus parameterization is still unavoidable. Hence, amidst the fierce emergence of convection-resolving models [Stevens *et al.*, 2019], various schemes to parameterize convection continue to develop. In particular, the recent decades have witnessed a surge of novel ideas that have accelerated this progress [Rio *et al.*, 2019, and references therein].

The "art" of tuning parameters used in convection parameterization schemes, or simply parameter tuning, plays a vital role in this development process [Hourdin *et al.*, 2017]. While deficiencies of convective parameterization are primary factors for model biases, it alone cannot alleviate all mode biases [Goswami *et al.*, 2017]. Hence, parameter sensitivity investigations are necessary not only to optimize the performance of a scheme but also to understand the extremities to which a scheme can be held responsible for biases in a simulation [Qian *et al.*, 2015; Goswami *et al.*, 2017]. In this study, we aim to contribute to understanding one tunable parameter, the convective adjustment timescale τ , by investigating the sensitivity of climate simulations to two dif-

44 ferent τ values for land and ocean in contrast to having one value globally in the Zhang-McFarlane
45 convective parameterization scheme [Zhang and McFarlane, 1995, ZM95 hereafter] in the Com-
46 munity Atmospheric Model (CAM), the atmospheric model of the Community Earth System Model
47 [Danabasoglu et al., 2020].

48 In CAM, deep convection is represented using the Zhang-McFarlane (ZM) convection pa-
49 rameterization scheme. The ZM is an adjustment-type convective parameterization scheme where
50 the atmospheric instability is removed via an adjustment towards a background state. In ZM, con-
51 vective available potential energy (CAPE) defines atmospheric instability, and τ is the CAPE con-
52 sumption time. In their paper, ZM95 used τ values of 2, 4, and 6 hours. To quote ZM95, "The
53 adjustment time scale determines the intensity and duration of convection for a given CAPE. With
54 small τ the convection is short-lived but intensity is high, on the other hand with larger τ the con-
55 vection is long-lived but of low intensity". ZM95 reported their scheme to be particularly sen-
56 sitive to the choice of τ . Since there is no strict range of τ , several studies investigated the sen-
57 sitivity of CAM simulations to different τ values. For example, Mishra and Srinivasan [2010]
58 used $\tau=[1,\infty]$. Contrasting water-vapor isotope simulations in a suite of CAM single-column sim-
59 ulations with a range of τ values, Lee et al. [2009] found their simulations to match better with
60 satellite observations with $\tau = 8$ hrs. Mishra [2011, 2012] prescribed $\tau = 8$ hrs in global climate
61 simulations and noted improvements in the simulations of tropical climate, especially the con-
62 vectively coupled equatorial waves. Evaluating 22 tunable parameters in CAM, Qian et al. [2015]
63 reported τ as one of the most critical tuning parameters. In all of the above studies, τ has a sin-
64 gular value globally.

65 One value of τ globally is not a logical choice because deep convection exhibits different
66 behaviors over continents and oceans [Hagos et al., 2013; Matsui et al., 2016; Roca et al., 2017;
67 Roca and Fiolleau, 2020]. Since the width of a thermal plume is steered by boundary layer height
68 [Williams and Stanfill, 2002], a deep continental boundary layer generates wider updraft veloc-
69 ities in deep convection [Lucas et al., 1994]. Matsui et al. [2016] provided a climatological view
70 of the contrast between oceanic and continental convective precipitating clouds from long-term
71 TRMM satellite multisensor statistics. They found large proportions of deep clouds over land.
72 Zipser et al. [2006] also found the most intense storms typically over continents. These obser-
73 vations suggest that the atmospheric deep convection over land is wider and stronger than those
74 over the oceans. In other words, atmospheric convection over land is shorter lived than that over
75 ocean [Roca et al., 2017]. It advocates for a shorter convection consumption time scale over land
76 than over oceans which motivated us to address the following question: although two different

77 τ values incorporating land-ocean inhomogeneity are logical, is it fruit-bearing in a model-simulated
78 climate? To answer this question, we investigate,

- 79 • response of the mean climate, and
- 80 • response of large-scale waves,

81 by contrasting 5-year-long climate simulations with and without incorporating land-ocean inho-
82 mogeneity via τ values.

83 Convective parameterization schemes, particularly adjustment-type schemes, are based on
84 the idea that convection takes some time to stabilize the atmosphere to a background state. Es-
85 sentially, this time taken is τ in the ZM scheme. Although numerically τ can have almost any value,
86 it is decided based on a scale separation between the convective activity of the individual clouds
87 and large-scale forcing. This concept is nicely depicted in Figure 1.1 of [Davies, 2008]. The graph
88 in that figure is a function of timescales associated with convection, and consists of a turbulent
89 initial segment indicating fluctuation of individual clouds, followed by a flat segment where these
90 fluctuations smooth out, and finally a segment corresponding to longer time-scales that shows
91 the evolution of the large scale forcing field itself. Conceptually, changing τ within a reasonable
92 range (within the flat segment of Figure 1.1 of [Davies, 2008]) should not result in a dramatic change
93 in the mean state of the simulated climate. We shall investigate it in detail in the first part of our
94 results section.

95 Some changes that we expect in our experiments are in the simulated organization of con-
96 vection. The organization of convection comes from the dynamic and thermodynamic impacts
97 of convection on the atmosphere. Simply put, it is the memory of convection [Davies *et al.*, 2009],
98 i.e. the fact that convection changes the large-scale properties, and can make their environment
99 favorable or unfavorable to subsequent convection. Identifying sources of convective memory
100 in cloud-resolving simulations, Colin *et al.* [2019] argued that the persistence of the state of con-
101 vection contributes to convective memory. Colin *et al.* [2019] also suggested that convective mem-
102 ory and organization interact mutually. By altering τ we essentially alter memory associated with
103 convection. Hence, we expect to see changes in convective organization. Taking a cue from Mishra
104 [2011], we anticipate improved convective organization in the tropics for longer τ . However, land-
105 ocean heterogeneity in τ is a unique feature of our experiments that we argue is essential based
106 on heterogeneity in the behavior of convection over land and ocean. As supporting evidence, we

107 shall present an analysis of equatorial waves focusing on the MJO to evaluate the organization
 108 of convection in the second part of our results section.

109 The paper is organized as follows. A brief description of the methodology is provided in
 110 Section 2. Section 3 evaluates the response of the model to different τ values. Finally, a few con-
 111 cluding remarks are provided in Section 4.

112 **2 Model and simulation details**

113 We used the atmospheric model of the Community Earth System Model, version 2.1.3 (CESM
 114 2.1.3) [Danabasoglu *et al.*, 2020], that is the Community Atmosphere Model, version 6 (CAM6),
 115 developed and maintained at the National Center for Atmospheric Research (NCAR), with lon-
 116 gitude and latitude specifications 1.25° and 0.9° , respectively, and 32 vertical levels. We forced
 117 the model by HadISST1 climatological monthly mean SST data provided by the Met Office Hadley
 118 Centre [Rayner, 2003]. In short, we performed CESM “F2000climo” simulations. In general,
 119 these are atmospheric simulations forced by present-day climatology. All simulations are 6 years
 120 long, and we analyzed the last 5 years of each simulation since, for atmosphere-only simulations,
 121 1-year spin-up is enough.

122 We performed 5 simulations. The one with out-of-the-box τ value of 1 hour globally is called
 123 the control (*CTRL*). In the next 3 simulations, we delayed the τ value over ocean (τ_O) to 2, 3 and
 124 4 hours keeping τ over land (τ_L) 1 hour. We called these 3 simulations *EXPT_{2h}*, *EXPT_{3h}* and
 125 *EXPT_{4h}*, respectively. We performed a last 5th experiment, named *EXPT_{slow}*, for which we used
 126 a τ value of 4 hours globally. Before starting our comparative analysis, we rename our first sim-
 127 ulation as *EXPT_{fast}*, which initially we had named CTRL, for clarity and better fluency of nar-
 128 ration of our findings. Table 1 depicts the τ values for different experiments.

129 Our analyses primarily show a comparison between the 5 aforementioned simulations. For
 130 some analyses we have used outgoing long-wave radiation (OLR) from NOAA ($2.5^\circ \times 2.5^\circ$; daily
 131 from 01-Jun-1974 to 12-Dec-2019) [Liebmann and Smith, 1996] as observational benchmark.

133 **3 Results**

134 **3.1 Mean Climate**

135 Since about 75% of the global surface is ocean, in the simulations of the mean climate, we
 136 expect a similar model response in our experiments by delaying τ only over the oceans, as ear-

Experiment Name	τ_L	τ_O
<i>EXPT_{fast}</i>	1hr	1hr
<i>EXPT_{2h}</i>	1hr	2hr
<i>EXPT_{3h}</i>	1hr	3hr
<i>EXPT_{4h}</i>	1hr	4hr
<i>EXPT_{slow}</i>	4hr	4hr

Table 1. τ values for different experiments

132

137 lier studies did by having a larger τ globally. An evaluation of some of the mean features of sim-
 138 ulated climate in our experiments confirm this. We find an increase in large-scale rainfall and a
 139 decrease in convective rain going from *EXPT_{fast}* to *EXPT_{slow}* (Fig 1 and Supplementary Fig
 140 S1). Similarly, we also notice warming in the lower levels, stronger warming in the upper lev-
 141 els, slight cooling in the mid-levels; moistening in the lower levels, and drying in the mid-levels
 142 (Fig 2 and Supplementary Fig S2). These features have been reported in earlier studies [for ex-
 143 ample, Fig 8 in *Mishra and Srinivasan, 2010*].

144 Investigating the mean features for land and ocean separately, we notice in addition, lower
 145 level (upper level) warming (cooling) is more (less) over land than over oceans (Fig 2). In the case
 146 of moisture, the letter "S" patterned vertical structure over the ocean is more curvy and squeezed
 147 down meaning lower level (middle level) moistening (drying) is stronger over oceans than over
 148 land and the respective peaks are vertically closer to the sea surface. These profiles, all together,
 149 indicate a model response to changes in τ in terms of the distribution of atmospheric convection
 150 and clouds, which impacts heating/cooling and moistening/drying of the air column (Supplemen-
 151 tary Fig S2). Essentially these responses indicate an accumulation of convective instability in the
 152 atmosphere with delaying of convective adjustment time scale. It is attributable to more low-level
 153 warming over the continents and more low-level moistening over the oceans. More moistening
 154 near the ocean surface is relatively straightforwardly understandable, and it is a consequence of
 155 the atmosphere taking longer to convect with larger τ . To a zero-order approximation, as a re-
 156 sult of the near-surface moisture pile-up in the oceanic regions, there is a moisture deficit in the
 157 lower levels over the continental regions (Fig 3 and Supplementary Fig S3a and S3b). Indeed it
 158 is apparent, in relative sense, in Fig 3. Although q_O does not exhibit a clear moistening signal,

159 the land drying in q_L is profound. The consequences are reflected in terms of changes in cloud
 160 cover. In an overall declining tendency of cloud cover, from $EXPT_{fast}$ to $EXPT_{slow}$, over the
 161 tropics high clouds decrease more steeply than low clouds. Low clouds decrease less rapidly over
 162 the ocean compared to those over land (Fig 4). It should be noted that cloud categories are ob-
 163 jectively defined in CESM. For example, low-level clouds are the ones below 700 hPa and high
 164 clouds are between 400 and 50 hPa. Cloud covers are integrated for each model level correspond-
 165 ing to respective cloud categories. In that regard, going from $EXPT_{fast}$ to $EXPT_{slow}$, low-cloud
 166 cover changes (Fig 4) are consistent with relative surface moistness over land and ocean (Fig 3).

167 Taken together, the altered vertical profiles of moisture and temperature, distribution of con-
 168 vective and large-scale rainfall, and associated clouds are consistent with the idea that convec-
 169 tion is short-lived and stronger for smaller τ values and long-lived and weaker for longer τ value.
 170 It is also evident from the solution of the CAPE equation in the ZM scheme, which can be ex-
 171 pressed as $CAPE(t) = CAPE_o \exp(-\frac{t}{\tau})$ in the absence of large-scale CAPE generation, where
 172 $CAPE_o$ is the values of CAPE at $t = 0$. A larger τ in this expression means a slower decay of
 173 CAPE. The duration of convection is essentially linked with its persistence and hence "mem-
 174 ory". We discuss its impact on the simulation of the equatorial waves in the following section.

175 3.2 Simulation of MJO variance and propagation

176 Organization is a primary feature of tropical convection. It essentially means a cluster of
 177 deep precipitating clouds tied together. An important question is, what brings these clouds to-
 178 gether? In other words, what causes convection to organize? One idea to see the organization of
 179 convection is through superpositions of convectively coupled equatorial waves (CCEWs). These
 180 atmospheric waves and tropical convection are entangled. In the tropics, the atmosphere responds
 181 to convective heating in terms of waves that, in turn, organize convection. Therefore, the fidelity
 182 of a model in simulating tropical climate is essentially its ability to simulate the CCEWs. A stan-
 183 dard metric to analyze CCEWs is the Takayabu-Wheeler-Kiladis (TWK) spectra [*Takayabu, 1994a,b;*
 184 *Wheeler and Kiladis, 1999*]. Figure 5 depicts the symmetric and asymmetric TWK-spectra for
 185 the observed and simulated outgoing long-wave radiation (OLR). Understandably, a striking fea-
 186 ture of the TWK-spectra of observed OLR shown in Fig 5a and b is the spectral power near the
 187 origin of the plots in the wavenumber range 1-5 and frequency 20-100 days, well known as the
 188 MJO. The MJO is a combination of or envelope of other waves in the equatorial atmosphere. Hence,
 189 the accuracy of MJO simulation is arguably a measure of the fidelity of accurate simulation of

190 waves in the atmosphere [Zhang *et al.*, 2020]. Guo *et al.* [2015] showed in detail that the accu-
 191 racy of CCEW simulation is critical for a realistic MJO simulation.

192 A comprehensive review of the science of MJO is available in Zhang *et al.* [2020]. Promi-
 193 nent observed features of MJO suggest that they are most active in the Indo-Pacific warm pool
 194 with an eastward propagation. An interesting fact, along its path from the Indian to the Pacific
 195 Ocean, is that an MJO passes over the Indonesian maritime continent (IMC). During this pas-
 196 sage, MJO and the prominent diurnal variabilities in the meteorology over the IMC islands in-
 197 teract and mutually influence each other. So much so that nearly half of the MJOs fail to prop-
 198 agate into the Pacific. It is critical, therefore, to represent the land-ocean heterogeneity as real-
 199 istically as possible in climate models. Hence, we expect our experiments with logically defined
 200 different values of τ for land and ocean to improve simulated MJO features. Here, we shall present
 201 analyses evaluating the simulation of MJO variance and propagation. We can draw some idea
 202 of MJO simulation in different experiments from Fig 5. In Fig 5, the foremost remarkable fea-
 203 ture is the increase in spectral power in the MJO wave number and frequency range for experi-
 204 ments with a longer τ . A closer visual inspection reveals that the MJO spectral power does not
 205 dramatically change from $EXPT_{2h}$ to $EXPT_{slow}$. For other waves, no one simulation is remark-
 206 ably better than the rest. Fig 5 loosely suggests that overall the symmetric signal waves are im-
 207 proved for longer time scales, but there are no clear improvement for the antisymmetric part.

208 To bring out the active region of MJO we applied space-time filtering on OLR data con-
 209 taining the signal corresponding to wavenumbers 1-5 and a period of 20–100 days. In Fig 6 the
 210 variance of the MJO-filtered daily OLR anomalies is shown. In observations (Fig 6a), the peak
 211 variance is over the Indo-Pacific warm pool. Feeble variance peaks are noted in the eastern sides
 212 of the Pacific (off the Gulf of California) and Atlantic (around the western coast of Sierra Leone).
 213 It is consistent with the fact that although MJO is most active in the Indo-Pacific warm pool re-
 214 gion, it has considerable influence modulating the convective activity over the eastern equato-
 215 rial Pacific [Maloney and Hartmann, 2000a,b; Maloney and Kiehl, 2002] and Atlantic [Klotzbach,
 216 2014]. For $EXPT_{fast}$ high variance is noted around the warm-pool region but widely spread and
 217 has multiple peaks. The strongest variance is around Northern Australia and the south-western
 218 Pacific region. The other secondary maxima are over the southern Bay of Bengal, the central equa-
 219 torial Indian Ocean, and the central Pacific regions.

220 The simulated MJO variance strength and pattern experience some changes with changes
 221 in τ values. In general, a slower τ_O keeping τ_L same yields more variance. In other words, it in-

222 creases convective activity in MJO space and time scales. In $EXPT_{2h}$ a pronounced peak is
 223 located over the western-central equatorial Pacific with two secondary maxima near the south-western
 224 equatorial Pacific and eastern equatorial Indian Ocean. In $EXPT_{3h}$ the variance is more concen-
 225 trated over the western equatorial Pacific, with a secondary peak south of the central equatorial
 226 Indian Ocean. With larger values of τ_L , the maximum variance gets more and more focused over
 227 the warm pool region, from $EXPT_{fast}$ to $EXPT_{3h}$ (comparing Fig 6b-d). It is noteworthy, that
 228 all the pronounced peaks for $EXPT_{2h}$ and $EXPT_{3h}$ are over oceans, in and around the Indo-Pacific
 229 warm pool region, but split unlike observations (Fig 6a). The model simulated MJO variance fur-
 230 ther slowing τ_O to 4 hours ($EXPT_{4h}$ shown in Fig 6e) suggests that MJO variance does not nec-
 231 essarily increase with increasing τ_O . The variance peak intensities are visibly weaker in $EXPT_{04}$
 232 compared to that in $EXPT_{2h}$ and $EXPT_{3h}$ and more only than that in $EXPT_{fast}$. However, a note-
 233 worthy feature of $EXPT_{4h}$, a fine detail missing in all other simulations, is the variance peaks
 234 near the eastern side of the equatorial Pacific and Atlantic oceans. Baring these subtle variance
 235 peaks, $EXPT_{slow}$ looks the best, although still a considerably weaker variance peak compared
 236 to observations. The variance fields normalized by the respective domain means are available
 237 in Supplementary Fig S4, which depicts a better visual illustration of the variance peaks.

238 A prominent feature of MJOs is eastward propagation. The propagation features of the MJO
 239 are arguably better characterized by Hovmöller plots averaged over the latitude band between 10°S
 240 and 10°N , shown in Fig 7. Each frame in Fig 7 depicts 10°S - 10°N averaged cross-correlations
 241 of OLR anomalies with MJO-index. The MJO-index is defined as the 20-100-day filtered OLR
 242 anomalies averaged over 5°S - 5°N , 75°E - 85°E following *Guo et al. [2015]*. It is noteworthy to
 243 mention, reiterating *Guo et al. [2015]*, the philosophy behind using such an MJO index. An in-
 244 dex based on a 20-100 day filter brings out the dominant intraseasonal signal in the data that ide-
 245 ally should be an MJO signal. The eastward propagating red and blue patches of correlation val-
 246 ues in observations (Fig 7a) confirm it. We note the phase speed is faster over the west Pacific
 247 (east of $\sim 120^\circ\text{E}$) than that over the Indian Ocean (west of $\sim 100^\circ\text{E}$). The relatively slow phase speed
 248 in the longitude range $\sim 100^\circ$ - 120°E is collocated with the Indonesian archipelago. These dif-
 249 ferent phase speeds over land and oceanic regions are consistent with MJO interaction with the
 250 profound diurnal variations of meteorology over the MC. It furthermore emphasizes the need to
 251 mimic land-ocean heterogeneity realistically in climate models.

252 To assess the performance of our different experiments in simulating MJO propagation fea-
 253 tures, we recall the "good" and "bad" models of *Guo et al. [2015]*. In Figure 2, *Guo et al. [2015]*
 254 showed that the "good" models simulated more realistic eastward propagation than the "bad" mod-

255 els. In Fig 7, $EXPT_{4h}$ is the only experiment with an eastward propagation and exhibits some
 256 resemblance with observations and the only "good" model, albeit with some key caveats. The
 257 positive anomalies almost abruptly died over the MC and reappeared over the western Pacific.
 258 Nonetheless, an intriguing observation, that contains the novelty of our research, is the more re-
 259 alistic eastward propagation simulated in $EXPT_{4h}$ than in $EXPT_{slow}$. An improved simulation
 260 of eastward propagation in $EXPT_{4h}$ supports our argument that using two τ s for land and ocean
 261 is a logical choice. It reconfirms our anticipation that representing land-ocean heterogeneity via
 262 τ in ZM in CAM alters convective memory and affects the organization of convection. A larger
 263 τ_O than τ_L , although reasonable, is only based on intuition. Detailed sensitivity analysis would
 264 be needed to investigate and pin down the best pair of τ values.

265 **4 Discussion and Conclusion**

266 Climate models continue to grow, fueled by a growing understanding of the earth system.
 267 Hence, it is only logical to include a fairly well-recognized and relatively old knowledge about
 268 land and ocean heterogeneity of atmospheric convection in the parameterization of convection.
 269 We argue that using two different τ in ZM in CAM can be one simple yet fruit-bearing way. In
 270 our experiments to investigate the model response to land-ocean heterogeneity in τ values, we
 271 used $\tau_L = 1$ hr, and $\tau_O = 2$ hrs, 3 hrs, 4 hrs. In two additional experiments, $EXPT_{fast}$ and $EXPT_{slow}$,
 272 we used $\tau_L = \tau_O = 1$ hr and $\tau_L = \tau_O = 4$ hrs, respectively, to complement the previous group
 273 of experiments. The τ values that we have used are informed by our knowledge of frequency, life-
 274 cycle, and behavior of atmospheric convection over land and ocean learned from previous stud-
 275 ies [Lucas *et al.*, 1994; Williams and Stanfill, 2002; Zipser *et al.*, 2006; Hagos *et al.*, 2013; Mat-
 276 sui *et al.*, 2016; Roca *et al.*, 2017; Roca and Fiolleau, 2020] and inspired by results of relevant
 277 model sensitivity experiments [Zhang and McFarlane, 1995; Lee *et al.*, 2009; Mishra and Srini-
 278 vasan, 2010; Mishra, 2011; Misra *et al.*, 2012].

279 Our findings regarding the model simulated mean state in different experiments are con-
 280 sistent with earlier studies [Lee *et al.*, 2009; Mishra and Srinivasan, 2010; Mishra, 2011; Misra
 281 *et al.*, 2012]. For example, total rainfall remained approximately the same while large-scale rain-
 282 fall increased and convective rain decreased for longer τ_L s. Consistency of the model response
 283 for a slow τ only over the oceans with slowing down τ globally is most likely a result of 75% of
 284 the global surface being ocean. However, since there is no physical barrier between the atmospheric
 285 columns over continents and oceans, having two τ values in our experiments, which essentially
 286 are prescribed to represent heterogeneity in the persistence of convection over the two different

287 surfaces, created a distinction between the intensities with which the model responses are felt over
288 land and ocean. For example, the oceanic boundary layer is moister and warmer than the con-
289 tinental boundary layer (Fig 3). Furthermore, the mid-troposphere is drier and cooler over oceans
290 than over the continents (Fig 2). These land-ocean heterogeneities inevitably create differences
291 in atmospheric instabilities. These instabilities are essentially realized in the form of atmospheric
292 convection that, by design in our experiments with slower τ , takes longer to bring the atmosphere
293 back to a background state. It is suggestive of a longer persistence of convective instability over
294 the ocean than that over the continents which essentially can be linked with memory of convec-
295 tion [Davies *et al.*, 2009; Colin *et al.*, 2019; Hwong *et al.*, 2023].

296 The conclusion that the model simulated better convectively coupled equatorial waves in
297 $EXPT_{2h}$ than in $EXPT_{slow}$ is a key. We conclude this based on our finding of a better MJO sim-
298 ulation in $EXPT_{2h}$, consistent with improved symmetric waves. Scientists had advocated in favor
299 of a slower τ in earlier studies [Mishra, 2011; Misra *et al.*, 2012]. We also noted a signifi-
300 cant increase in MJO power for $\tau = 4$ hrs than $\tau = 1$ hr (comparing Fig 5b and Fig 5f). However,
301 an evaluation of the model simulated intraseasonal zonal propagation reveals that $EXPT_{4h}$ per-
302 forms considerably better than $EXPT_{slow}$. This confirms that having one τ globally is not only
303 unphysical but also slowing down tinkering persistence of convection to improve simulation of
304 equatorial waves, and may result in model responses that might look improved, but only super-
305 ficially.

306 Our results, in general, serve as proof of concept that a realistic representation of convec-
307 tive adjustment time scale over land and ocean is a logical requirement that properly implemented
308 shall lead to improvements in climate model simulations. In specific, we advocate at least two
309 τ values, one for the continents and one relatively slower for the oceans in ZM in CAM. The fact
310 that we did not perform a rigorous model sensitivity analysis [e.g., Qian *et al.*, 2015; Lin *et al.*,
311 2016; Goswami *et al.*, 2017] nor did we perform any cloud-resolving simulation targeting the
312 life-cycle of atmospheric convection [Davies *et al.*, 2013; Colin *et al.*, 2019; Daleu *et al.*, 2020,
313 e.g.,] leaves a scope as well as the requirement for future research to determine the best values
314 of τ_L and τ_O for ZM in CAM. It will hopefully guide convection parameterization schemes, es-
315 pecially the adjustment types, to address land-ocean heterogeneity. Specifically, we recommend
316 that future developments of CAM should consider prescribing different τ_L and τ_O in ZM in CAM.

5 Open Research

- 318 • Model : We used the atmospheric model of the Community Earth System Model, version
319 2.1.3 (CESM 2.1.3) [Danabasoglu *et al.*, 2020]
- 320 • Description of the model simulations is provided in Section 2 of the manuscript. A source
321 file of CESM 2.1.3, zm_conv.F90, modified for our experiments is provided in https://github.com/bidyutbg/CESM_Tau_experiment.git.
322
- 323 • Data analysis software: Figures 1-5 are produced in Python and the details of the method-
324 ology is provided in the relevant sections of the text. Figure 5 is produced using script avail-
325 able at https://github.com/bidyutbg/CESM_Tau_experiment/blob/main/WK_spectra_FINAL-NEW.ipynb. Figure 6 is produced using script available at https://github.com/bidyutbg/CESM_Tau_experiment/blob/main/CCEW_variance-compare_FINAL.ipynb. Figure 7 is produced using script available at https://www.ncl.ucar.edu/Applications/Scripts/mjoclivar_9.ncl.
326
327
328
329
- 330 • Model Output Data: Data archival is underway in Zenodo. Archival will be completed
331 soon. A sample of the data is provided as Supporting Information for review purposes.

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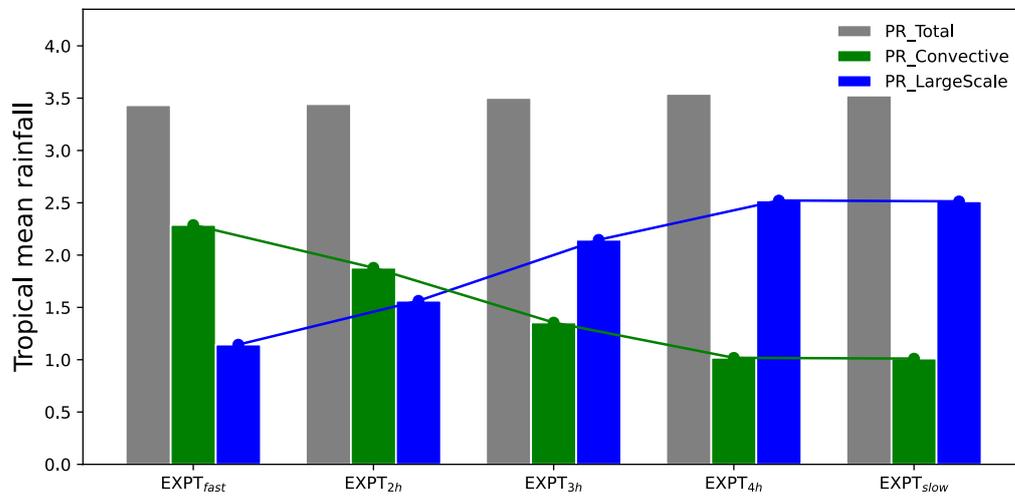
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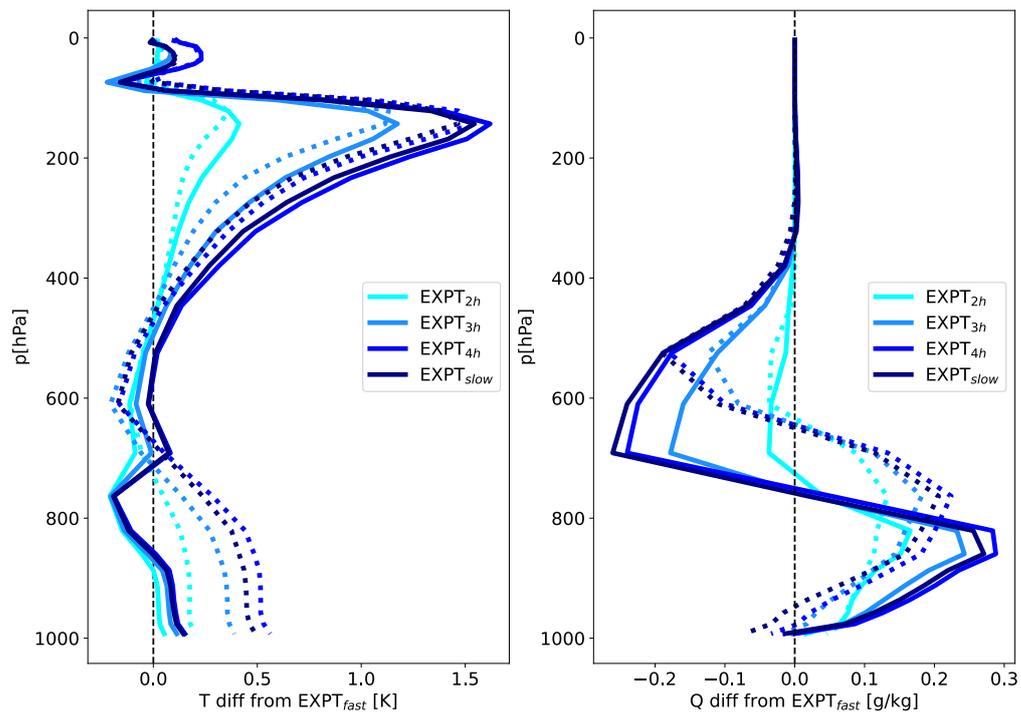
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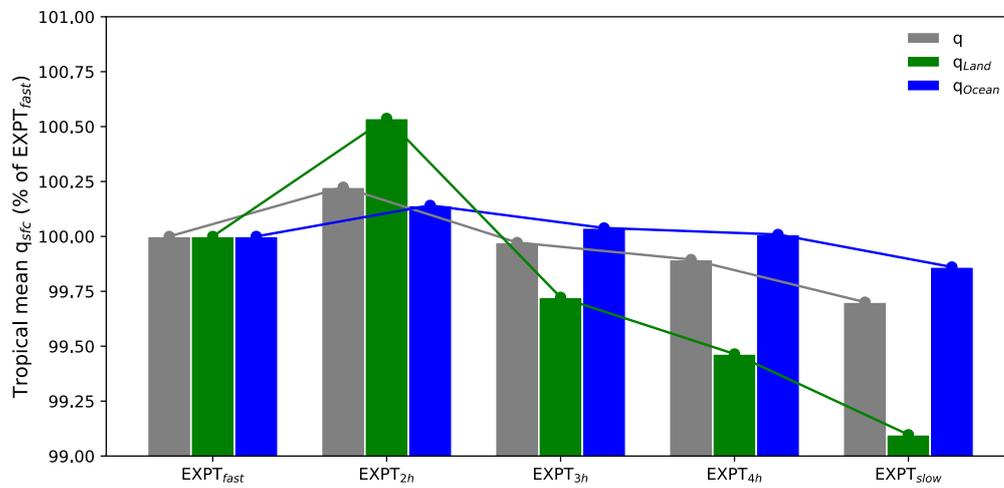
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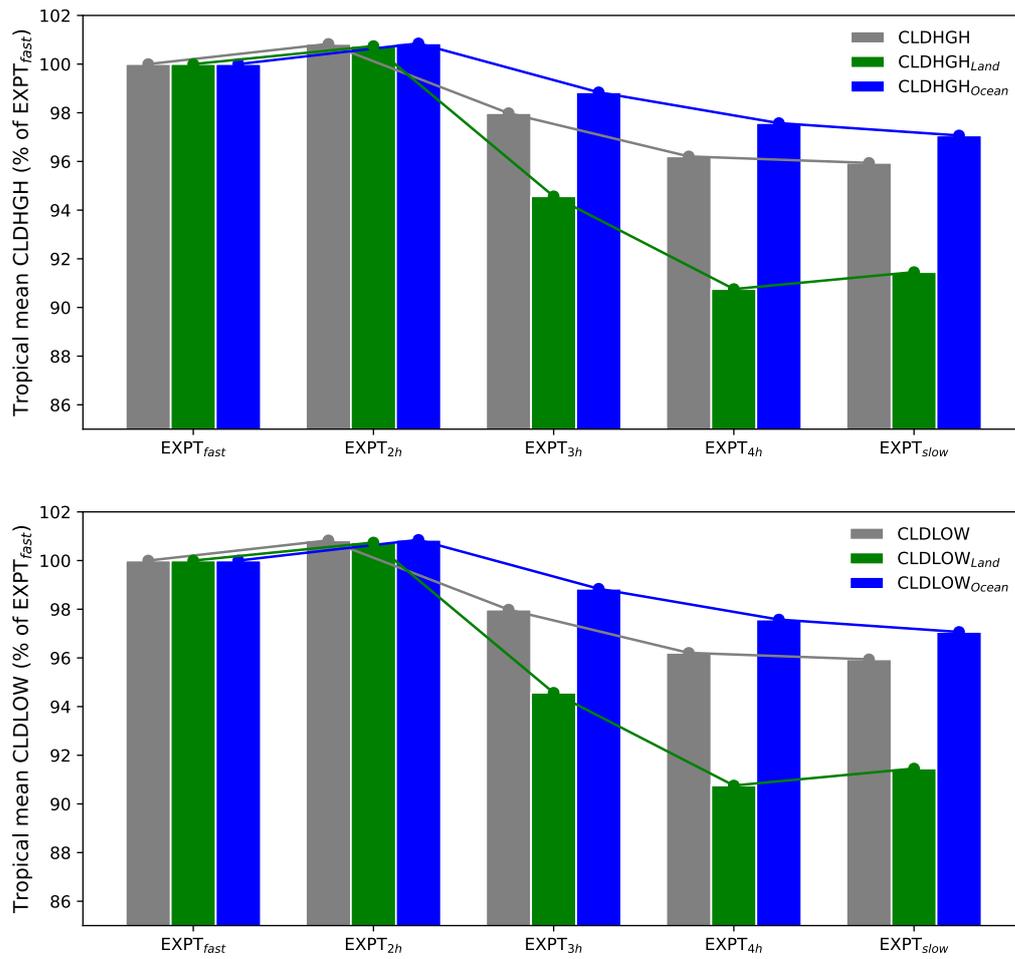
472 **Figure 1.** Tropical (tropics defined as the zonal belt between 30°S-30°N) annual mean daily rainfall
 473 (mm/day) for different experiments mentioned in Table 1.



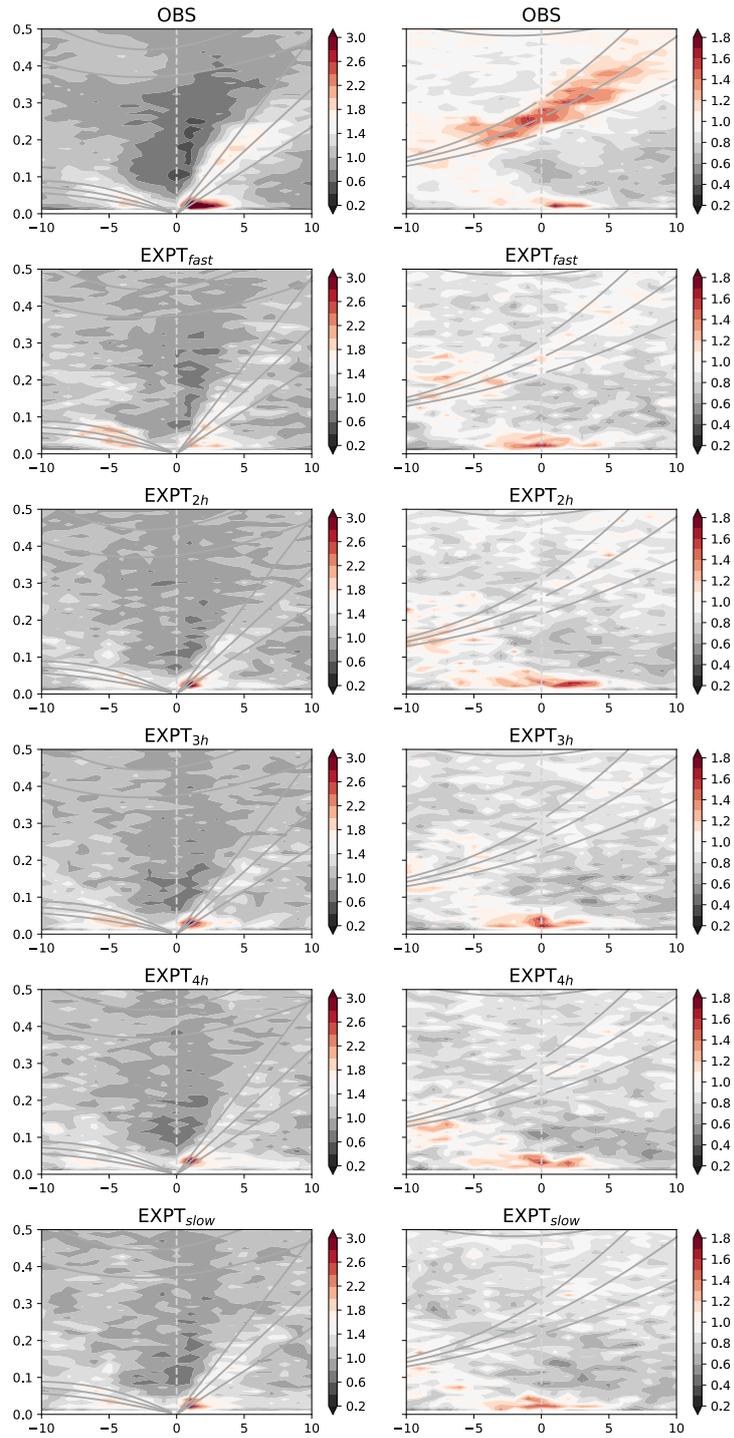
474 **Figure 2.** Tropical (tropics defined as the zonal belt between 30°S-30°N) mean vertical profiles of tem-
 475 perature (T) and specific humidity (Q). Departures of different experiments, as indicated in the legends, from
 476 $EXPT_{fast}$ (Land: Dotted, Ocean: Solid). The vertical dashed line indicate the zero line.



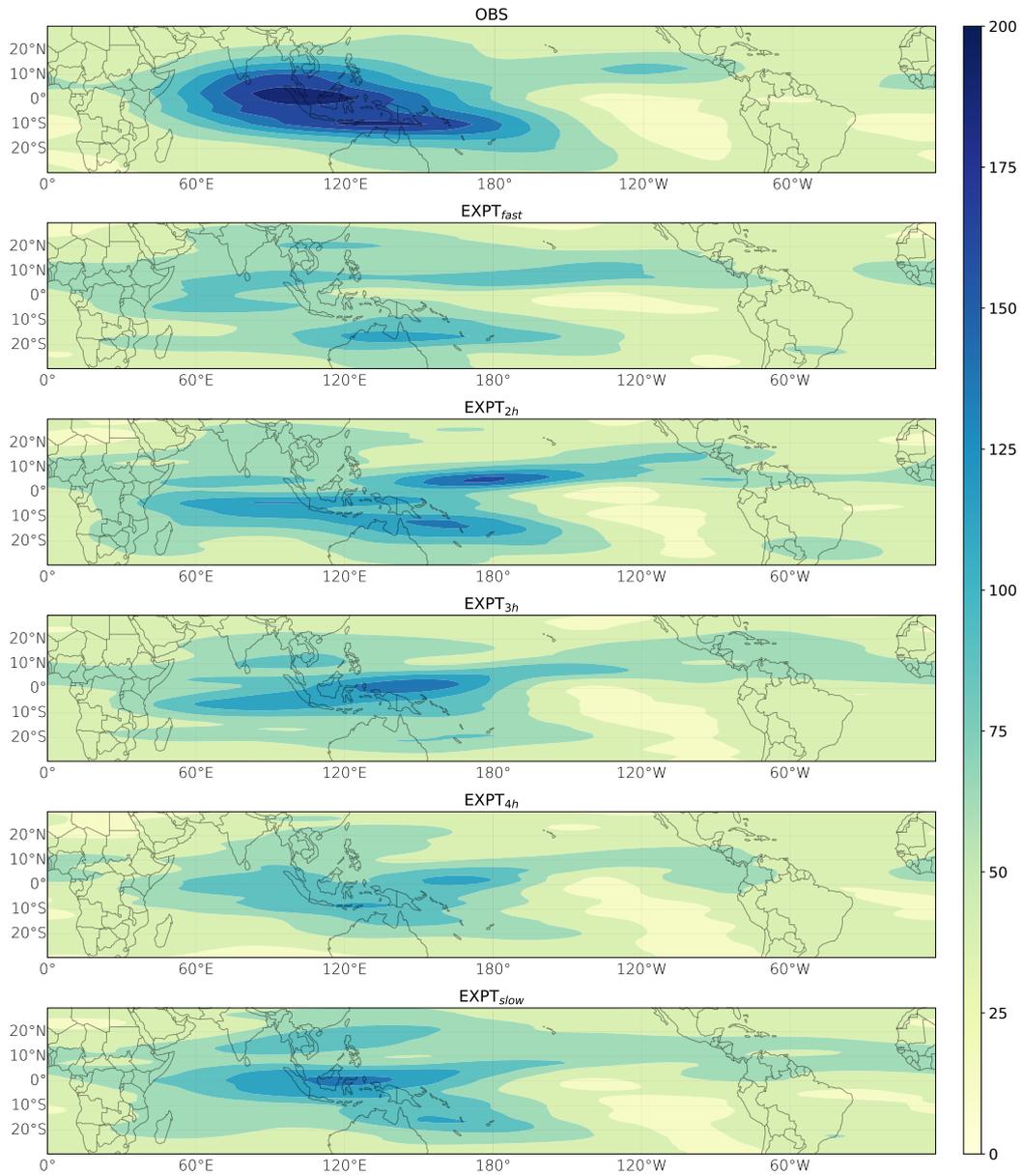
477 **Figure 3.** Tropical (tropics defined as the zonal belt between 30°S-30°N) annual daily mean specific hu-
 478 midity as surface depicted as % of $EXPT_{fast}$.



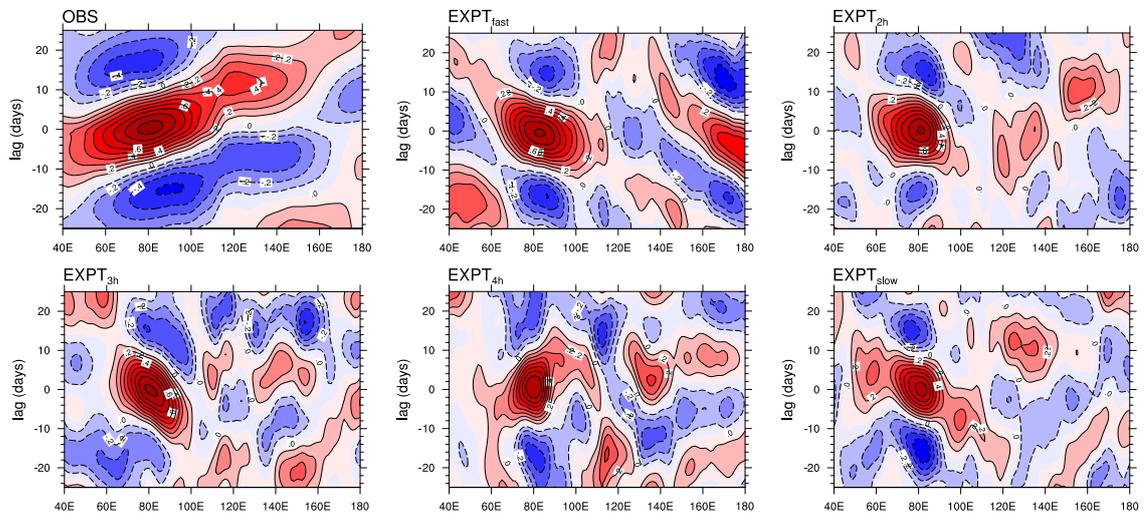
479 **Figure 4.** Tropical (tropics defined as the zonal belt between 30°S-30°N) annual daily mean High and Low
 480 cloud cover depicted as % of $EXPT_{fast}$.



481 **Figure 5.** Takayabu-Wheeler-Kiladis spectra of OLR for OBS (from NOAA) and different experiments (as
 482 named above each panel), for the symmetric component (left-hand side panels) and antisymmetric component
 483 (right-hand side panels).



484 **Figure 6.** MJO variance computed as the daily variance of OLR data filtered for 1-5 wavenumber and 20-
 485 100 day frequency, for OBS (from NOAA) and different experiments (as named above each panel).



486 **Figure 7.** MJO propagation: Hovmoller (averaged from 10°S to 10°N) plots of MJO-filtered OLR (W m⁻²)
487 anomalies (Winter), for OBS (from NOAA) and different experiments (as named above each panel).

Supplementary Materials for
An assessment of representing land-ocean
heterogeneity via convective adjustment
timescale in the Community Atmospheric Model
6 (CAM6)

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This PDF file contains:

- Figure S1. Annual tropical daily mean rainfall (mm/day).
- Figure S2. Tropical mean vertical profiles of DTCOND and DCQ.
- Figure S3a. Annual daily mean surface specific humidity (g/kg).
- Figure S3b. Annual daily mean mid-low-level specific humidity (g/kg).
- Figure S4. MJO variance.

Supplementary Materials

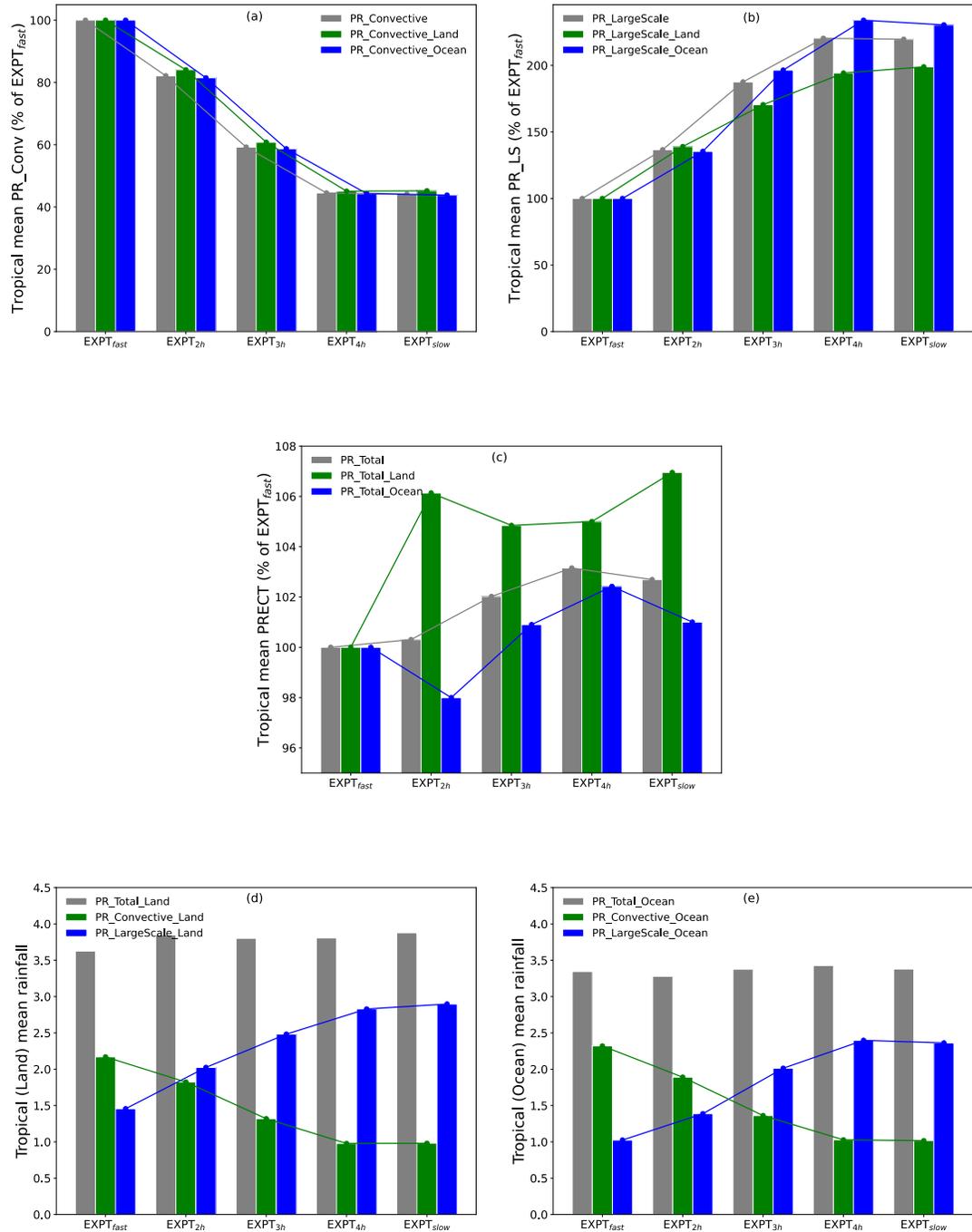


Figure S1. Annual tropical daily mean rainfall (mm/day). Same as Figure 1 of the main manuscript, except the Total, Convective and Large-scale rainfalls are plotted for land and ocean in addition to their total over the whole tropics.

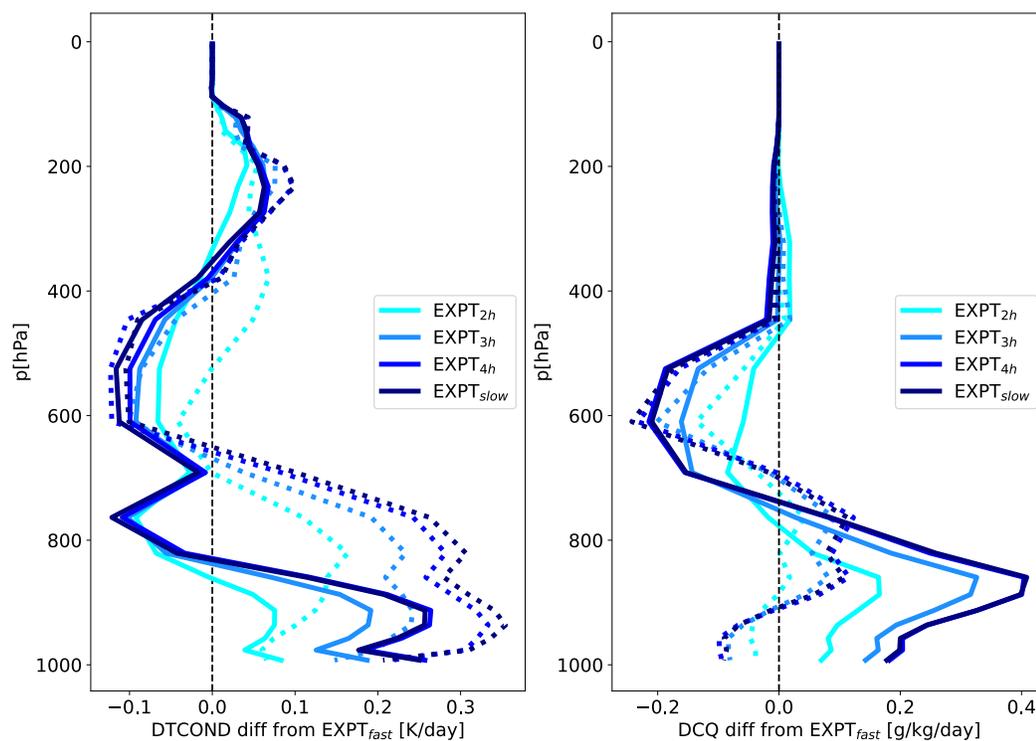


Figure S2. Tropical mean vertical profiles of temperature tendency due to moist processes (DTCOND) and specific humidity tendency due to moist processes (DCQ). Same as Figure 2 of the main manuscript, except for DTCOND and DCQ.

Supplementary Materials

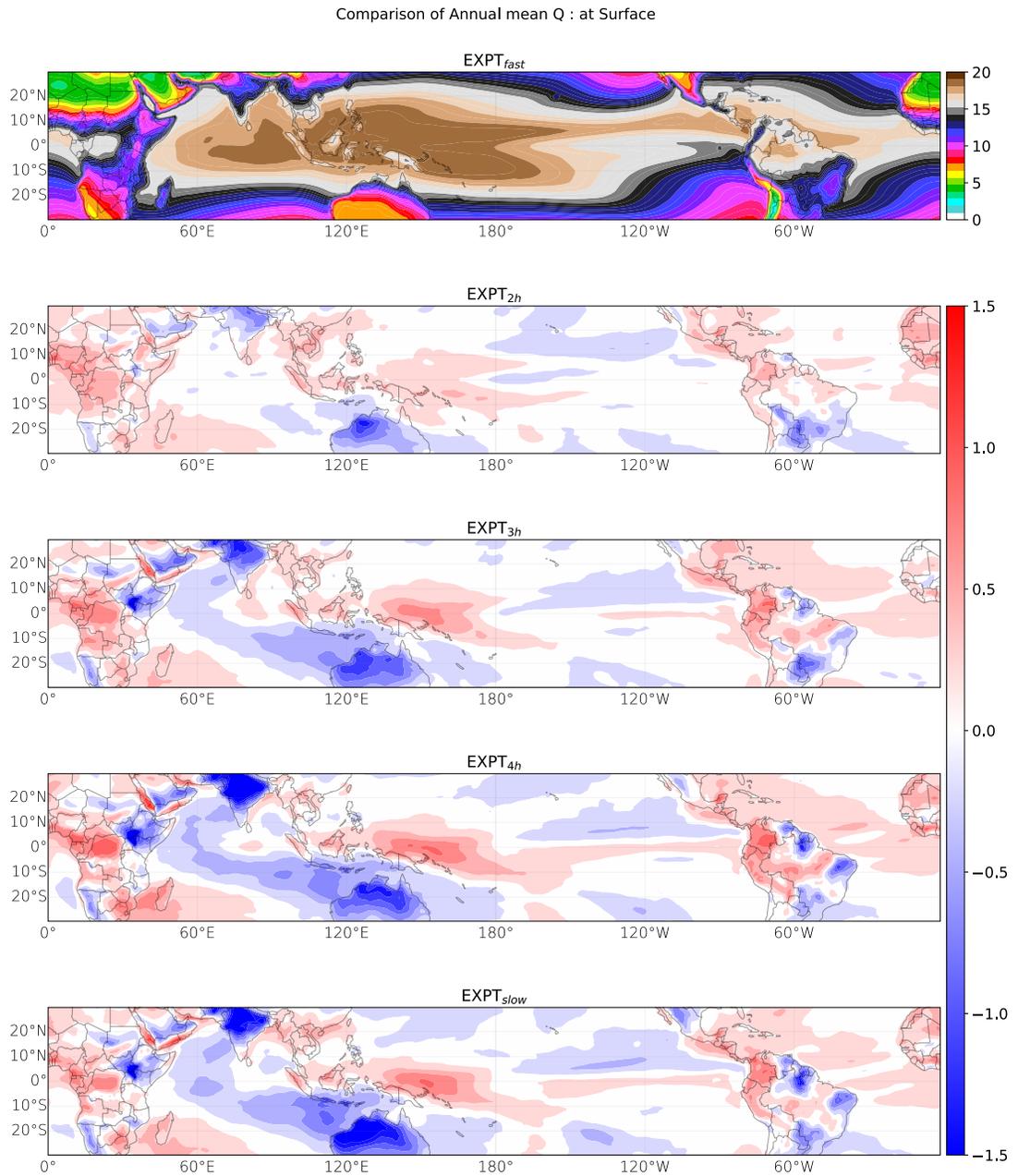


Figure S3a. Annual daily mean surface specific humidity (g/kg). Top panel shows the absolute values for $EXTP_{fast}$ and the remaining panels show departures of other simulations, simulation names as indicated by the panel headings, from $EXTP_{fast}$.

Supplementary Materials

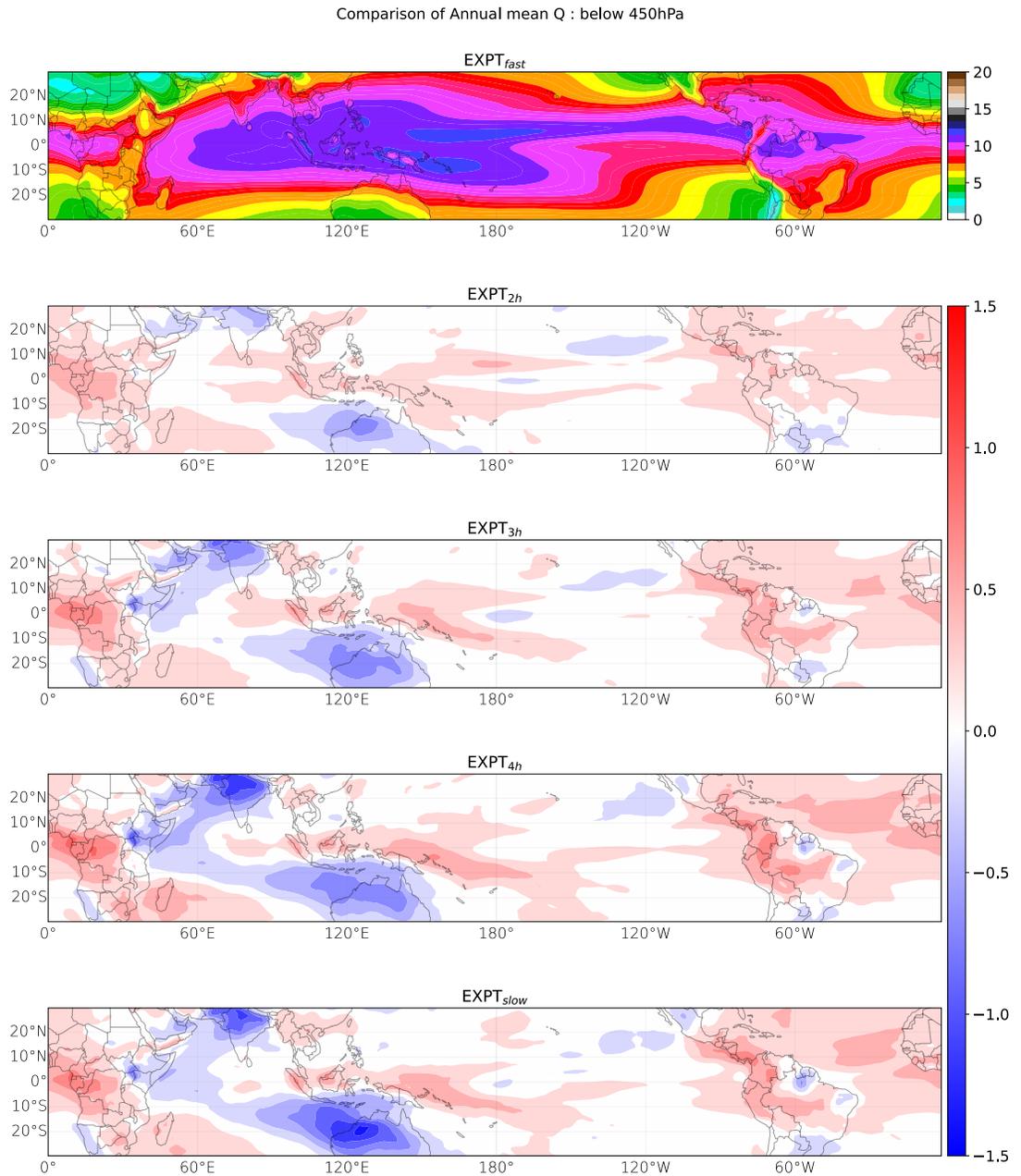


Figure S3b. Annual daily mean mid-low-level specific humidity (g/kg). Same as Figure S3a, except averaged over surface to 450hPa.

Supplementary Materials

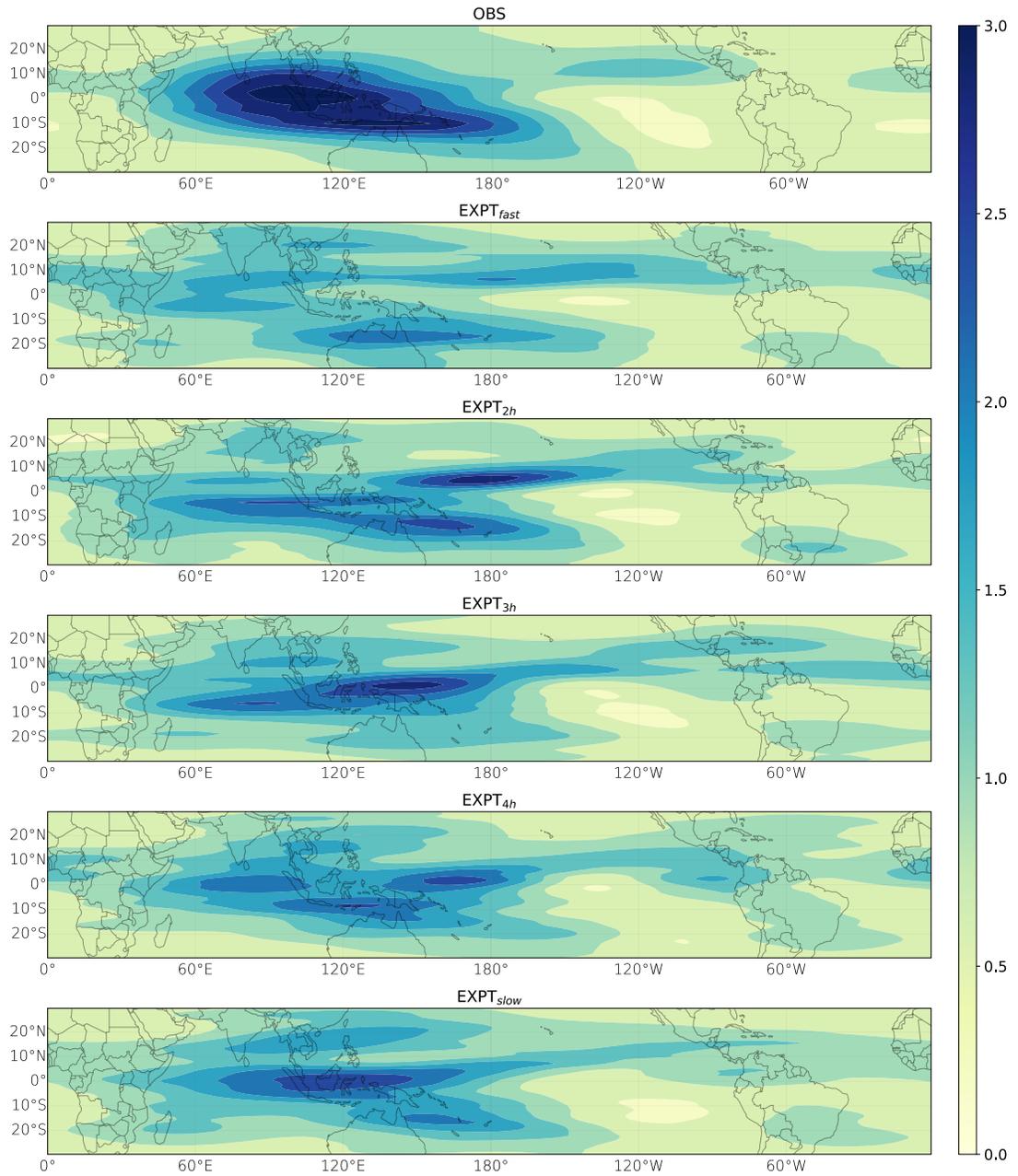


Figure S4. MJO variance. Same as Figure 6 of the main manuscript, except the variance fields are normalized by the respective domain means.